

The Proton Structure from HERA and the Tevatron

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Topics

- $F_2, F_L, F_2^{c\bar{c}}$
- High Q^2 Neutral & Charged Current Cross Sections
- Gluon Momentum Density
- Inclusive Jet Cross Section at all η and Dijet Mass
- Ratio of Jet Cross Sections at Different \sqrt{s}
- BFKL vs DGLAP

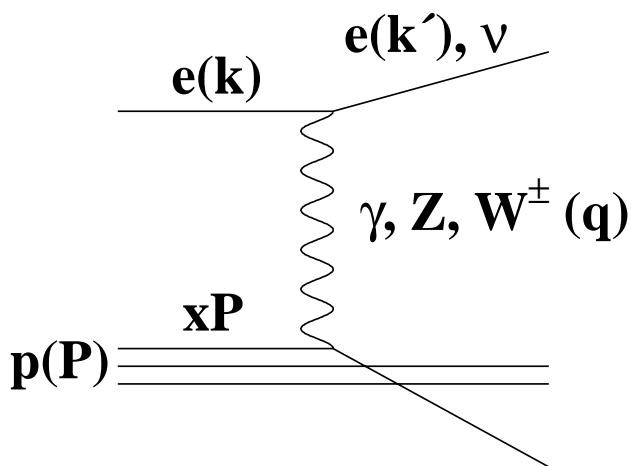
2000 Aspen Winter Conference on Particle Physics

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Some History

- powerful tool: Deep Inelastic Scattering of leptons off protons
- if virtual γ carries “large” 4-momentum transfer,
it can probe internal p structure
- Structure Functions scale with momentum transfer,
similarly to e scattering off μ
 \Rightarrow *pointlike constituents of p*
- 30 years of DIS experiments \rightarrow
impressive understanding of p structure
- complimentary: $p\bar{p}$ collisions
- p structure understanding crucial for (new) physics
in current and future colliders

Deep Inelastic Scattering at HERA



$$Q^2 = -q^2 = -(k - k')^2$$

(4-momentum transfer)²

$$x = Q^2 / (2P \cdot q)$$

parton momentum fraction

$$y = P \cdot q / (P \cdot k)$$

fractional energy transfer

$$\frac{d^2\sigma^{NC}(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2^{NC}(x, Q^2) - y^2 F_L^{NC}(x, Q^2) \mp Y_- x F_3^{NC}(x, Q^2)]$$

$$F_2^{NC}(x, Q^2) = \sum_{i=1}^{n_f} A_i x [q_i(x, Q^2) + \bar{q}_i(x, Q^2)]$$

$$xF_3^{NC}(x, Q^2) = \sum_{i=1}^{n_f} B_i x [q_i(x, Q^2) - \bar{q}_i(x, Q^2)]$$

A_i, B_i : electroweak quark couplings

$xq(x, Q^2), x\bar{q}(x, Q^2)$: momentum distributions of (anti)quarks

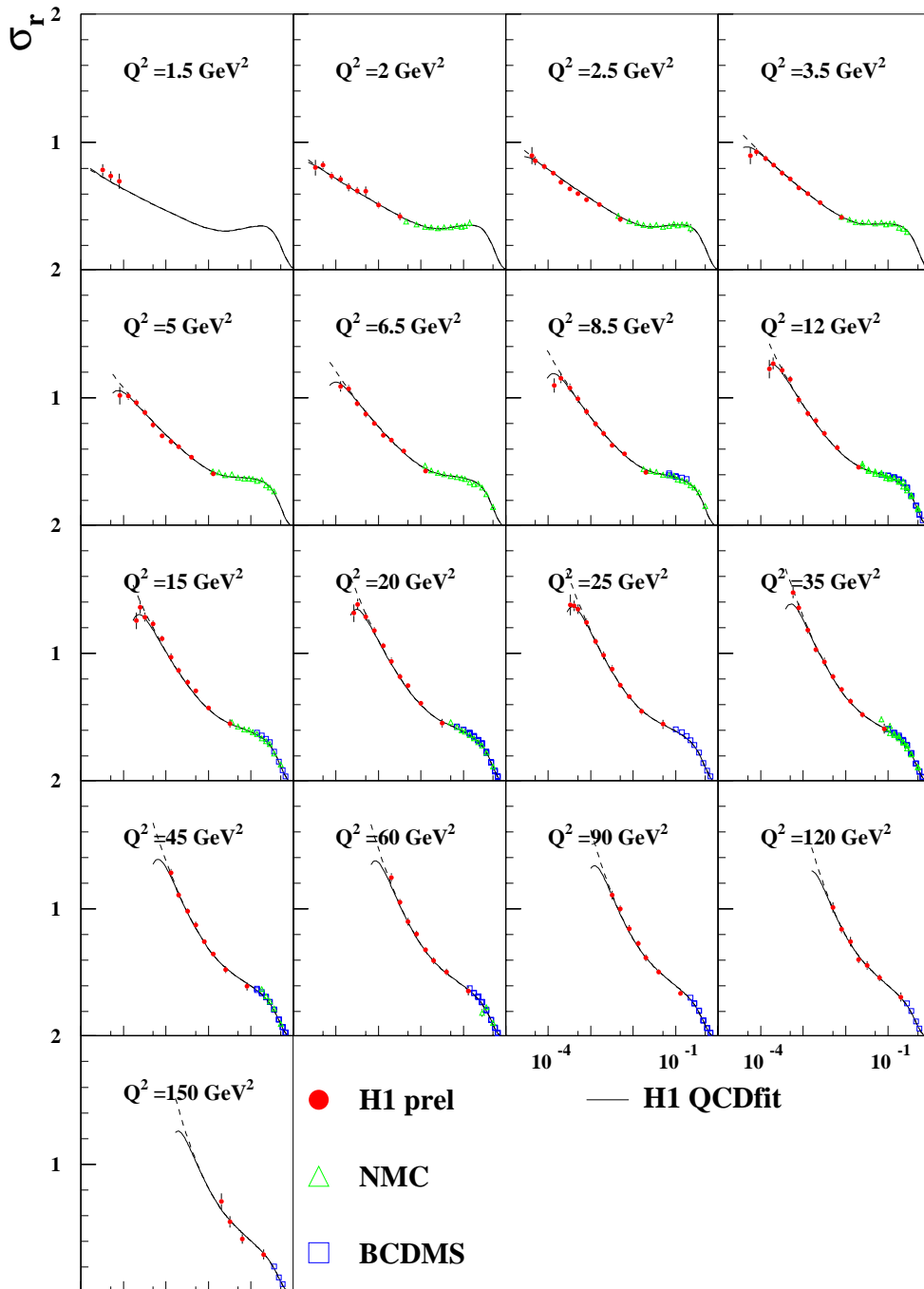
$F_L(x, Q^2)$: longitudinal S.F. (zero at LO, a few % at NLO)

$$Y_\pm = 1 \pm (1 - y)^2$$

Cross Section at Moderate Q^2

$$\sigma_r(x, Q^2) = \frac{d^2\sigma}{dx dQ^2} \cdot \frac{xQ^4}{2\pi\alpha^2} \cdot \frac{1}{Y_+} = F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2)$$

H1 96-97



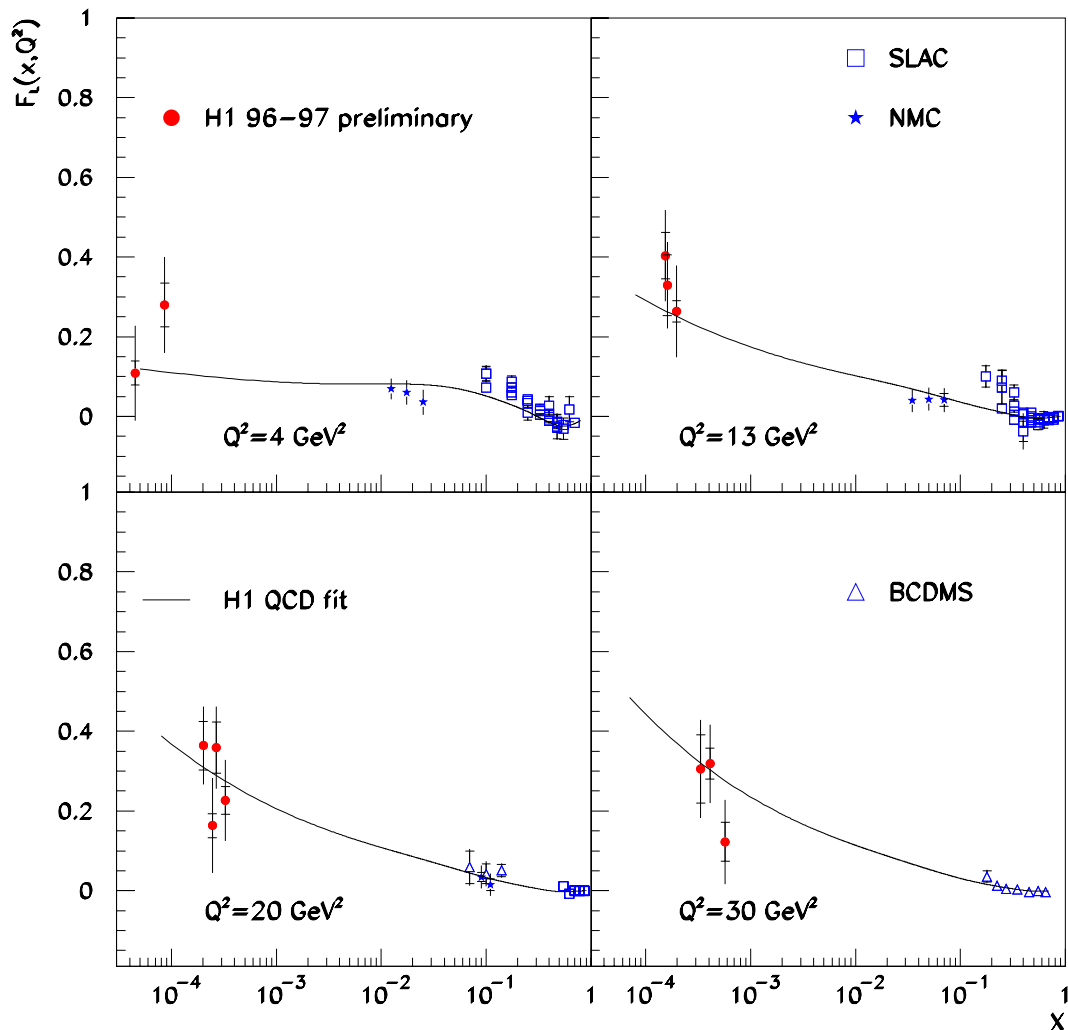
σ_r rise with decreasing x persists down to $Q^2 \approx 1 \text{ GeV}^2$

pQCD describes the data at $Q^2 \geq 1 \text{ GeV}^2$

$F_L(x, Q^2)$

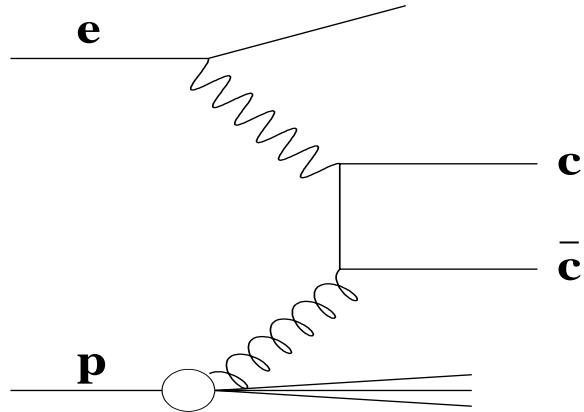
$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \cdot \{[1 + (1 - y)^2]F_2(x, Q^2) - y^2 F_L(x, Q^2)\}$$

- at high y , $[1 + (1 - y)^2]$ and y^2 comparable
- QCD fit to F_2 at low $y \rightarrow$ extract PDF's \rightarrow evolve in $Q^2 \rightarrow$ predict F_2 at high y
- subtract F_2 from $d^2\sigma/dx dQ^2$

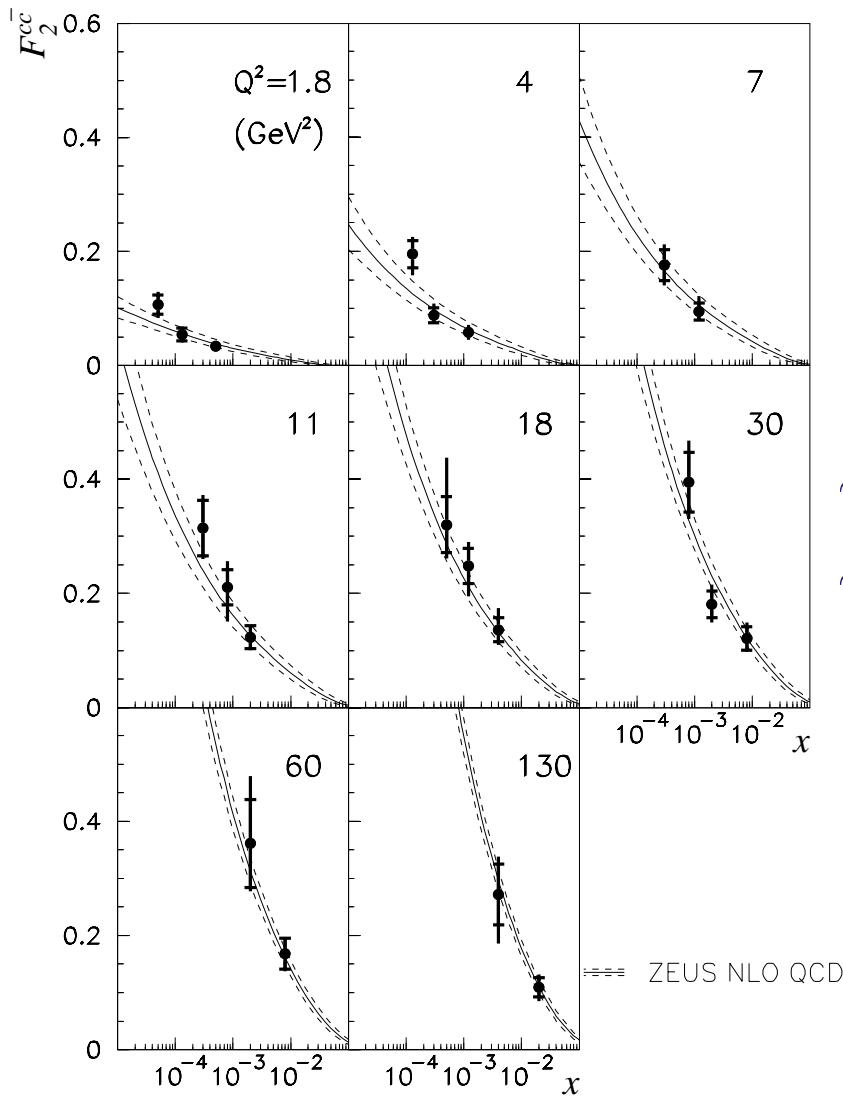


$$F_2^{c\bar{c}}(x, Q^2)$$

Boson-Gluon Fusion



ZEUS 1996–97

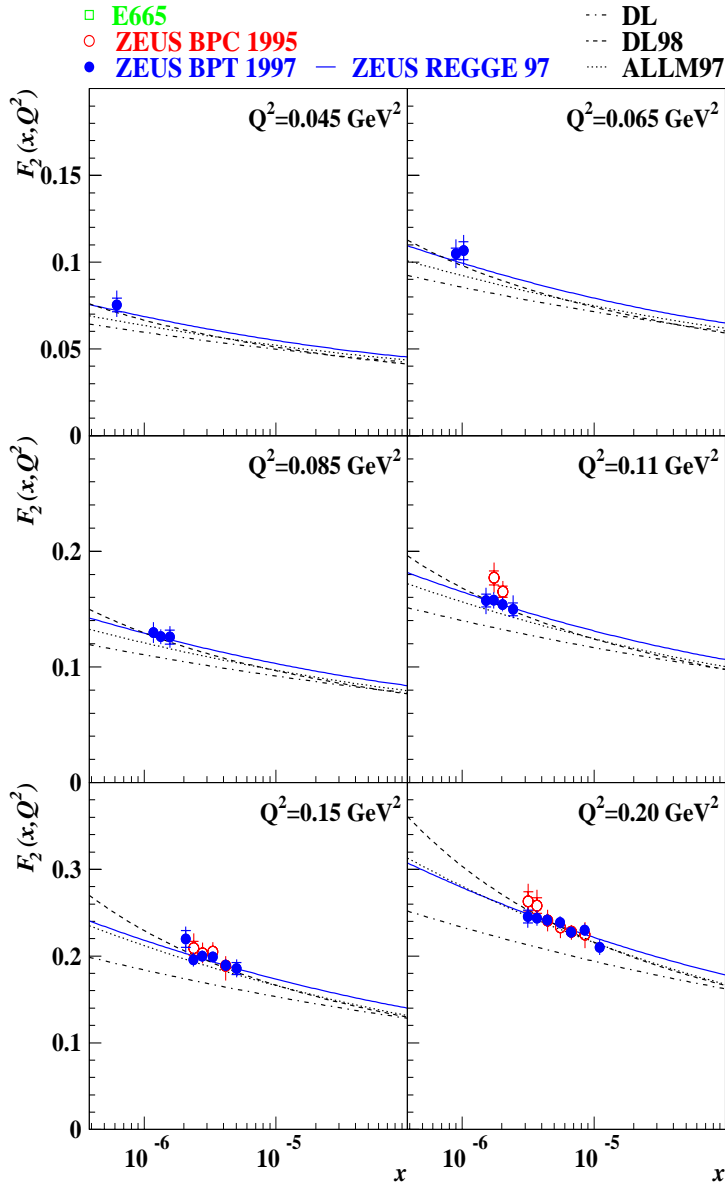


- $F_2^{c\bar{c}} \uparrow$ when $x \downarrow$
faster than F_2
(gluonic origin of c)

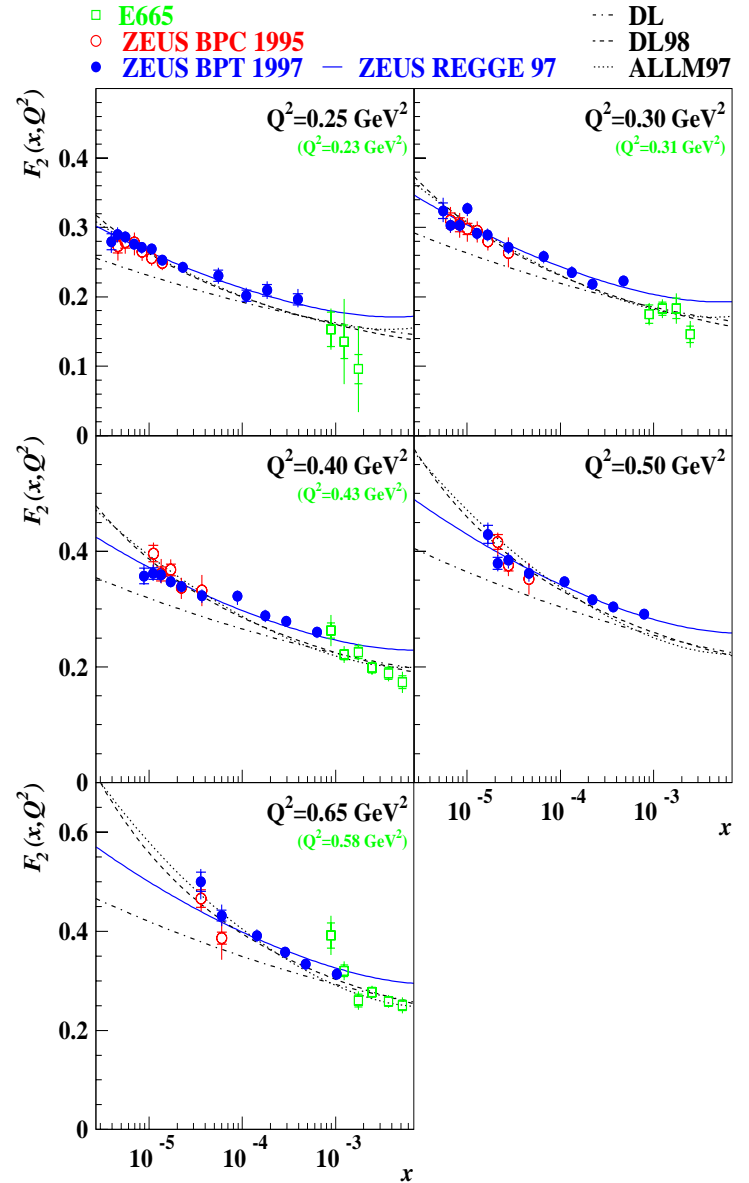
- $F_2^{c\bar{c}}/F_2$ is
 $\sim 10\%$ at $Q^2 = 1.8 \text{ GeV}^2$
 $\sim 30\%$ at $Q^2 = 130 \text{ GeV}^2$
 for $10^{-4} < x < 10^{-3}$

$F_2(x, Q^2)$ at very low Q^2

ZEUS 1997 (Preliminary)



ZEUS 1997 (Preliminary)

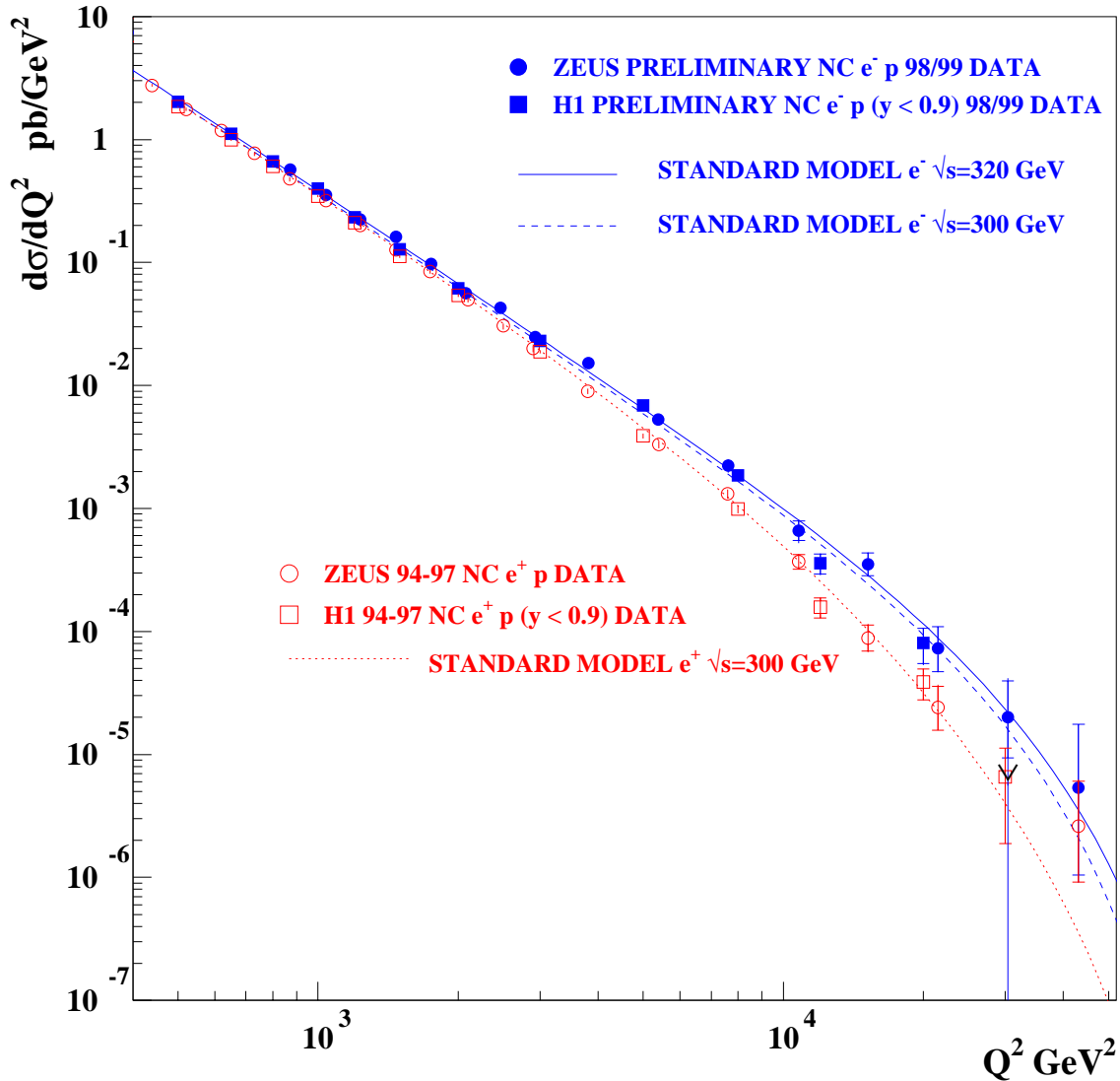


- F_2 still rises at $Q^2 \approx 0.1 \text{ GeV}^2$

High Q^2 Neutral Current

$$\frac{d^2\sigma^{NC}(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2(x, Q^2) - y^2 F_L(x, Q^2) \mp Y_- x F_3(x, Q^2)]$$

H1 + ZEUS NC DATA

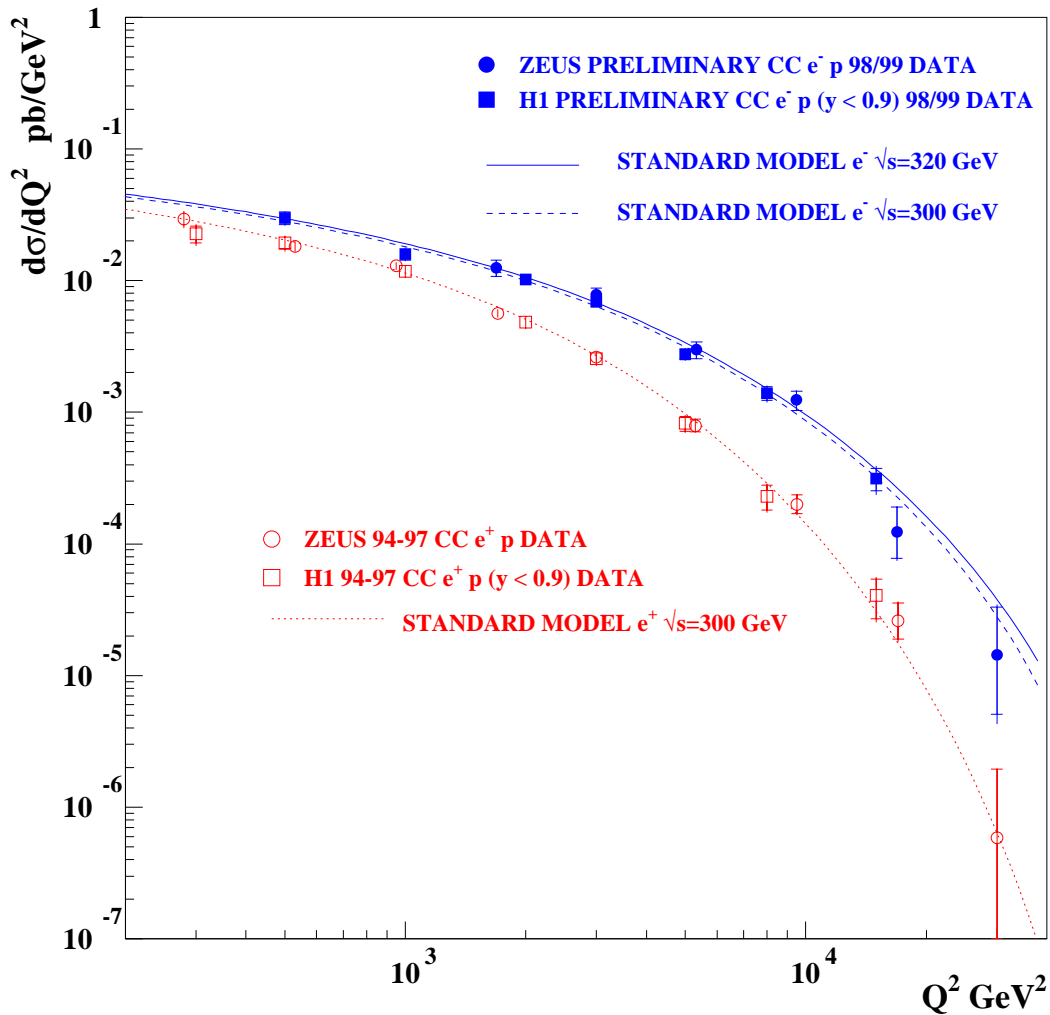


- SM predictions agree with data
- Parity-violating contribution from Z exchange visible for $Q^2 \geq M_Z^2$

High Q^2 Charged Current

$$\frac{d^2\sigma^{CC}(e^\pm p)}{dx dQ^2} = \frac{G_F^2}{2\pi} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \begin{array}{l} [\bar{u} + \bar{c} + (1 - y^2)(d + s)] \\ [u + c + (1 - y^2)(\bar{d} + \bar{s})] \end{array}$$

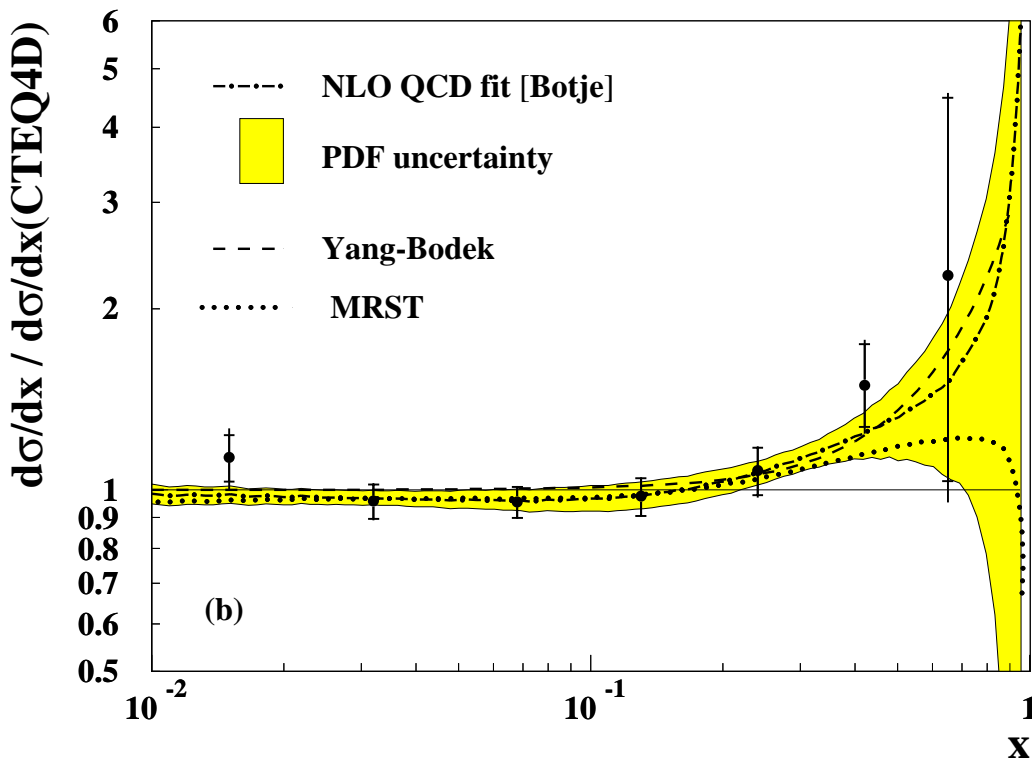
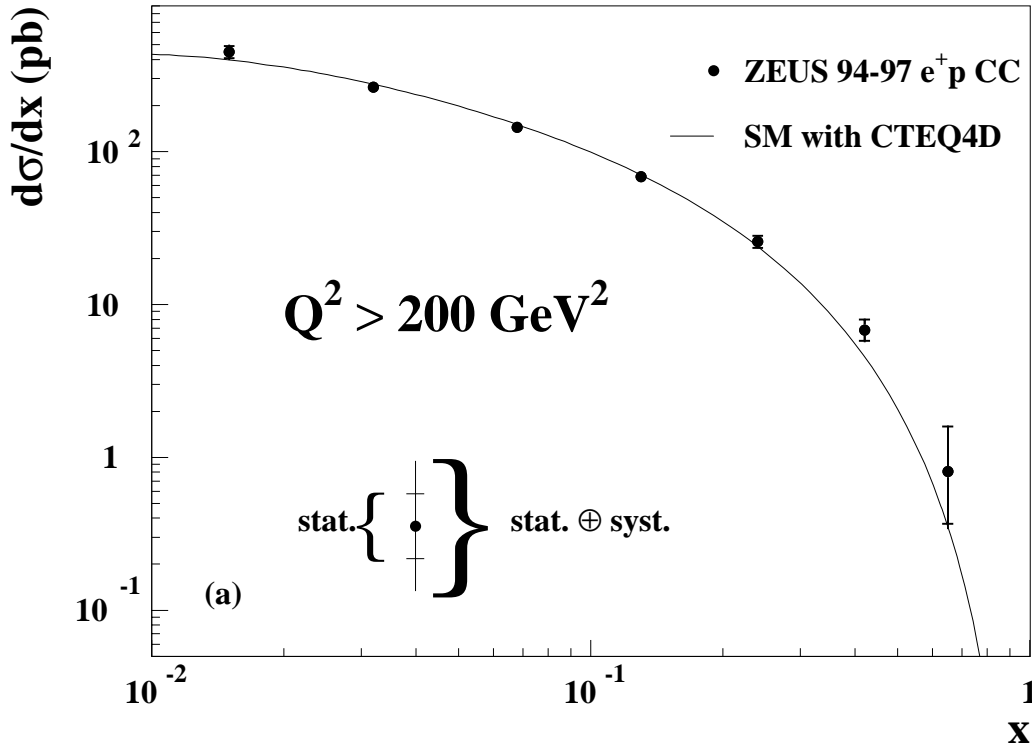
H1 + ZEUS CC DATA



- SM predictions agree with data
- $e^- p$ dominated by u , $e^+ p$ dominated by $(1 - y)^2 d$
- $M_W = 80.8_{-4.5}^{+4.9}(\text{stat})_{-4.3}^{+5.0}(\text{syst})_{-1.3}^{+1.4}(\text{PDF})$ GeV
agrees w/ direct determination of M_W at Tevatron

High Q^2 CC (and the d/u ratio)

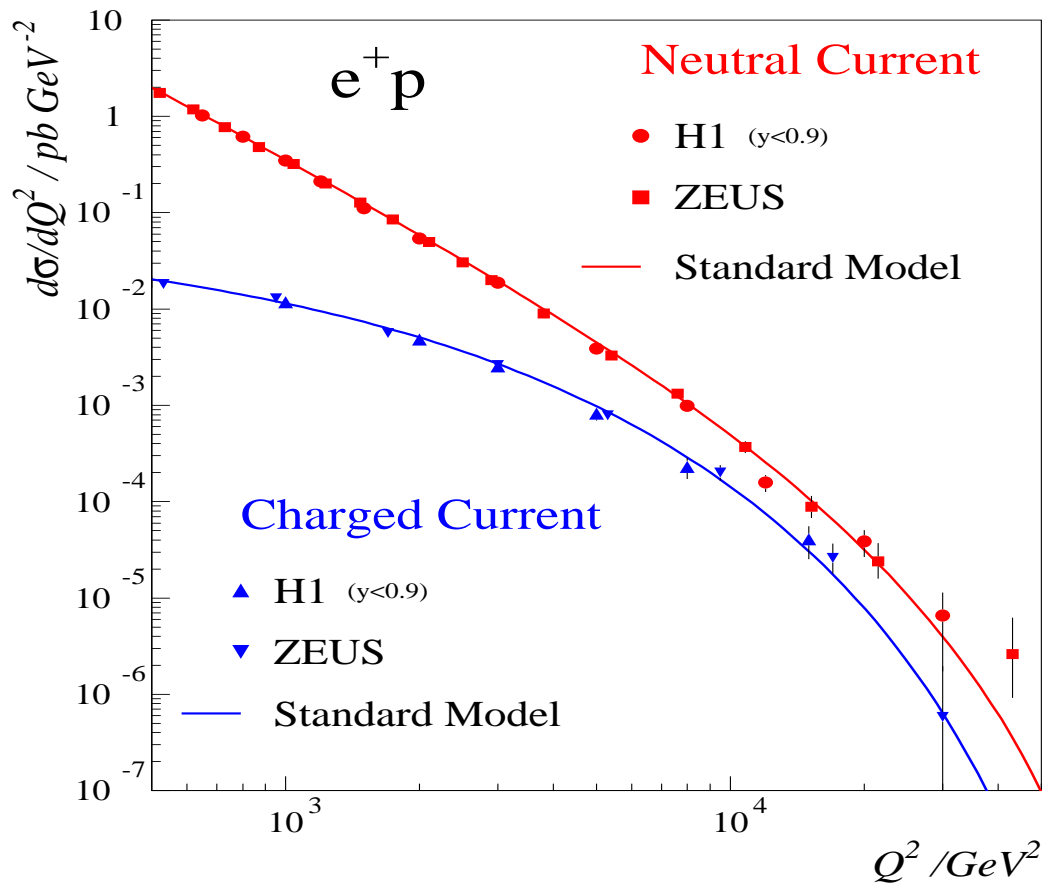
ZEUS CC 1994-97



High Q^2 NC and CC (from e^+p)

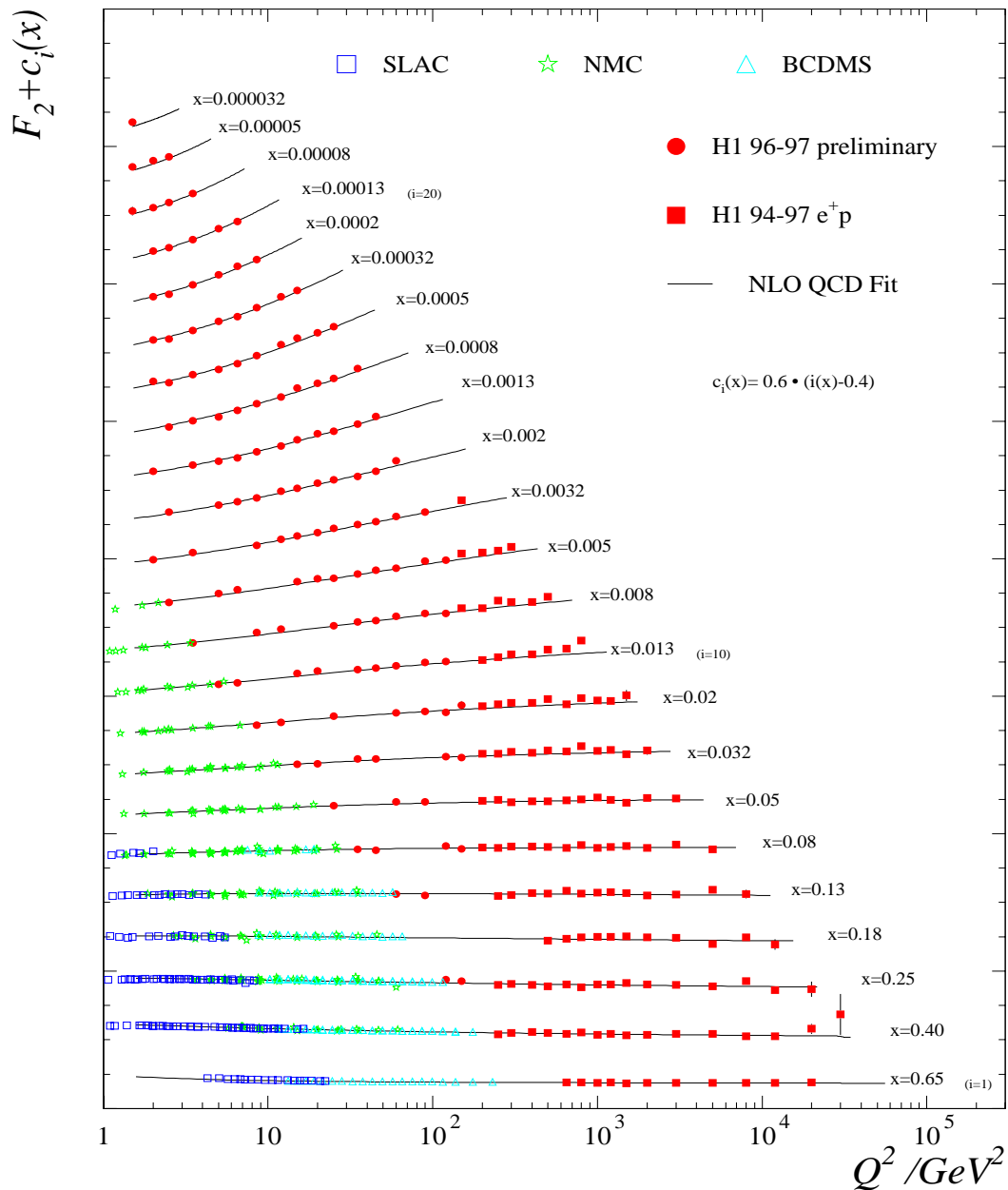
$$\frac{d^2\sigma(e^+p)}{dx dQ^2} = \mathcal{P} [Y_+ F_2(x, Q^2) - y^2 F_L(x, Q^2) - Y_- x F_3(x, Q^2)]$$

$$\mathcal{P}^{NC} = \frac{2\pi\alpha^2}{xQ^4}, \quad \mathcal{P}^{NC} = \frac{G_F^2}{4\pi x} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2$$



- EM force stronger than weak for $Q^2 < M_W^2, M_Z^2$
- EM & Weak forces comparable for $Q^2 \geq M_W^2, M_Z^2$

$F_2(x, Q^2)$: Scaling Violations

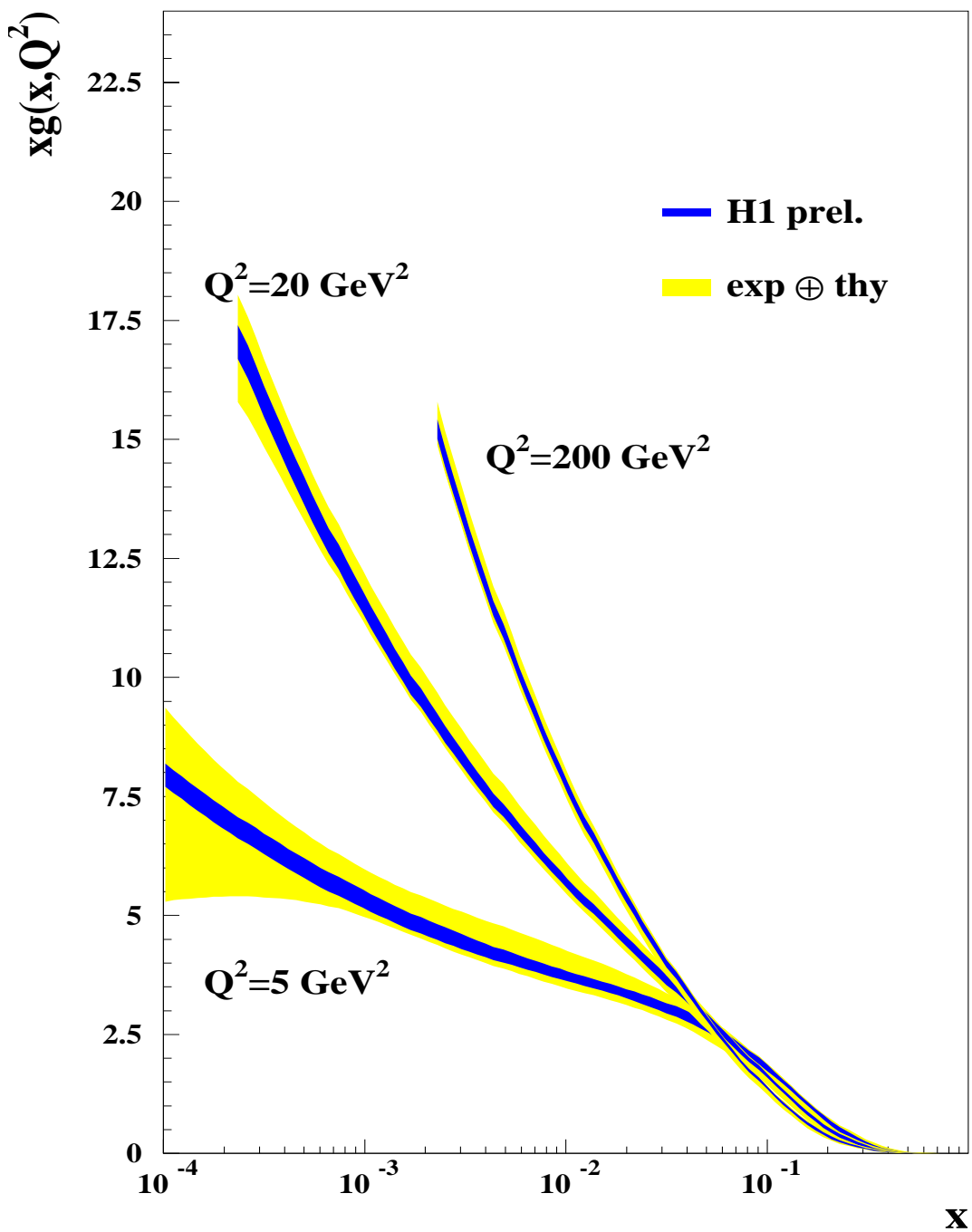


- due to quark-gluon interactions, $\propto \alpha_s \log Q^2$
- as $Q^2 \uparrow$: gluons and sea quarks \Rightarrow more partons of smaller x
- Large scaling violations described by NLO DGLAP for $3 \cdot 10^{-5} < x < 0.65$ and $1 < Q^2 < 3 \cdot 10^4 \text{ GeV}^2$

$xg(x, Q^2)$ from F_2 Scaling Violations

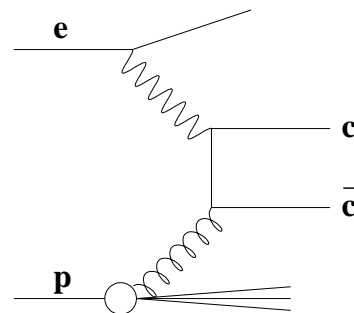
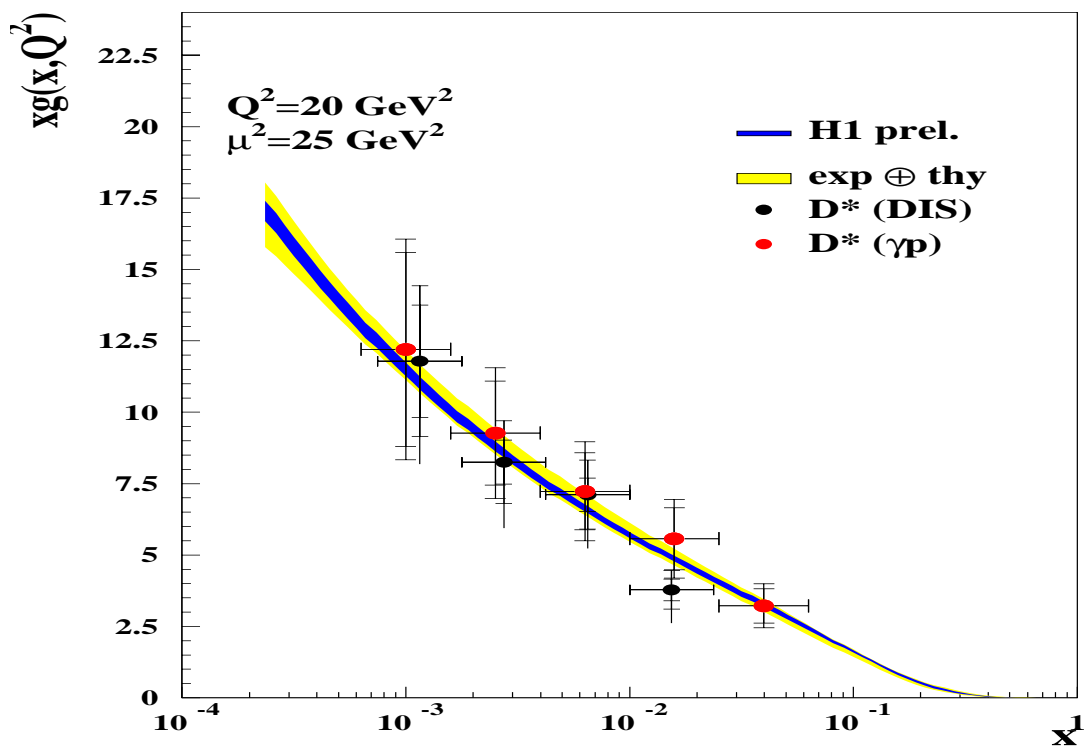
extract $xg(x, Q^2)$ from $dF_2/d\ln Q^2$

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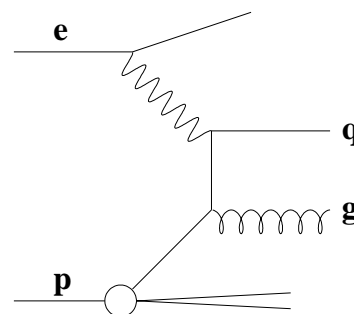
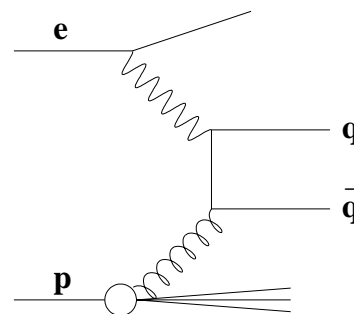
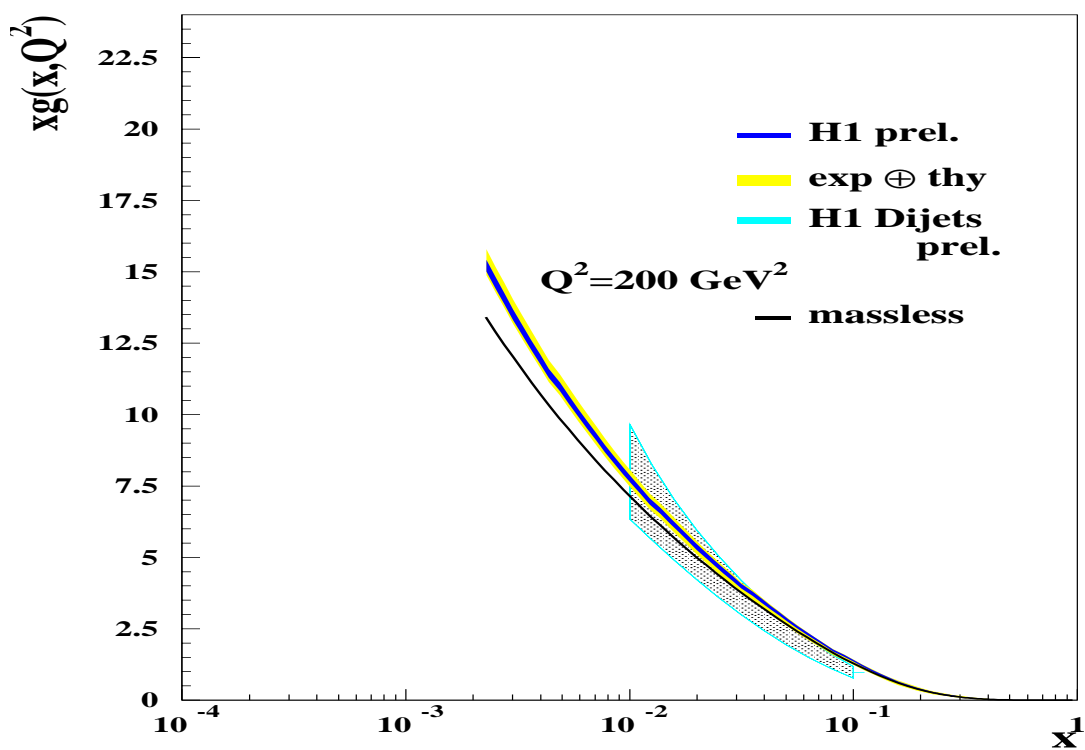


$xg(x, Q^2)$ from $F_2^{c\bar{c}}$ and Dijets

H1 96-97



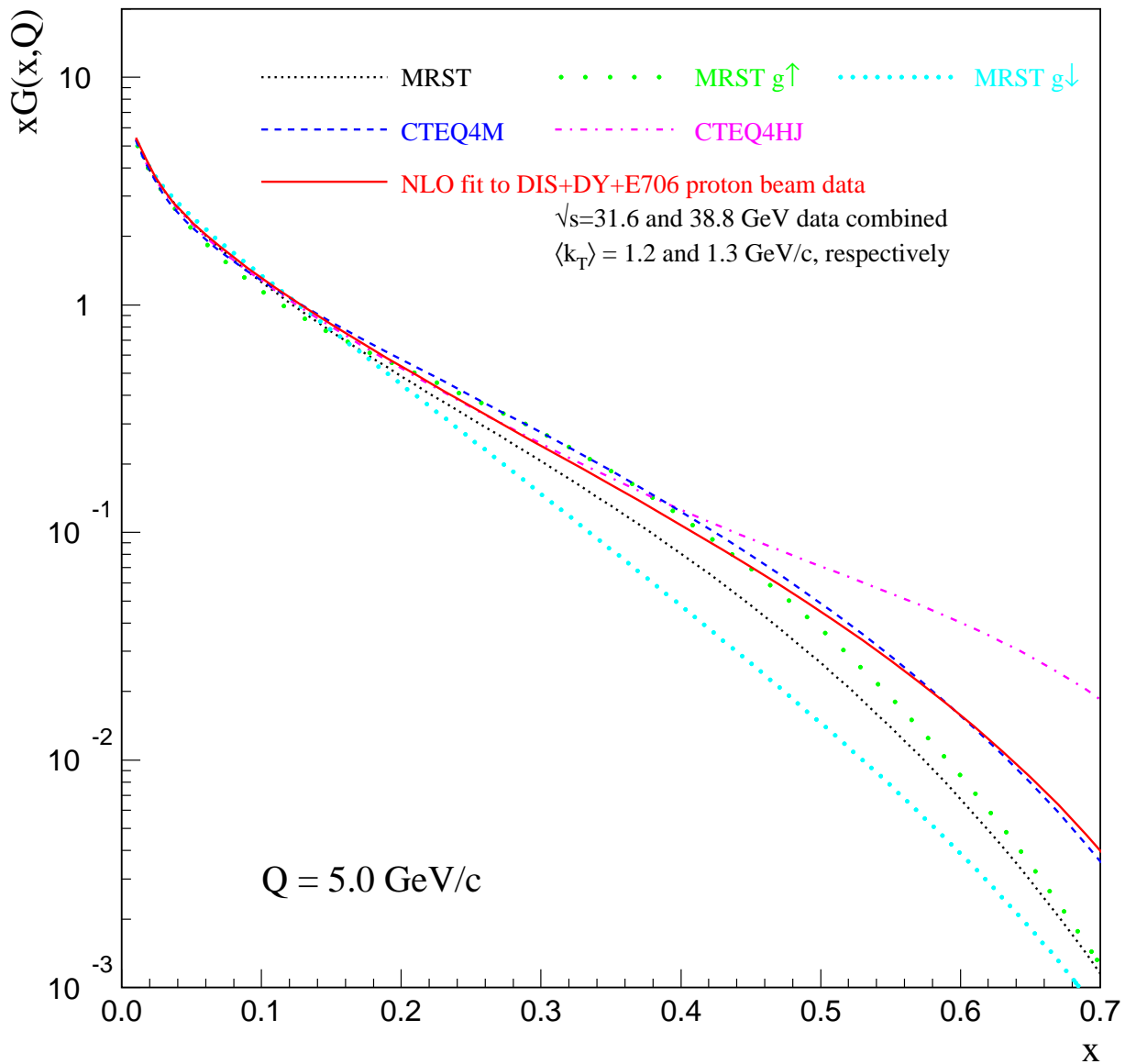
H1 96-97



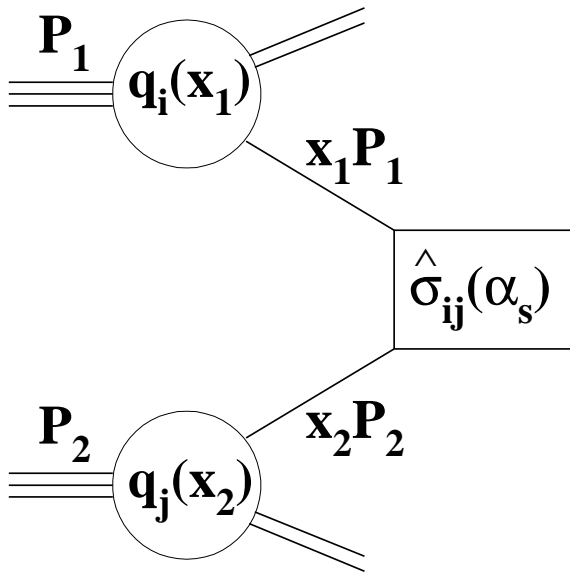
$xg(x, Q^2)$ at high x

MRST: direct photons from WA70,
range of k_T corrections from E706

CTEQ: inclusive jets from Tevatron



Jet Physics in $p\bar{p}$ collisions



$q_{i,j}$: parton densities

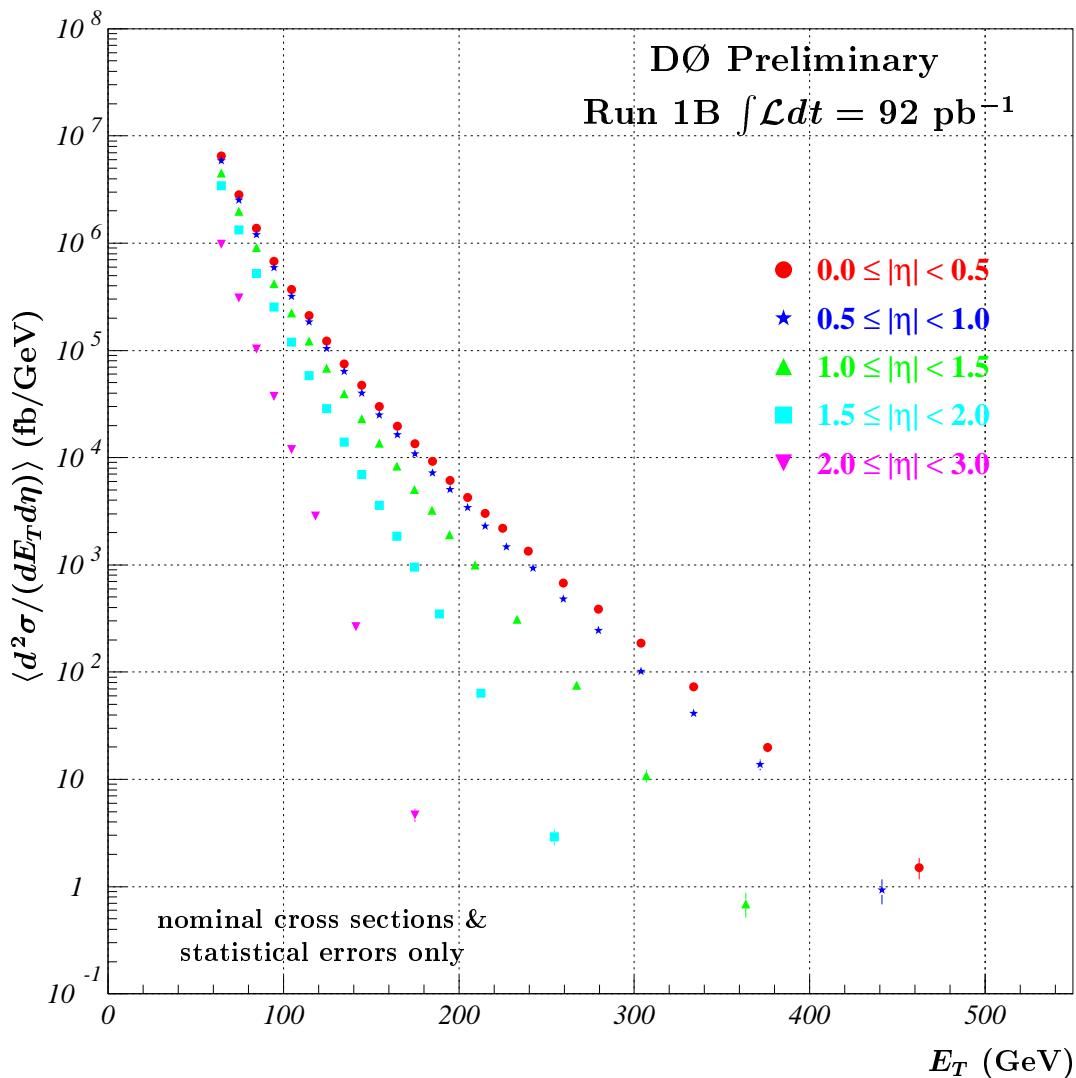
$\hat{\sigma}_{ij}$: partonic cross section

$\hat{\sigma}_{ij}(x_1 P_1, x_2 P_2, Q^2, \mu_F^2, \mu_R^2)$

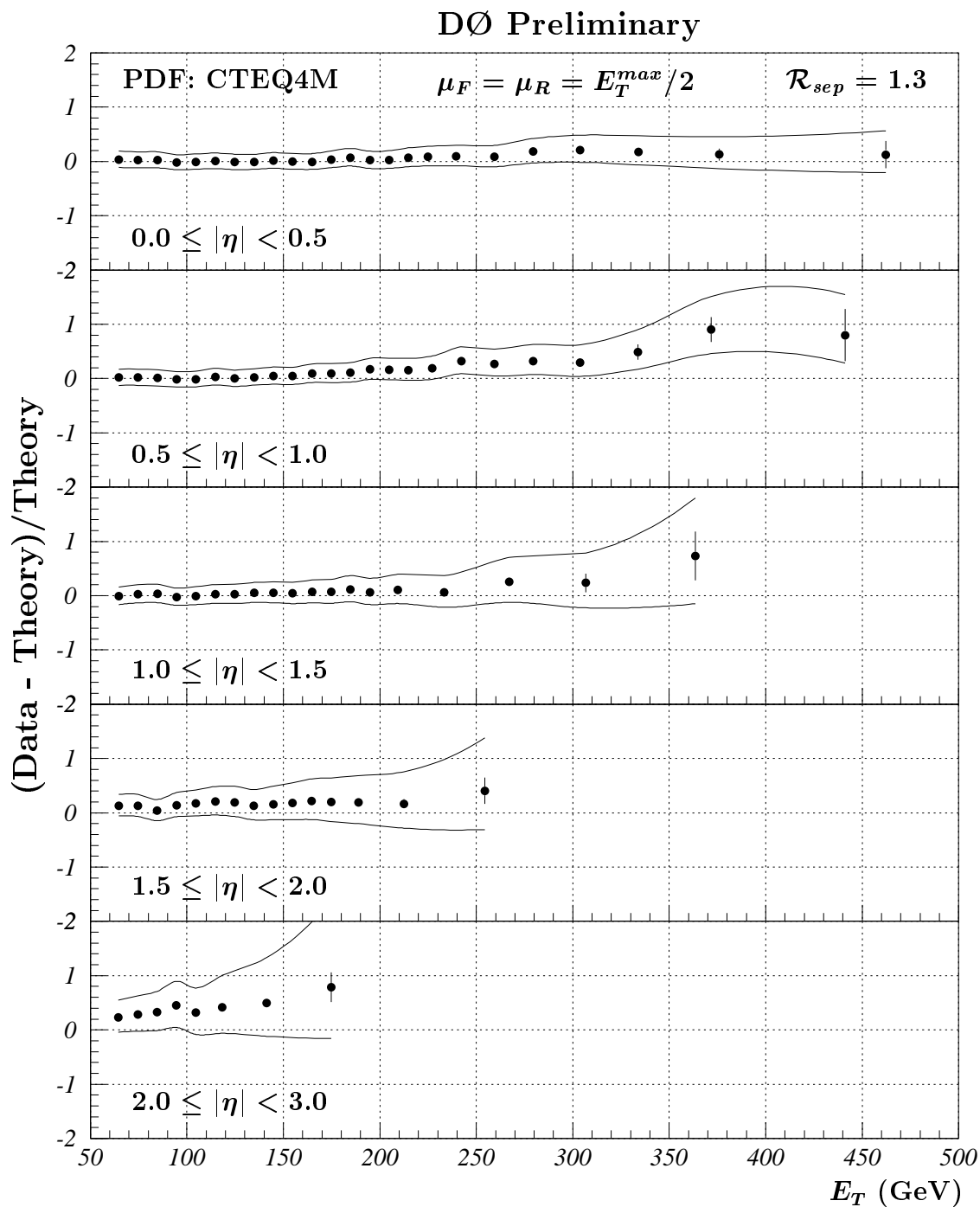
- Q : characteristic scale of the hard scattering
- Factorization scale μ_F : separates long- from short-distance
- Renormalization scale μ_R : to avoid divergent integrals in $\hat{\sigma}$
- scheme: $\mu_F = \mu_R = Q$
- Q of order jet- E_T
- probe perturbative QCD (hard scattering)
- test and use in PDF's
- search for quark substructure

Inclusive Jet Cross Section

- open question at high jet E_T 's
- new measurement at all rapidities \Rightarrow
better discriminating power:
 - 4x data points
 - correlations (χ^2)
 - forward $\eta \Rightarrow$ higher x

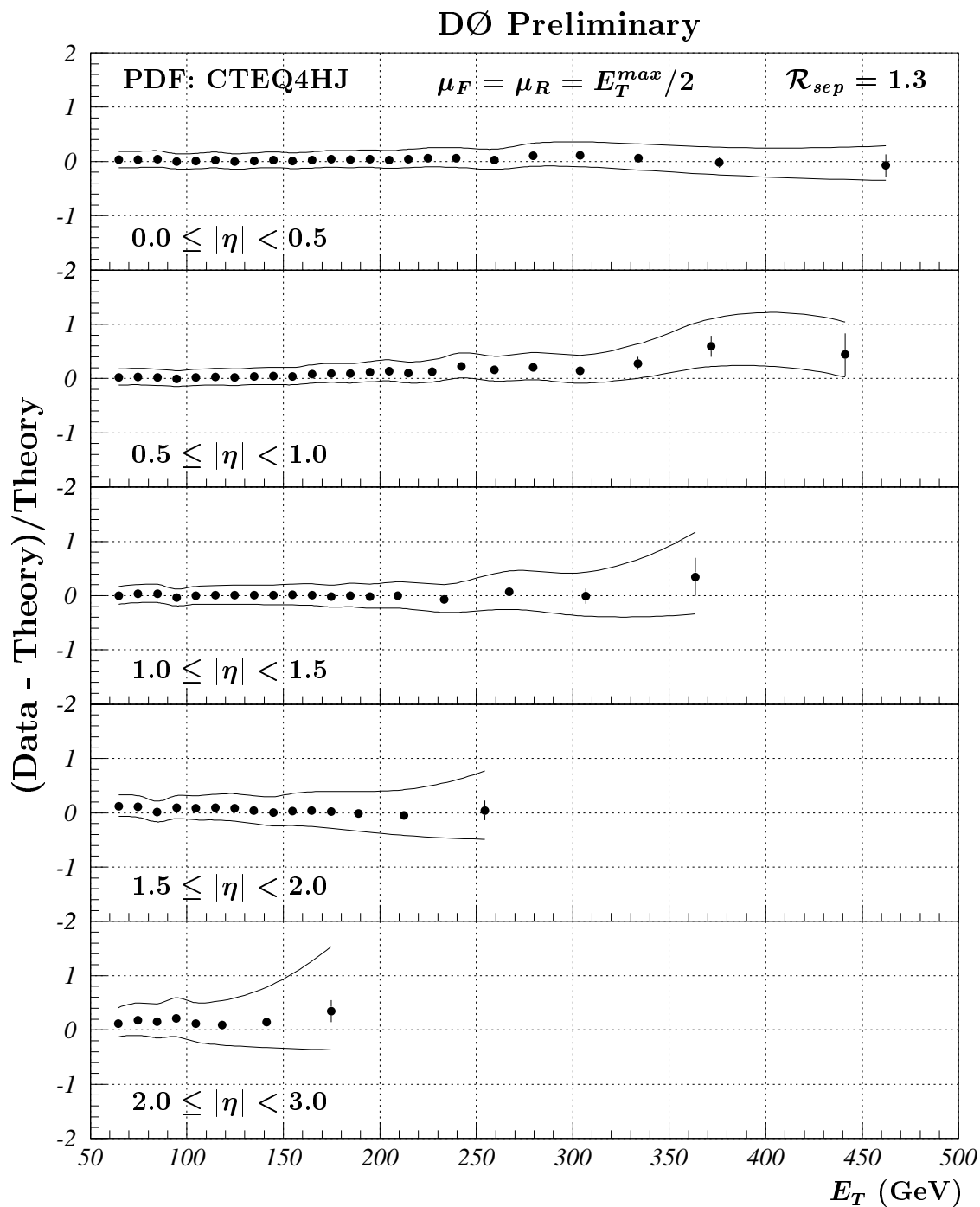


Inclusive Jet Cross Section at All Rapidities



- NLO QCD describes the data over seven orders of magnitude
- trend at higher E_T 's?

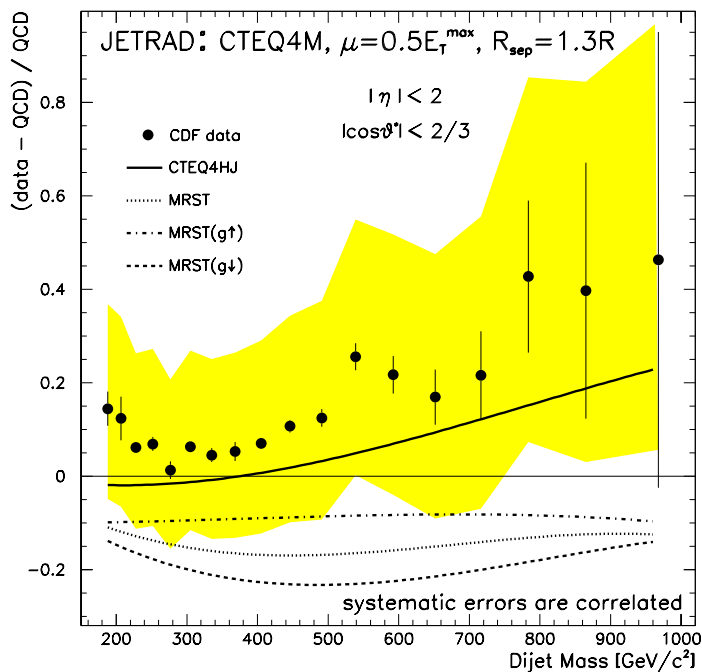
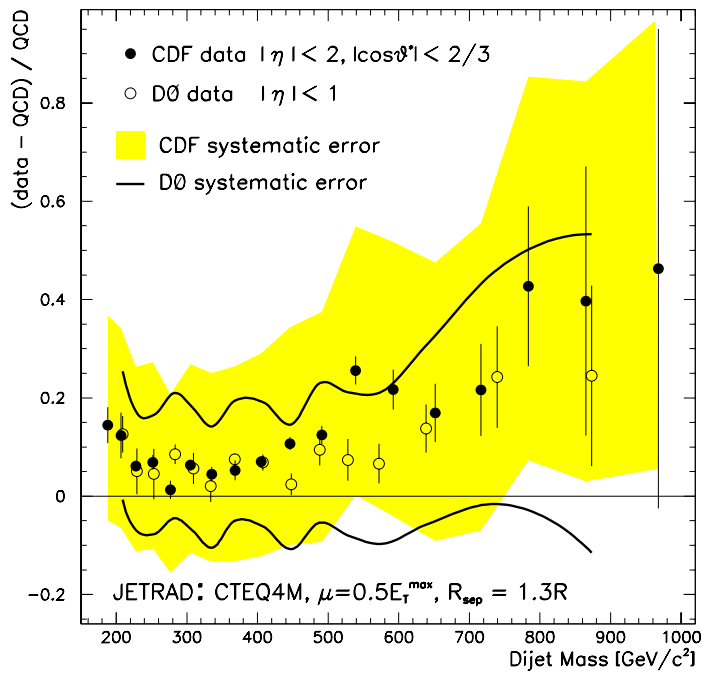
Inclusive Jet Cross Section at All Rapidities



- enhanced gluon at higher x or other effects?
- what about prompt photon data? (k_T , resummation)

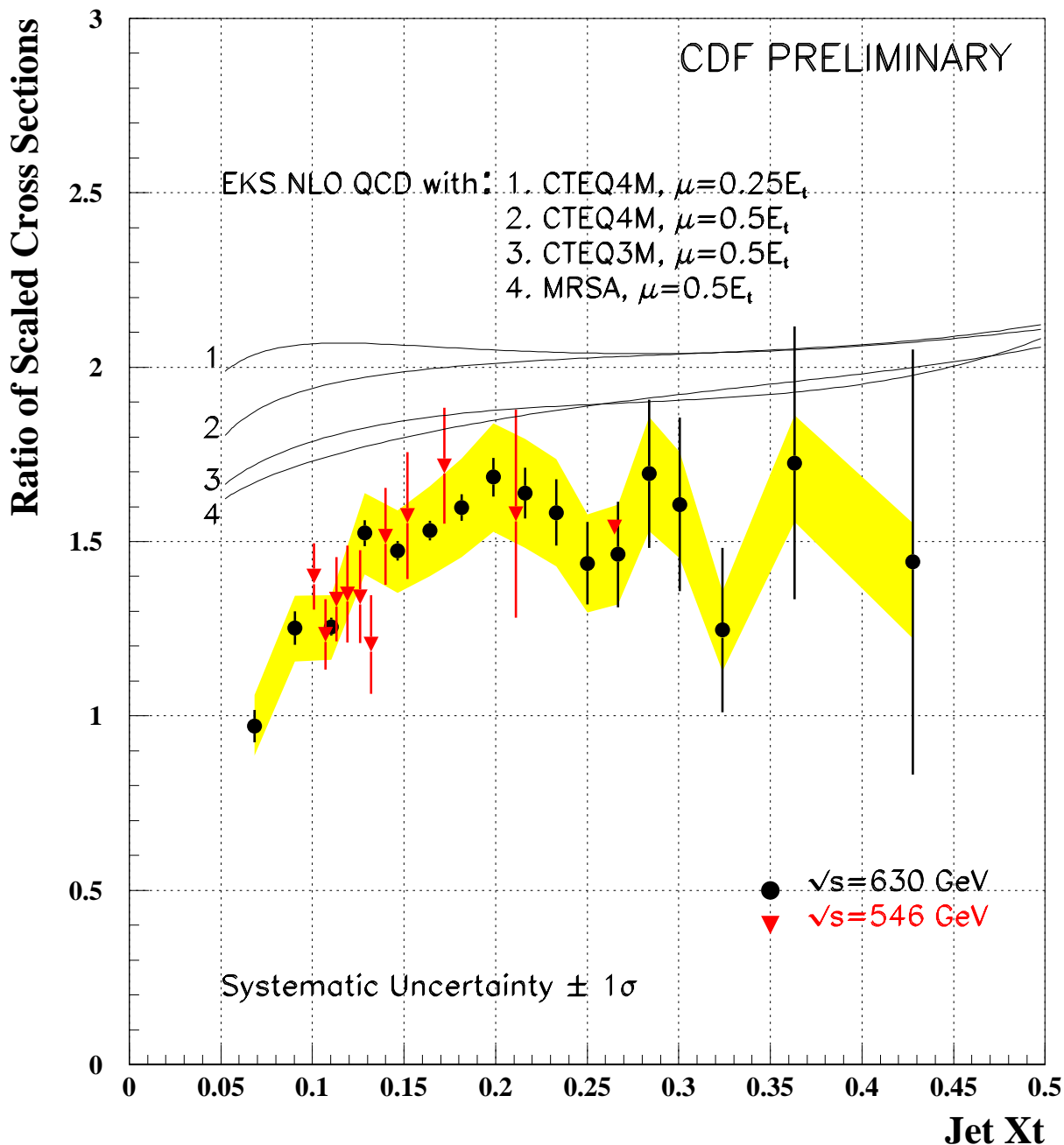
Dijet Mass

$$M \equiv \sqrt{(E_1 + E_2)^2 - (\vec{P}_1 + \vec{P}_2)^2} = \sqrt{2E_{T1}E_{T2}[\cosh(\Delta\eta) - \cos(\Delta\phi)]}$$

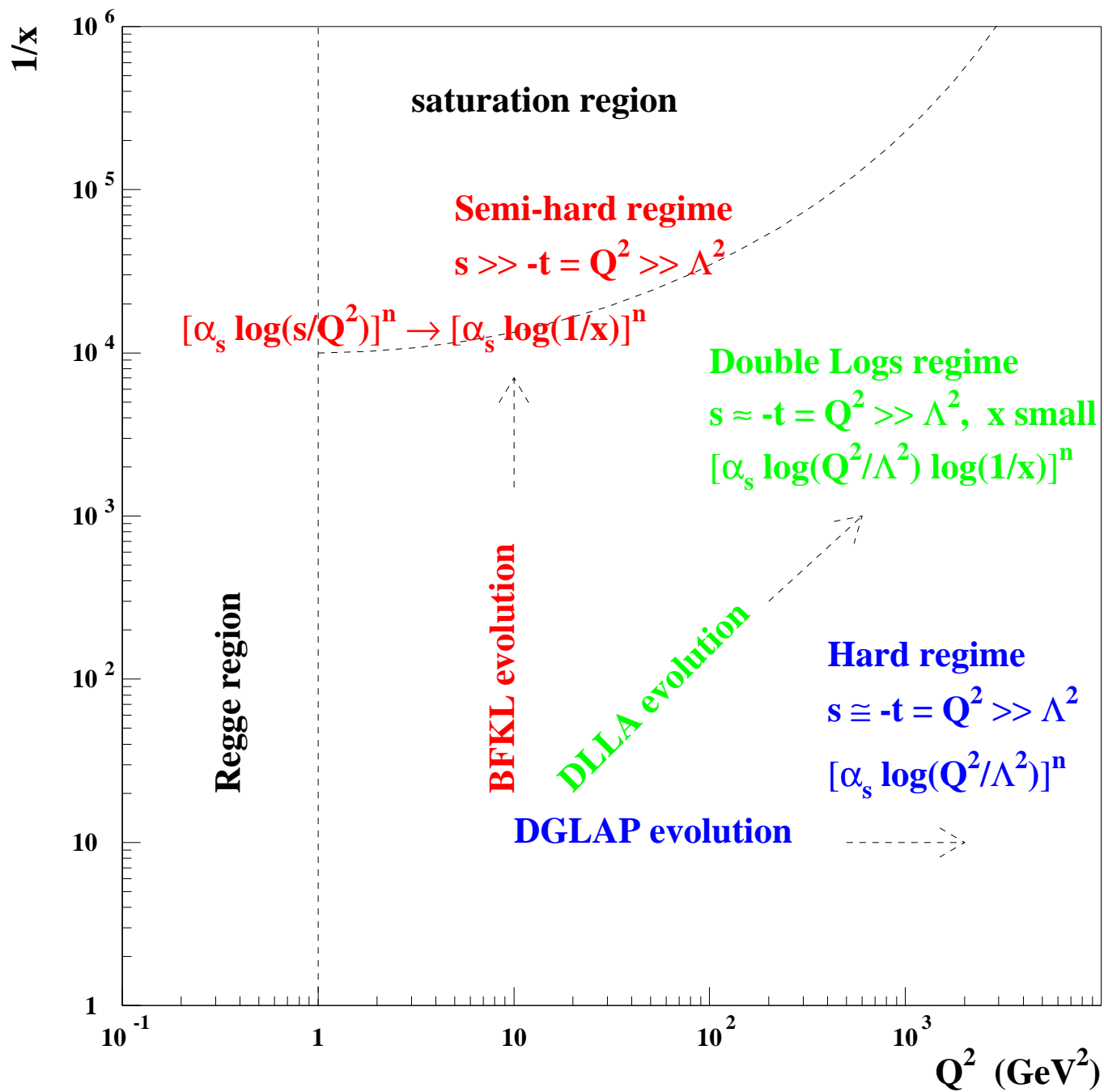


Ratio of Inclusive Jet Cross Sections at $\sqrt{s} = 1800$ and 630 GeV

Scaling violations: at intermediate/high $x_T = 2E_T/\sqrt{s}$,
cross section should decrease with $(Q^2)^{-1} \sqrt{s}$



DGLAP vs BFKL



Dijets at Large Rapidity Intervals

- for $\Lambda_{QCD} \ll Q \ll \sqrt{s} \Rightarrow$ large $\ln(s/Q^2)$ in σ_{dijet}
- $\hat{s} = x_1 x_2 s \Rightarrow$ for large $x_i \Rightarrow$ large $\ln(\hat{s}/Q^2) \approx \Delta\eta$ in $\hat{\sigma}$
- use BFKL to sum all powers of $(\alpha_s \Delta\eta) \Rightarrow$

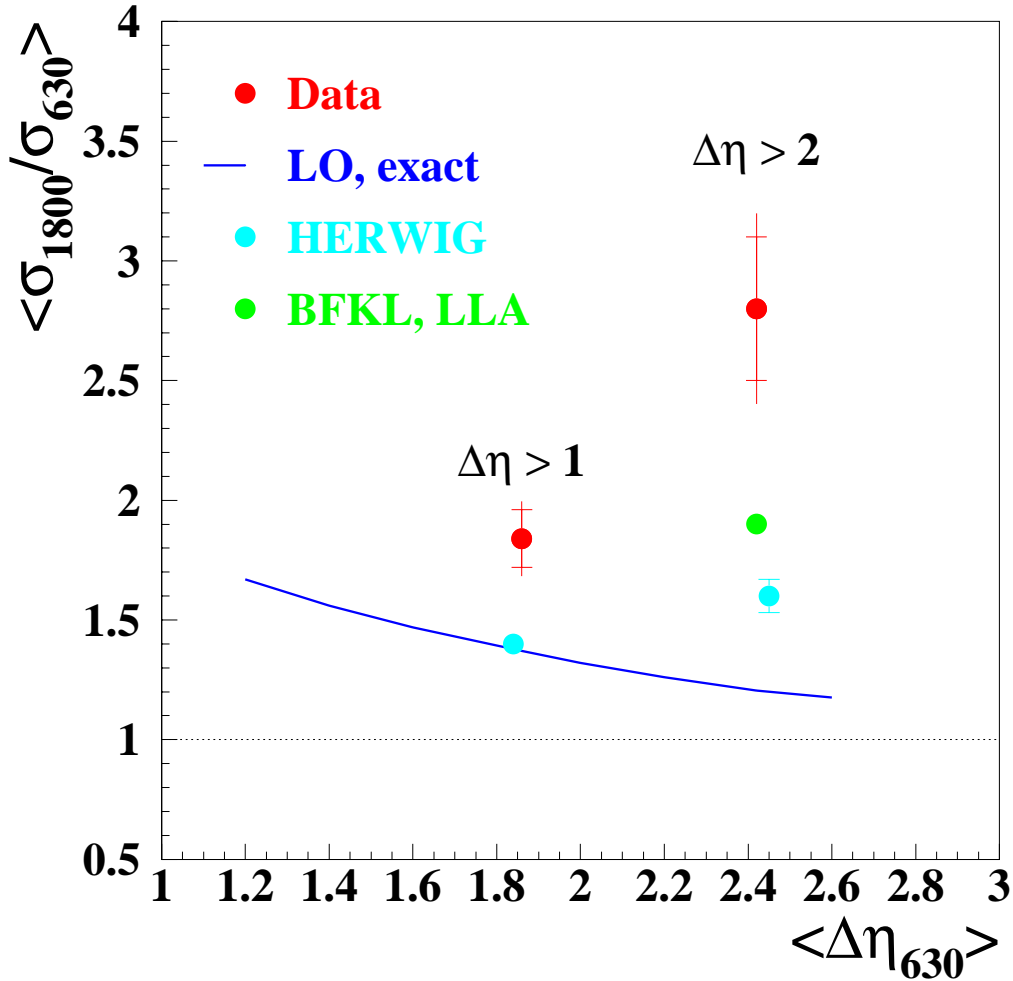
$$\hat{\sigma} \sim \frac{e^{(\alpha_{\text{BFKL}}-1)\Delta\eta}}{\sqrt{\alpha_s \Delta\eta}} \quad \text{with} \quad \alpha_{\text{BFKL}} = 1 + \frac{4 \ln 2 N_c \alpha_s}{\pi} \simeq 1.5$$

- PDF's steeply falling at large $x \Rightarrow$
not possible to see exponential growth in σ_{dijet}
- take ratio at $\sqrt{s_A} = 1800$ GeV, $\sqrt{s_B} = 630$ GeV
- fix $x_1, x_2, Q^2 \Rightarrow$ PDF's cancel \Rightarrow

$$R = \frac{\sigma(x_1, x_2, Q^2, s_A)}{\sigma(x_1, x_2, Q^2, s_B)} = \frac{e^{(\alpha_{\text{BFKL}}-1)(\Delta\eta_A - \Delta\eta_B)}}{\sqrt{\Delta\eta_A / \Delta\eta_B}}$$

- rise with $\Delta\eta \Rightarrow$ rise with s
- extract an effective value for α_{BFKL}

Dijets at Large $\Delta\eta$ (cont'd)



- R increases with $\Delta\eta \Leftrightarrow \sqrt{s}$:

for $\Delta\eta > 2 \Rightarrow \langle \Delta\eta_{630} \rangle = 2.4$ and $\langle \Delta\eta_{1800} \rangle = 4.6$

$$\langle \sigma_{1800}/\sigma_{630} \rangle = 2.8 \pm 0.3 \text{ (stat)} \pm 0.3 \text{ (syst)}$$

- effective BFKL intercept:

$$\alpha_{\text{BFKL}}(20 \text{ GeV}) = 1.65 \pm 0.05 \text{ (stat)} \pm 0.05 \text{ (syst)}$$

- difference in σ between different \sqrt{s} ($\Delta\eta$) increases w/ $\Delta\eta$

Summary

- proton structure function measured
for $0.04 < Q^2 < 3 \cdot 10^4 \text{ GeV}^2$ and $6 \cdot 10^{-7} < x < 0.65$
- DGLAP evolution describes data down to $Q^2 \sim 1 \text{ GeV}^2$
- c content of proton 10-30%
- d/u ratio uncertainty at high x
- high Q^2 CC and NC data agree with SM predictions
- $xg(x, Q^2)$ well constrained at low- x ,
but large uncertainties at high- x
- jet cross section from $p\bar{p}$ at forward $\eta \Rightarrow$ gluon at high- x
- sign of BFKL evolution at dijets with large $\Delta\eta$ in $p\bar{p}$
- await increased L , high Q^2 at HERA and
high E_T jets at Tevatron Run II !