PHYSICS 715

Problem Set 6

Due Friday, March 10, 2006

Reading: Landau and Lifshitz, Secs. 30, 33, 78, 79

LD 18: Thermal properties of rubber – a model.

A simple one-dimensional model of rubber treats the rubber molecule (a long-chain polymer) as a limp chain with n links of length a which can be oriented either to the right (n_+) or to the left (n_-) . The separation ℓ between the ends of the folded molecule is $\ell = (n_+ - n_-)a$. Determine the number of ways $\Gamma(n_+, n_-)$ of obtaining a configuration with a given ℓ and show that, when only the average value of ℓ is specified, the appropriate grand partition function for the molecule is

$$\mathcal{Y}(\tau, T) = \sum_{\substack{n_+, n_- \\ n_+ + n_- = n}} \Gamma(n_+, n_-) e^{\beta \tau (n_+ - n_-)a} \, .$$

 τ may be identified as the tension in the chain through the first law, $dG = -SdT - \ell d\tau + \mu dn$ with Gibbs free energy $G = -kT \ln \mathcal{Y}$.

Evaluate the sum above and use the result to determine $\bar{\ell} = \langle \ell \rangle$ as a function of T and τ . Determine the coefficient of thermal expansion $(1/\ell)(\partial \bar{\ell}/\partial T)_{\tau}$ and comment on the result.

LD 19: Electric dipole interactions of polar molecules

The dipole interaction between polar diatomic molecules with electric dipole moments \mathbf{p} is given approximately by

$$V_{12}(\mathbf{r}) = \begin{cases} -\frac{1}{r^3} \left[3(\mathbf{p_1} \cdot \hat{\mathbf{r}})(\mathbf{p_2} \cdot \hat{\mathbf{r}}) - (\mathbf{p_1} \cdot \mathbf{p_2}) \right], & r > r_0, \\ +\infty, & r < r_0 \end{cases}$$

Here $\mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2$ is the relative coordinate giving the separation of the centers of mass of molecules 1 and 2 and $\hat{\mathbf{r}} = \mathbf{r}/r$. The moments \mathbf{p}_i have the same magnitude but may have different orientations. The canonical partition function for the *N*-particle classical gas confined in a volume *V* can be reduced to the form

$$Z_N = \frac{Z_{\text{rot}}^N}{N! \lambda^{3N}} \frac{1}{(4\pi)^N} \int d^3 r_1 \cdots d^3 r_N \int d\Omega_1 \cdots d\Omega_N e^{-\beta \sum_{i < j} V_{ij}(\mathbf{r}_i - \mathbf{r}_j)}$$

where $d\Omega_i = d(\cos \theta)_i d\phi_i$ with θ_i and ϕ_i the polar and azimuthal angles that specify the orientation of \mathbf{p}_i , here averaged over.

Calculate the second virial coefficient for this gas correct to order $(V_{12}/kT)^2$.

LD 20: The Coulomb energy and equation of state of a neutral plasma. Definition and use of the radial distribution function.

The equation of state of a neutral plasma with an average of $N_{\rm e}/V = N_{\rm ion}/V = \frac{1}{2}n$ electrons and singly charged positive ions per unit volume is given by the virial theorem as

$$P = nkT + \frac{1}{3} \frac{1}{V} \sum_{i < j} \left\langle \frac{e_i e_j}{|\mathbf{x}_i - \mathbf{x}_j|} \right\rangle$$
$$= nkT + \frac{1}{6} nN \left(\left\langle \frac{e^2}{|\mathbf{x}_1 - \mathbf{x}_2|} \right\rangle_{ee} - \left\langle \frac{e^2}{|\mathbf{x}_1 - \mathbf{x}_{N+1}|} \right\rangle_{e-ion} \right).$$

The average is over the exact canonical distribution. We have counted the numbers of pairs of each type, hence the factor nN, and have used the identity of the electronelectron and ion-ion distributions to combine the e-e and ion-ion terms. Electrons are labeled $1, \ldots, N$, and ions, $N + 1, \ldots, 2N$. (See Prob. LD 14.)

(a) Write out explicitly what $\langle \cdots \rangle$ means in terms of integrals involving the potentials, and give an argument to show that P can be written exactly as

$$P = nkT + \frac{1}{6}n\left(\int d^{3}r \,\frac{e^{2}}{r}[n_{\rm ee}(r) - n_{\rm e,ion}(r)]\right),\,$$

where $n_{ee}(r)$ and $n_{e,ion}(r)$ are the radial distributions of electrons and ions relative to a given electron, that is, the numbers of particles per unit volume a distance rfrom the chosen particle. Give the definition of n_{ee} .

(b) Evaluate the integrals using the Debye-Hückel expressions for the distribution functions, and determine the temperature dependence of the correction to the ideal gas expression for P.

LD 21: Properties of the quantum density matrix

A quantum system has the probabilities p_i of being in the normalized pure states $|\Psi_i\rangle$, i = 1, 2, ..., with $\sum_i p_i = 1$. The state $|\Psi_i\rangle$ are not necessarily orthogonal, but the Ψ 's can be expanded using a complete orthonormal set of basis states $|\phi_m\rangle$, $|\Psi_i\rangle = \sum_m |\phi_m\rangle c_m^{\ i}$ with $c_m^{\ i} = \langle \phi_m | \Psi_i \rangle$. The density matrix for this system is $\rho_{mn} = \sum_i p_i c_m^{\ i} c_n^{\ i*}$ Show that the following relations hold.

- (i) $\text{Tr} \rho = 1.$
- (ii) $\operatorname{Tr} \rho^2 \leq 1$ with $\operatorname{Tr} \rho^2 = 1$ if and only if ρ describes a pure state.
- (iii) $\hat{\rho}$ is a positive Hermitian operator, $\langle \chi | \hat{\rho} | \chi \rangle \geq 0$ for any $| \chi \rangle$, and $\hat{\rho}^{\dagger} = \hat{\rho}$ where $\hat{\rho}$ is the operator with matrix elements ρ_{mn} .
- (iv) $\langle \hat{f} \rangle \equiv \sum_{i} p_i \langle \Psi_i | \hat{f} | \Psi_i \rangle = \text{Tr} \rho f$ for any observable, where $f_{nm} = \langle \phi_n | \hat{f} | \phi_m \rangle$.
- (v) the trace formula in (d) is unchanged by a change in basis $\{\phi_n\} \to \{\chi_m\}$, that is, $\operatorname{Tr} \rho' f' = \operatorname{Tr} \rho f$ where ρ' , f' are calculated in the new basis.