## PHYSICS 715

### Problem Set 12

## Reading: Landau and Lifshitz, Secs. 148-150, 152-153 Huang, Secs. 14.1-14.4, 16.1-16.4 (suggested)

## FINAL EXAM MONDAY, MAY 8, 12:25 pm

#### LD 39: Bose-Einstein condensation in an atomic trap.

Bose-Einstein condensation was observed directly in 1995 in dilute atomic gases confined in magnetic traps (see M. H. Anderson et al., Science **269**, 198 (1995); C. C. Bradley et al., Phys. Rev. Lett. **75**, 1687 (1995); K. B. Davis et al., Phys. Rev. Lett. **75**, 3969 (1995)). The traps are approximately harmonic with characteristic oscillation frequencies  $\nu$  of about 150 Hz. The critical temperature  $T_c$  at which the condensation starts varies with the experimental conditions, but is typically  $\approx 150$  nK.

(a) A mean field model for an ideal Bose gas in an external potential  $V(\mathbf{r})$  describes the gas in a volume  $d^3r$  at a point  $\mathbf{r}$  as being in local equilibrium and having a local chemical potential  $\mu(\mathbf{r}) = \mu - V(\mathbf{r})$ , where  $\mu$  is the constant chemical potential for the entire system. When is this a reasonable approximation? The number distribution of the particles in the volume  $d^3r$  is given in this picture by a standard Bose number distribution with  $\mu$  replaced by  $\mu(\mathbf{r})$ , and the usual volume factor replaced by  $d^3r$ :

$$dN = \frac{1}{e^{\beta(E(\mathbf{p}) - \mu(\mathbf{r})} - 1} \frac{d^3 p \, d^3 r}{h^3} = \frac{1}{e^{\beta(H(\mathbf{p}, \mathbf{r})) - \mu} - 1} \frac{d^3 p \, d^3 r}{h^3}.$$

Integrate over all positions and momenta to obtain an expression for the maximum number of particles N that can be accomodated in a harmonic trap with  $V = \frac{1}{2}m\omega^2 \mathbf{r}^2$ , and use the result to determine the critical temperature for Bose-Einstein condensation in terms of N and the parameters in H. Evaluate  $T_c$  for  $N = 4 \times 10^4$ <sup>87</sup>Rb atoms in a 150 Hz trap (5S<sub>1/2</sub> electron and a nucleus with spin s = 3/2 in a total angular momentum state with F = 2,  $m_F = 2$ ). Compare the result to that in the first reference above. [Hint: the integral can be reduced to a standard form by introducing scaled variables  $r' = r/\lambda$ ,  $p' = \lambda p$ ,  $\lambda = 1/\sqrt{m\omega}$ , introducing the the 6-dimensional vector  $\mathbf{x} = (\mathbf{r}', \mathbf{p}')$ , writing the volume element as  $d^6x$ , and changing to  $x^2$  as the integration variable after performing the angular integration.]

Determine the number  $N_0$  of the atoms in the ground state of the system as a function of  $T/T_c$ , and plot of the ratio  $N_0/N$  versus  $T/T_c$ .

(b) The ground state wave function for an isotropic oscillator in three dimensions is

$$\psi_0(r) = \frac{1}{\pi^{3/4} r_0^{3/2}} e^{-r^2/2r_0^2}, \qquad r_0 = \sqrt{\frac{\hbar}{m\omega}}.$$

Determine the scaled number distribution  $\frac{1}{N} \frac{dN}{d^3(r/r_0)}$  of atoms in the trap and plot it as a function of  $r/r_0$  for  $N_0/N = 0$ , 0.1, 0.3, 1.0 and  $0 \le r/r_0 \le 10$ . Evaluate  $d^3N/dr^3$  at r = 0. Is this a high density? Explain. Compare your results qualitatively with those in the references above. [Hint: the momentum integral which appears must be evaluated numerically. Express the integrand in terms of  $\hbar \omega/kT_c$ ,  $T/T_c$ , and  $r/r_0$  to see its structure before doing the integration. Note that for a given  $T_c$ ,  $N = N_c$ .]

# LD 40: The Bragg-Williams approximation for the 3-dimensional Ising model as a Landau-Ginzburg mean field theory.

Mean field theory (or the Bragg-Williams approximation) gives the expression

$$m = \mu_0 \tanh\left[\frac{\mu_0 \mathcal{H}}{kT} + \frac{T_c}{T}\frac{m}{\mu_0}\right]$$

for the magnetic moment per spin in an Ising model with particles with magnetic moment  $\mu_0$ .  $\mathcal{H}$  is the applied magnetic field. The total magnetization per unit volume is M = nm, where n is the density of spins.

- (a) Show that the magnetic Gibbs function has the form assumed in the Landau-Ginzburg mean field approach when calculated to order  $m^4$  with  $\mathcal{H}$  retained as a free variable. [Hint: rewrite the relation above as an equation for  $\mathcal{H}$ , expand the inverse hyperbolic tangent which appears to order  $m^3$ , and determine  $F_M$  by integrating the relation  $\mathcal{H} = \frac{\partial F}{\partial M}$ .  $G_M = F M\mathcal{H}$ . Do not eliminate  $\mathcal{H}$  in  $G_M$ .]
- (b) Show that  $T_c$  is in fact the critical temperature at which spontaneous magnetization appears at  $\mathcal{H} = 0$  in the mean field theory. Determine the temperature dependence of the magnetization and the magnetic susceptibility  $\chi = \partial M/\partial \mathcal{H}$  for  $\mathcal{H} \to 0$  to lowest order in  $|T_c - T|$  for  $T > T_c$  and for  $T < T_c$ , and find the critical exponents  $\beta$  and  $\gamma$ . [Hint: minimize  $G_M$  at fixed  $\mathcal{H}$  to determine M. Recall that, in the Landau approach, one is dealing with a Taylor series expansion in T as well as M, and replace T by  $T_c$  wherever it does not appear in the difference  $|T - T_c|$ . Show that your solution for M actually minimizes  $G_M$  in the two temperature regions.]
- (c) Plot separate figures giving the approximate expressions for  $M/n\mu_0$  and  $\chi kT_c/n\mu_0^2$ for  $0 \leq T/T_c \leq 2$  and values in the the ranges  $0 \leq M/n\mu_0 \leq 2$  and  $0 \leq \chi kT_c/n\mu_0^2 \leq 4$ . Include on the same figures the exact results for these quantities obtained by solving the equation above and the corresponding equation for  $dM/d\mathcal{H}$  (evaluated for  $\mathcal{H} = 0$ ). [Hint: To find the nontrivial solution for Mnumerically for  $T < T_c$ , find the zero of  $1 - \frac{1}{x} \tanh \frac{T_c}{T} x$ ,  $x = m/\mu_0$ .]

#### LD 41: First order phase transitions in the Landau theory.

Consider a mean field theory with order parameter m and Landau free energy

$$G = \frac{1}{2}am^2 - \frac{1}{3}bm^3 + \frac{1}{4}cm^4, \qquad a, b, c > 0.$$

Sketch the possible shapes for G(m) as b is increased from zero with a and c fixed, and show that a first order phase transition occurs for b equal to a critical value  $b_c$ . Determine  $b_c$  and the change in the equilibrium value of m at the transition. [Hint: The scaling  $m = (2a/c)^{1/2}m'$ ,  $G = (a^2/c)G'$  is useful for seeing the structure of G. Make your sketches for m' small. To find  $b_c$ , determine when G has two quadratic minima with G = 0, but at different values of m.]

## LD 42: Critical exponents from a Landau free energy.

The Landau free energy for a hypothetical system is of the form

$$F(t,m,h) = -hm + \frac{1}{4}atm^4 + \frac{1}{6}bm^6, \qquad a, b > 0,$$

where *m* is the order parameter,  $t = (T - T_c)/T_c$ , and the external field *h* is the variable conjugate to *m*. Show that there is a "second order" phase transition at  $T = T_c$  for h = 0. Determine the critical exponents  $\alpha, \beta, \gamma, \delta$ , and show directly that

$$\alpha + 2\beta + \gamma = 2, \qquad \delta = 1 + \frac{\gamma}{\beta}.$$

Show explicitly that your solution for m actually minimizes F for t > 0 and t < 0. Keep only the leading behavior in t for  $t \to 0$ . [Hint: Calculate the thermodynamic derivatives which determine S and  $\chi = \left(\frac{\partial h}{\partial m}\right)_t^{-1}$ , and evaluate them for the value of m that minimizes F.  $C = T \frac{\partial S}{\partial T}$ . To determine  $\delta$ , minimize F on the critical isotherm t = 0 for  $h \neq 0$ .]