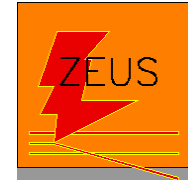




Dijets at HERA: A QCD Story



Douglas Chapin
ZEUS Collaboration
University of Wisconsin–Madison



OUTLINE

HERA and ZEUS

Deep Inelastic Scattering and pQCD

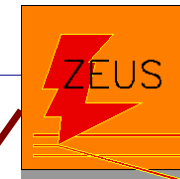
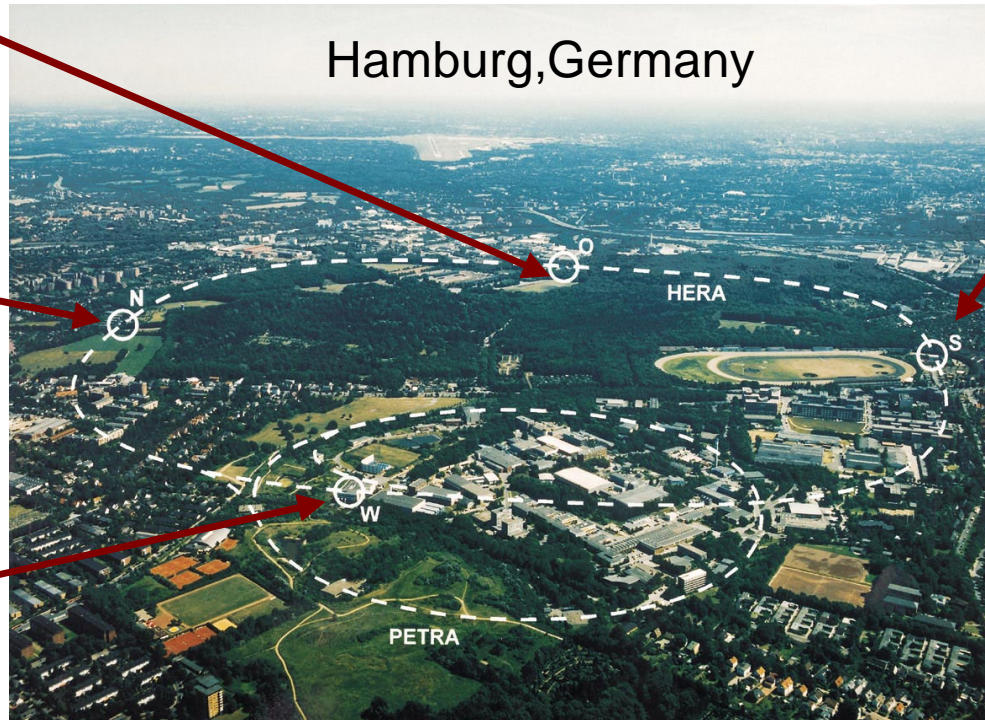
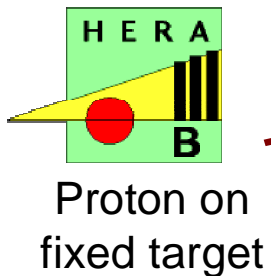
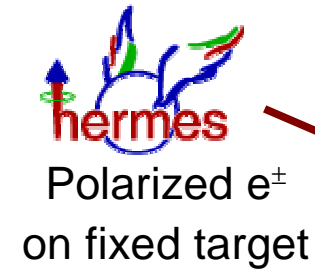
Structure functions and the gluon

Dijets and pQCD

Dijet cross section measurement

The Future

The HERA Collider at DESY



Collider
Detector

6.4km
circumference

920 (820) GeV
protons

27.5 GeV
electrons
or positrons

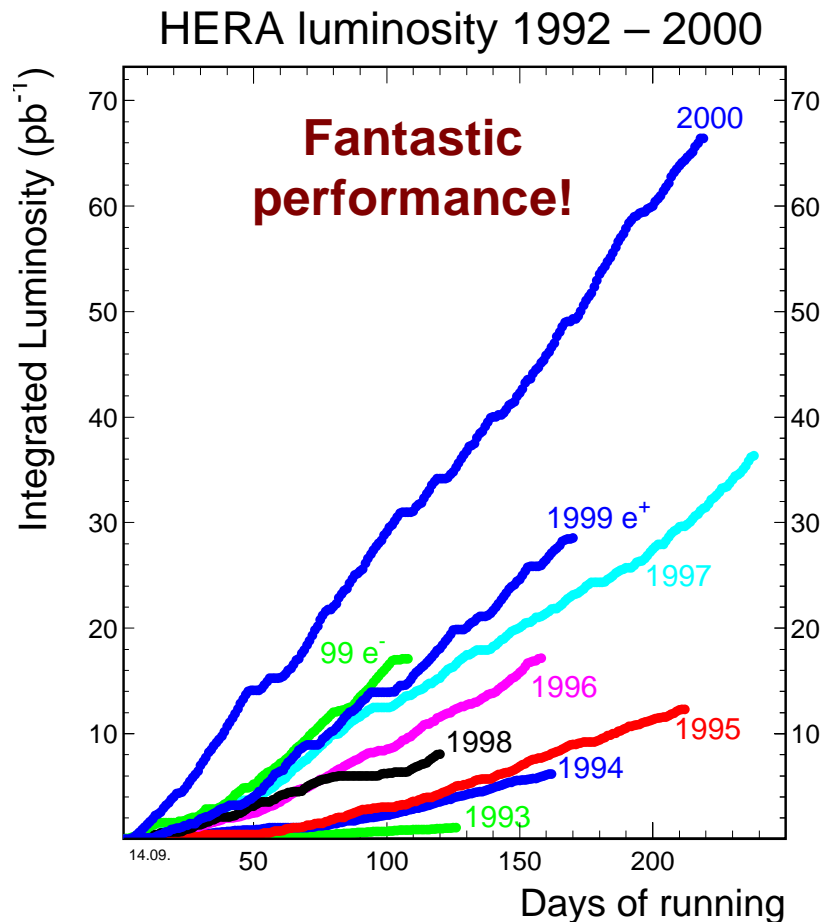
220 bunches
96ns crossing
interval

$$\sqrt{s} \approx 320 \text{ GeV}$$

equivalent to a 50TeV fixed target beam

$$\text{Instantaneous luminosity} \approx 1.8 \times 10^{31} \text{ cm}^{-2} \text{ s}^{-1}$$

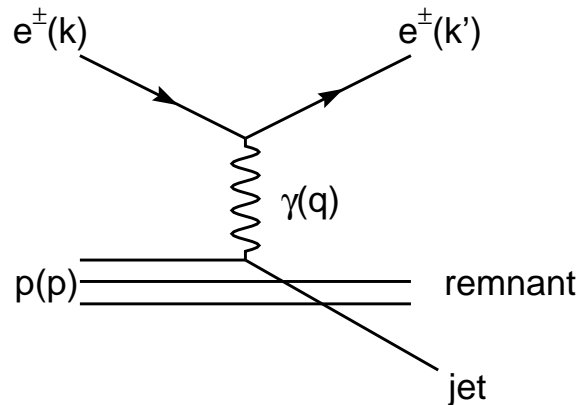
HERA Delivered Luminosity



- e^- in 92–94,98,99: 27pb^{-1}
- e^+ in 95–97,99–00: 166pb^{-1}
- 820GeV protons through 1997
920 GeV since 1998
- ZEUS integrated lumi since 1992:
 $\sim 130\text{pb}^{-1}$ (70% of delivered)
- **Currently undergoing a luminosity upgrade**
 - ready Summer 2001
 - **expect 1fb^{-1}** by end of 2005
 - 5 times current integrated total

Deep Inelastic Scattering

electron-proton scattering



Can also exchange Z, W^\pm

Any lepton – hadron pair

$e-p$ (HERA)

$e-A$ (SLAC)

$\nu-Fe$ (CCFR)

$\mu-A$ (E665, NMC, BCDMS)

DIS kinematic variables

$$Q^2 = -q^2 = -(k - k')^2 \quad \text{Momentum transfer}$$

$$x = \frac{Q^2}{2p \cdot q}$$

Fraction of the proton's momentum that participates in the hard scatter

$$y = \frac{p \cdot q}{p \cdot k}$$

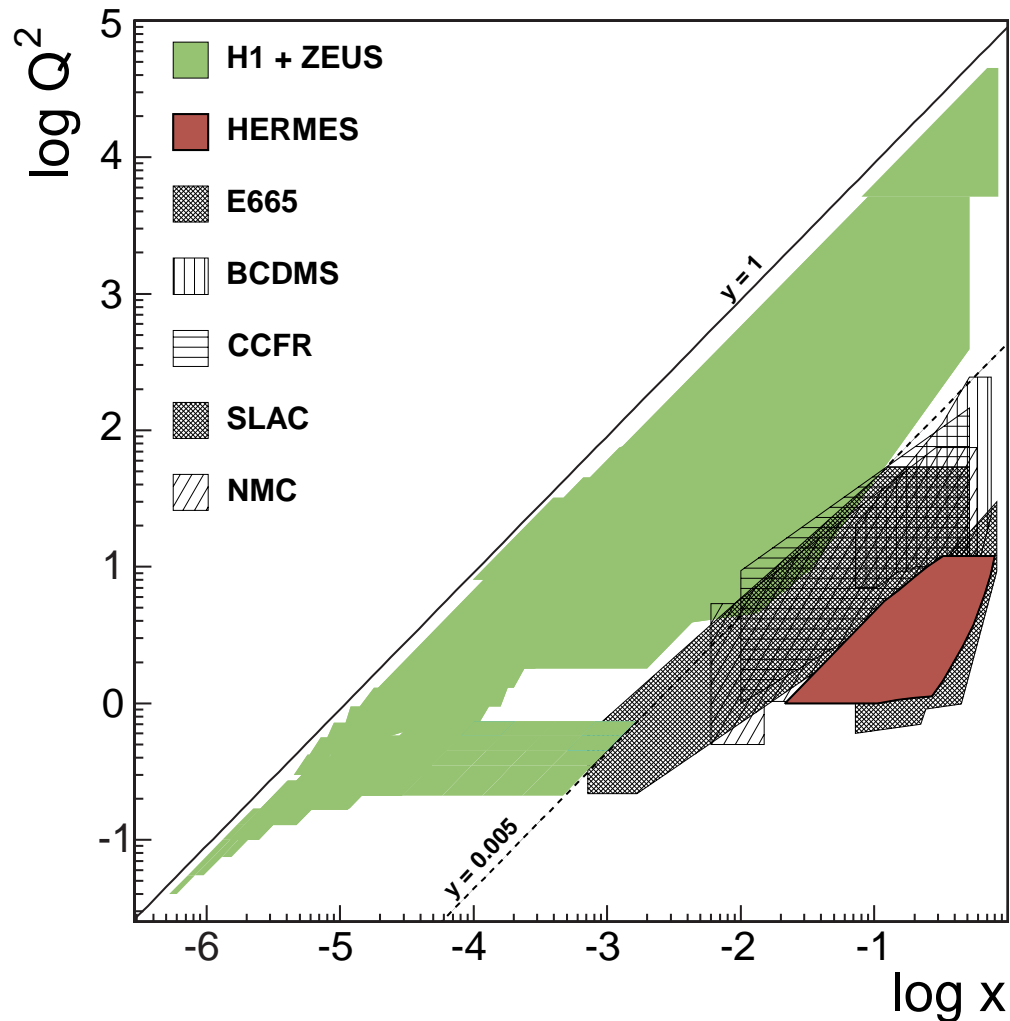
Fraction of the electron's energy available in the proton's rest frame

$$Q^2 = sxy$$

s = center of mass energy squared

HERA Kinematic Range

$$Q^2 = sxy$$



Extended kinematic region
available at HERA

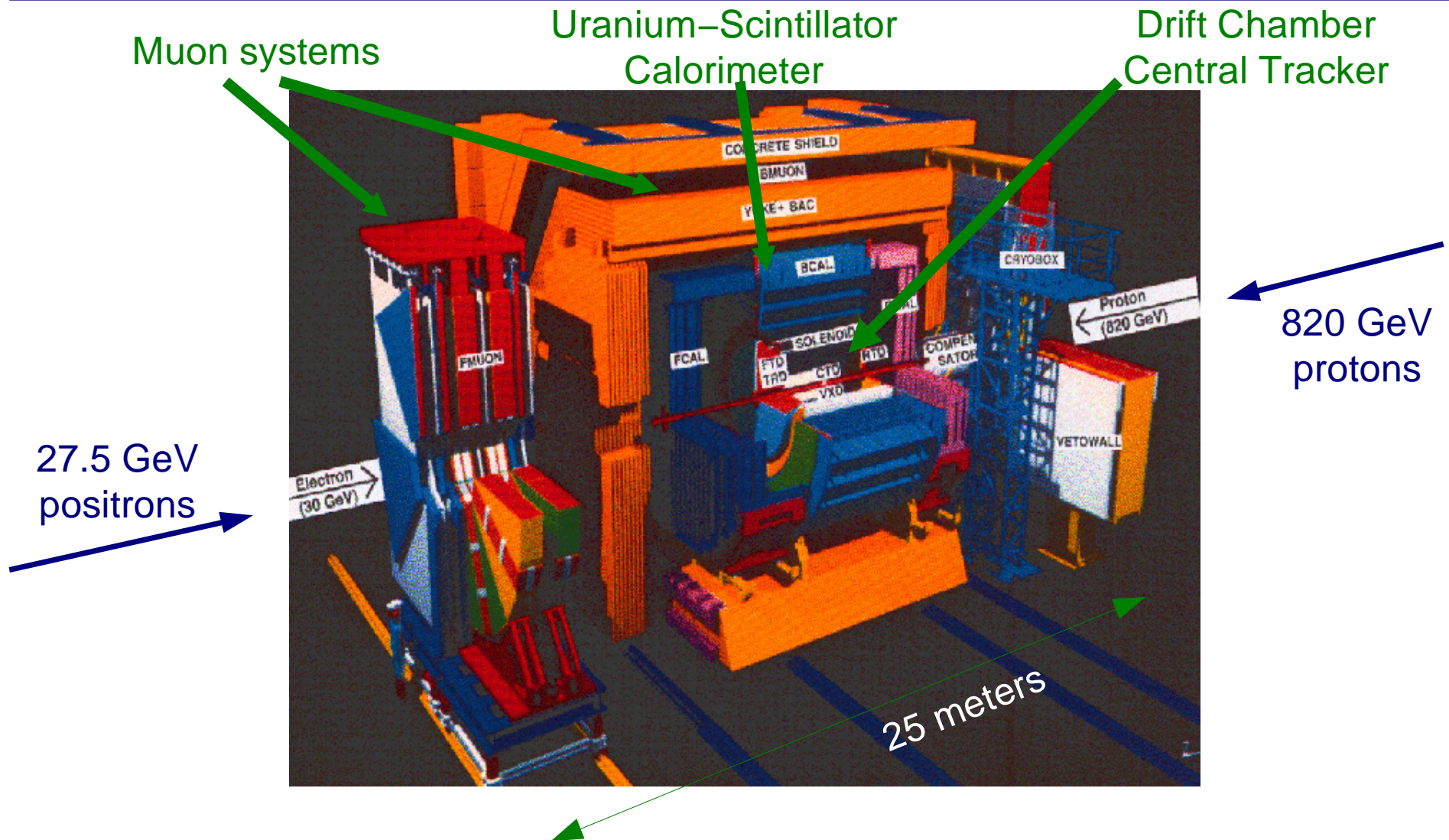
Additional ZEUS components
provide overlap with fixed target
experiments

$$0.45 \text{ GeV}^2 < Q^2 < 20000 \text{ GeV}^2$$

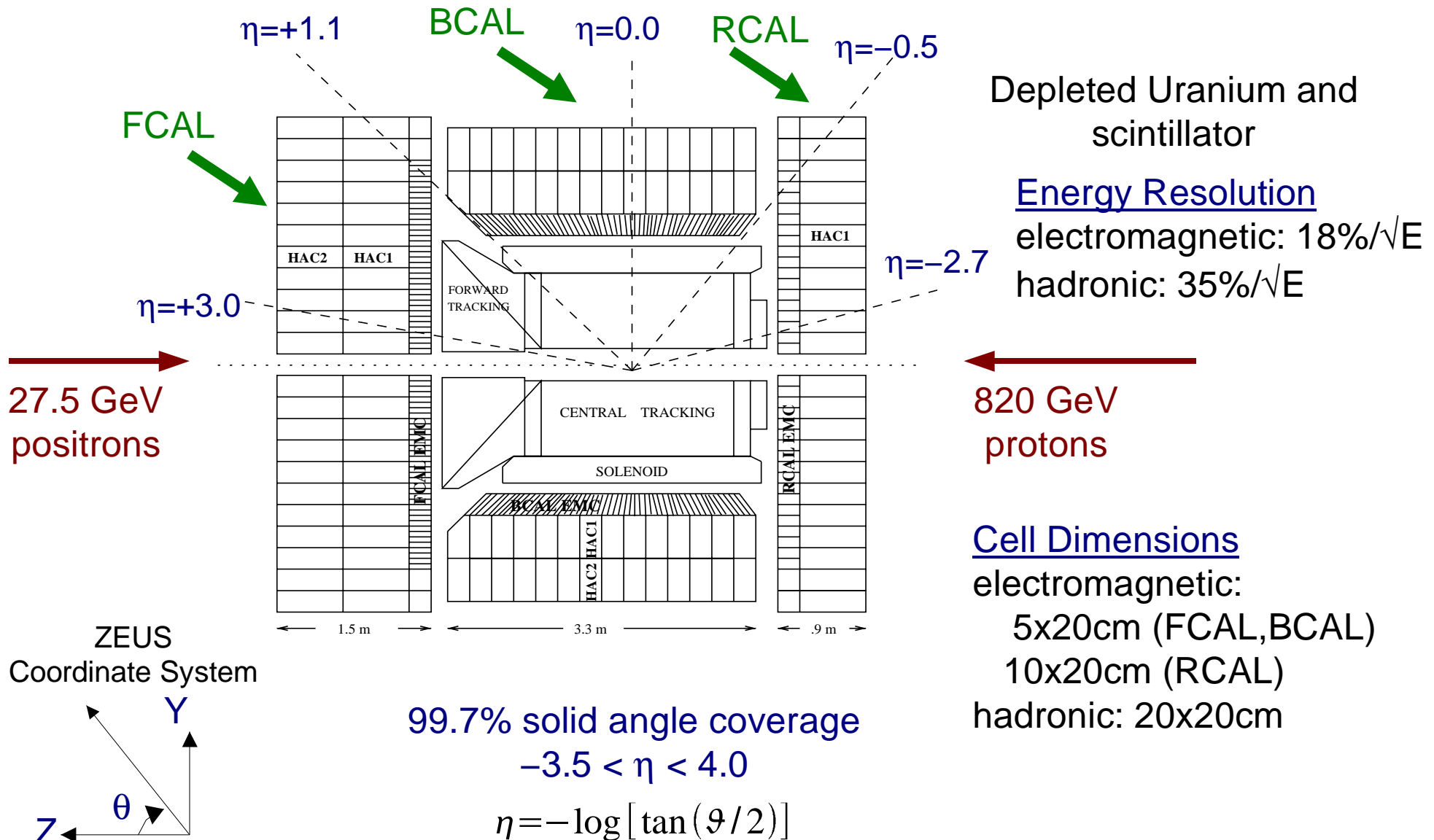
$$10^{-6} < x < 0.9$$

H1 and ZEUS: DESY e-p
 HERMES: DESY e-A
 E665: Fermilab μ -A
 BCDMS: CERN μ -A
 CCFR: Fermilab ν -A
 SLAC: many experiments e-A
 NMC: CERN μ -A

The ZEUS Detector at HERA

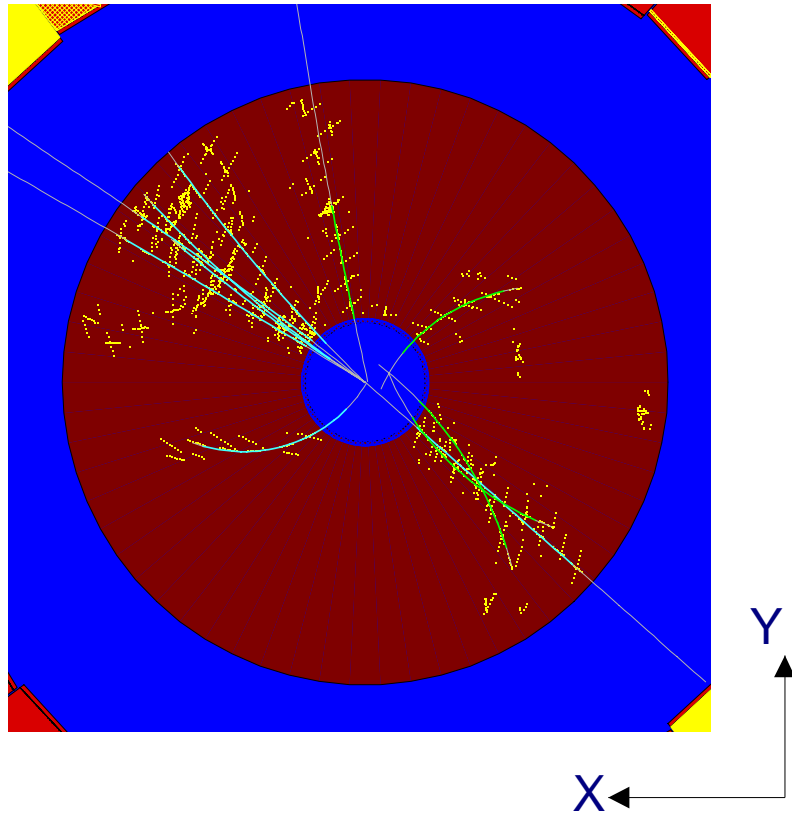


ZEUS Calorimeter



ZEUS Central Tracking Detector

View along beamline



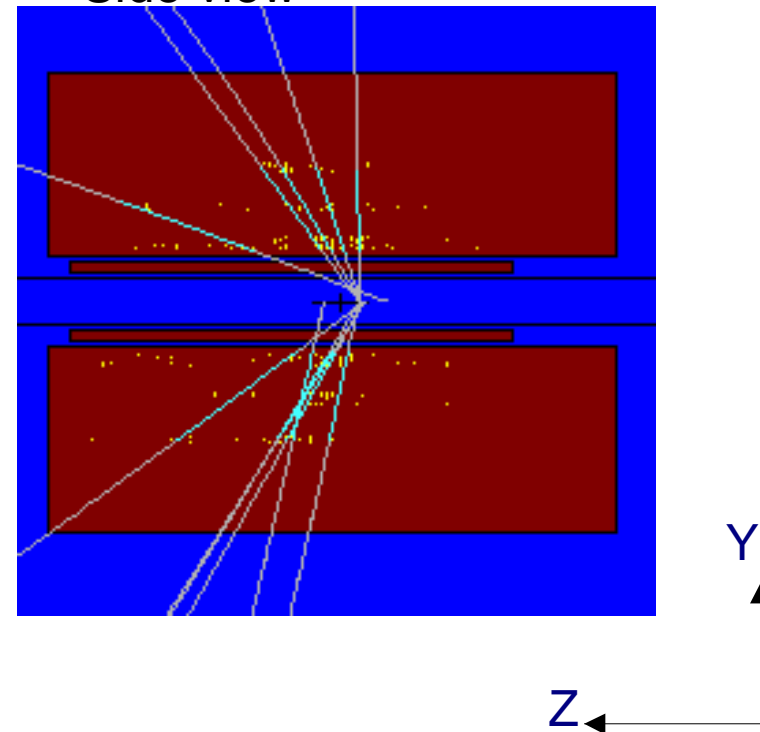
Vertex Resolution

longitudinal (z): 4mm

transverse (x-y): 1mm

Drift Chamber
inside 1.43T solenoid

Side view

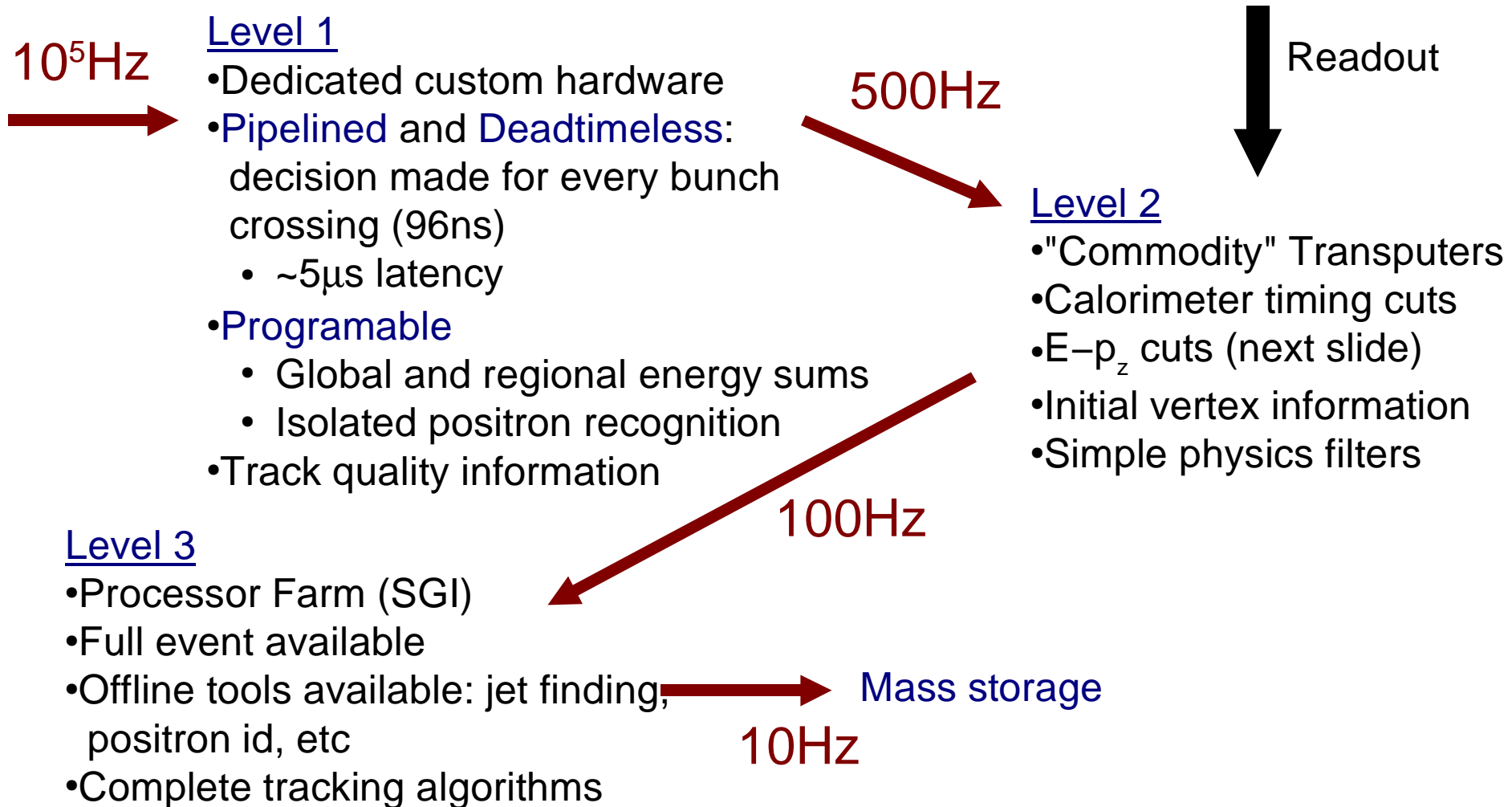


Micro Vertex Detector to be
installed next year

ZEUS Trigger

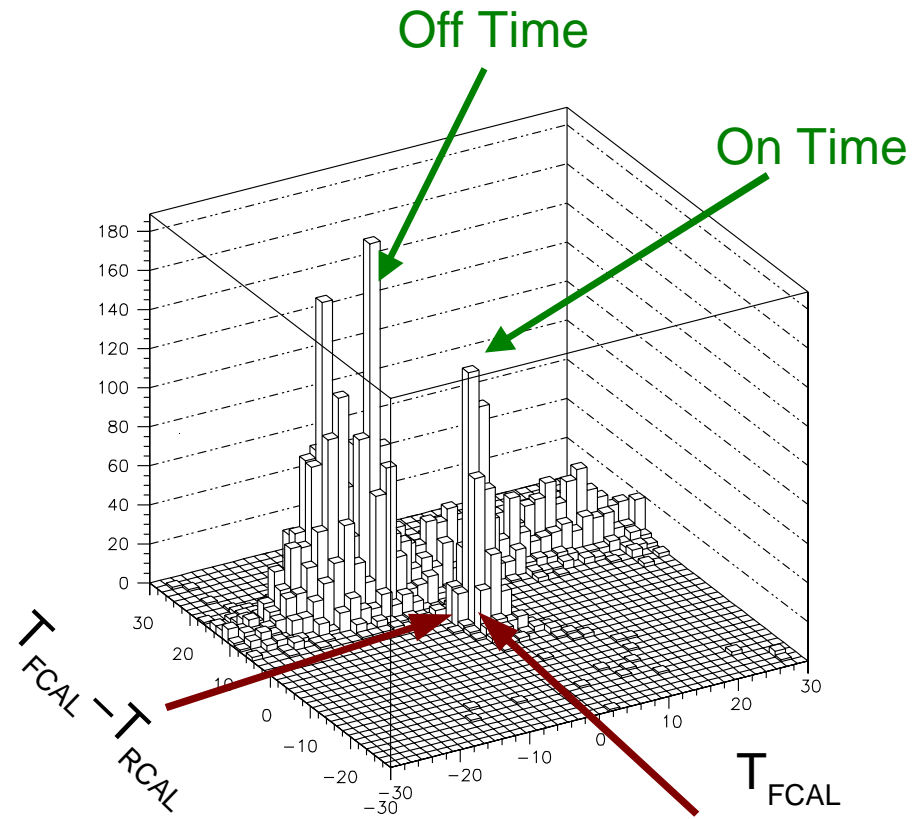
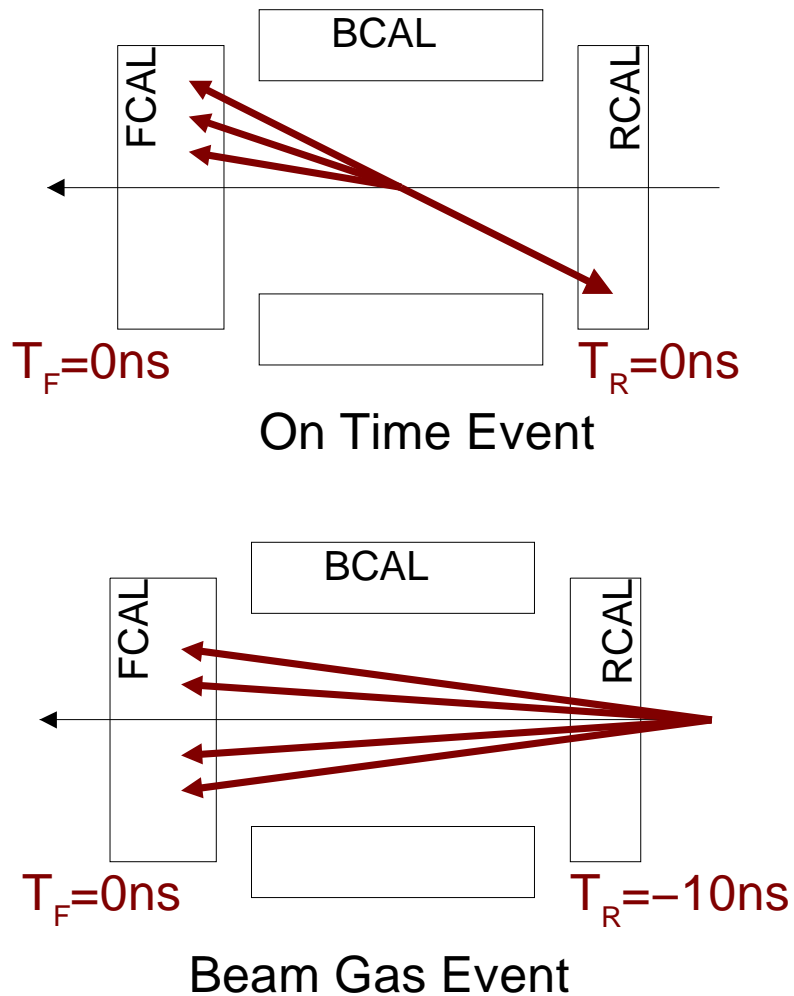
Challenge

10 MHz bunch crossing rate
Extract 10Hz Physics from 100kHz background



Beam Gas Background Rejection

"Distance" between FCAL and RCAL is $\sim 10\text{ns}$



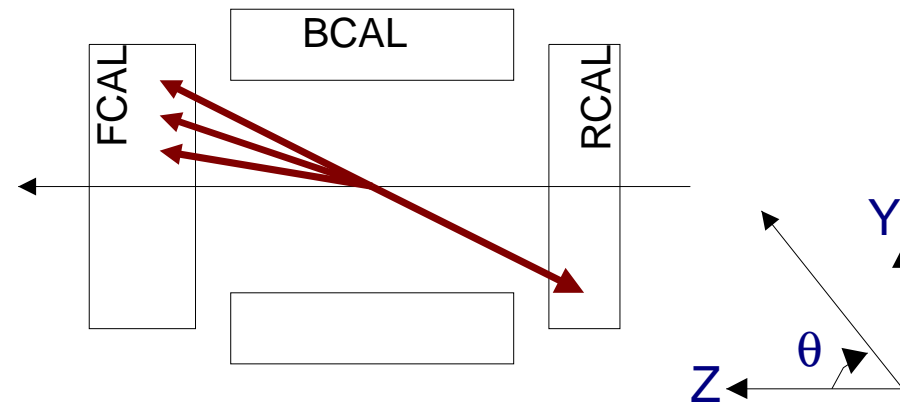
Calorimeter timing at Level 2

ZEUS Calorimeter timing
resolution $< 1\text{ns}$

Background Reduction: $E-p_z$

$$E - p_z = \sum_i E_i (1 - \cos \vartheta_i)$$

Sum runs over calorimeter cells

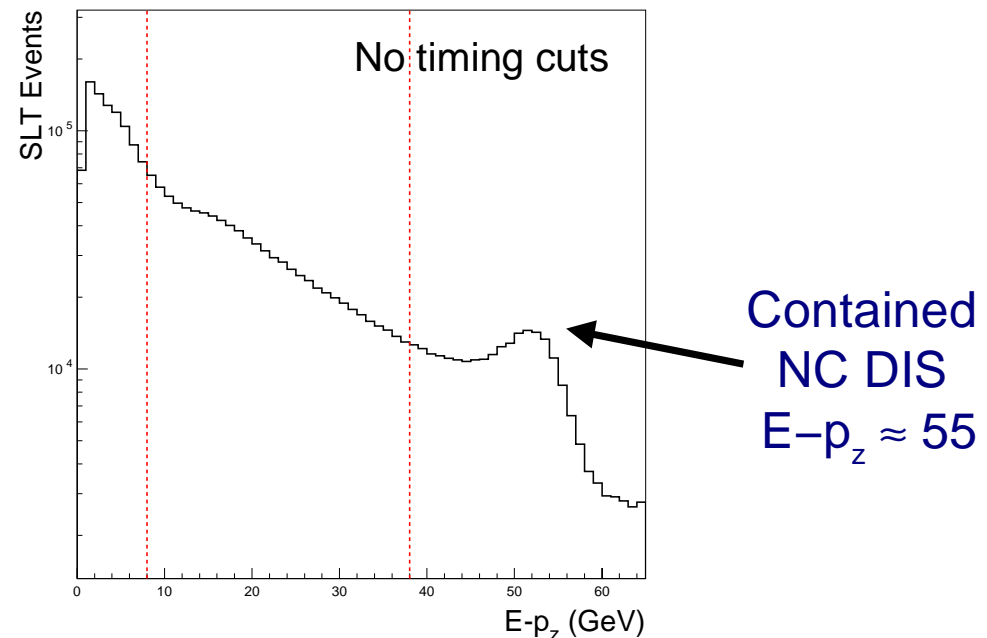


In a given frame, $E-p_z$ is conserved

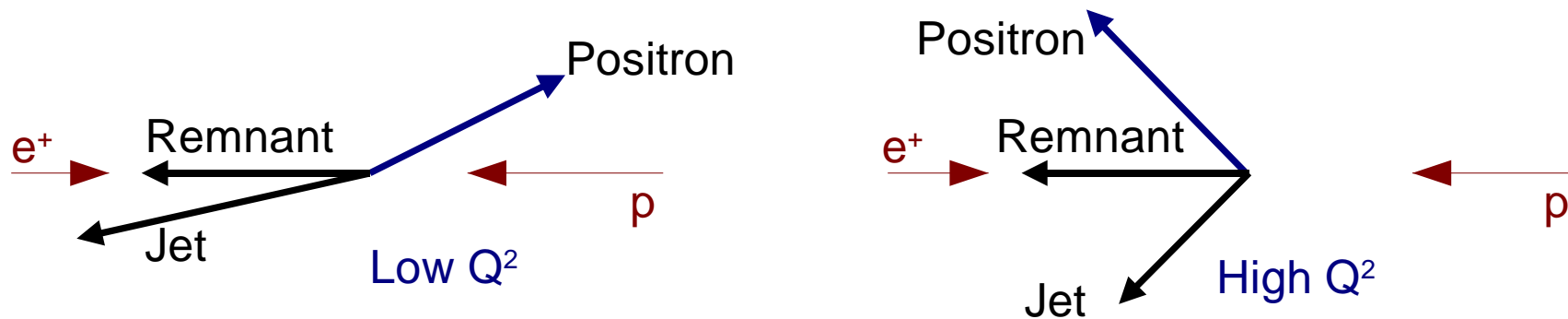
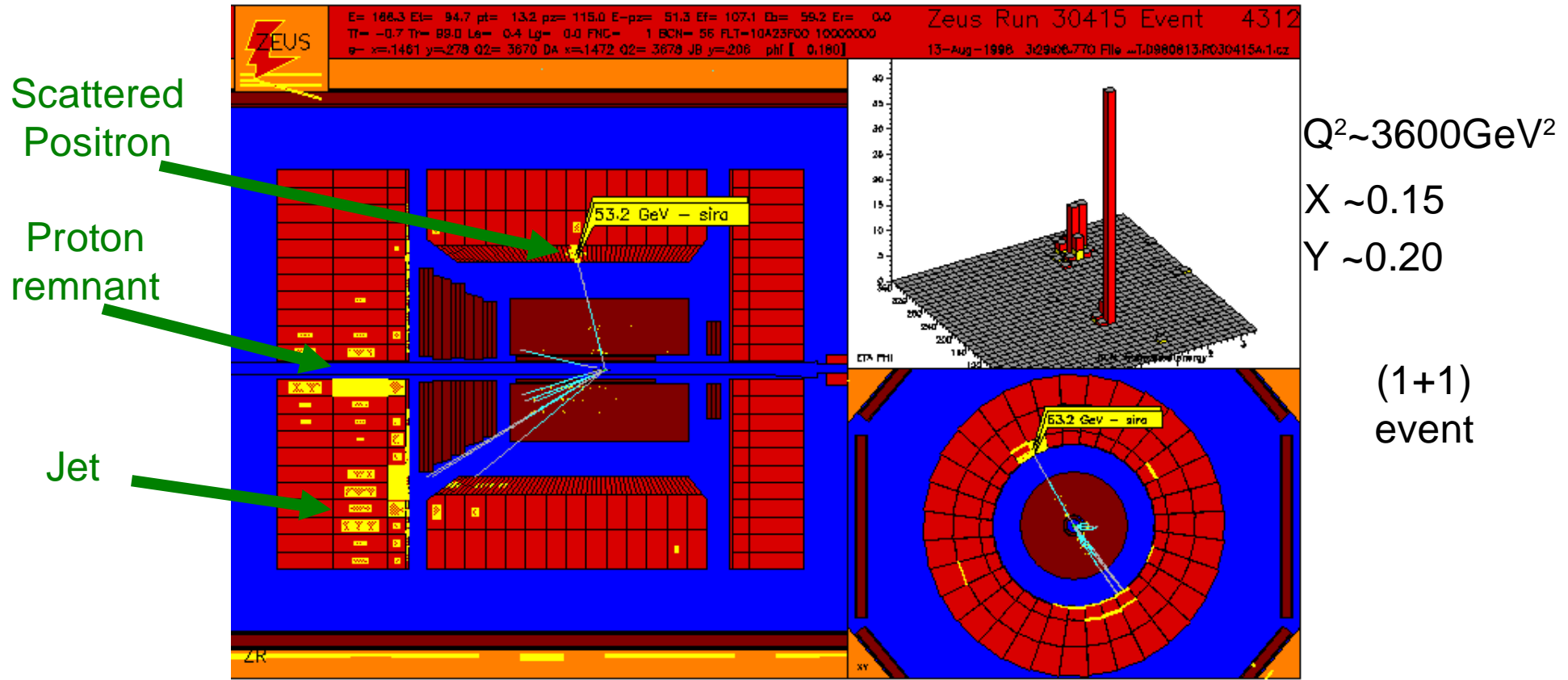


Before: $E-p_z = 2E_{\text{beam}} = 55\text{GeV}$

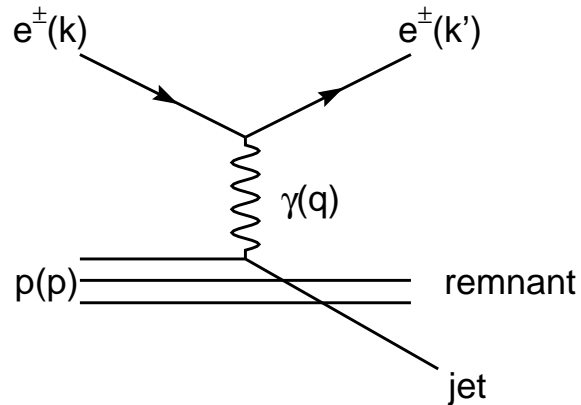
Unless energy escapes down rear beam pipe, $E-p_z$ after collision will be near $2E_{\text{beam}}$ for interesting physics at the nominal interaction point



Deep Inelastic Scattering Event



DIS Cross Section



$$Q^2 = -q^2 = -(k - k')^2$$

$$x = \frac{Q^2}{2p \cdot q} \quad y = \frac{p \cdot q}{p \cdot k} \quad Q^2 = sxy$$

Neutral Current: $e^\pm p \rightarrow e^\pm + X$

$$\frac{d^2 \sigma}{dx dQ^2} = \frac{2\pi\alpha_{em}}{xQ^4} [Y_+ F_2 \mp Y_- xF_3 - y^2 F_L] \quad Y_\pm = 1 \pm (1-y)^2$$

F_2 due to photon exchange with spin $\frac{1}{2}$ partons.

Related to quark densities f_i

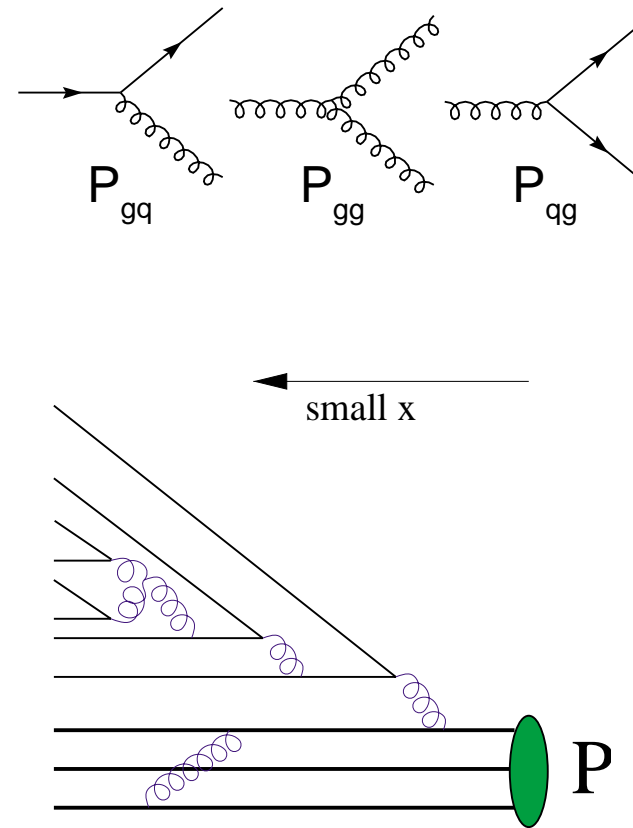
$$F_2(x) = \sum_i x Q_i^2 f_i(x)$$

F_3 contribution due to Z exchange.

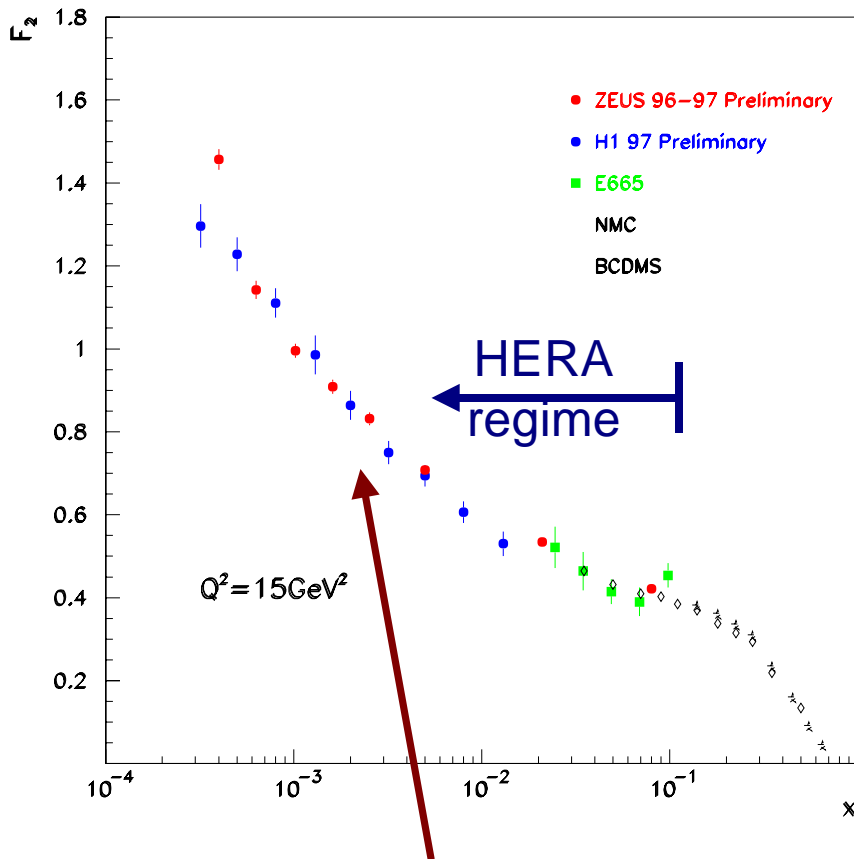
F_L contribution due to exchange of longitudinally polarized photons.

The Role of the Gluon

Splitting Functions from QCD

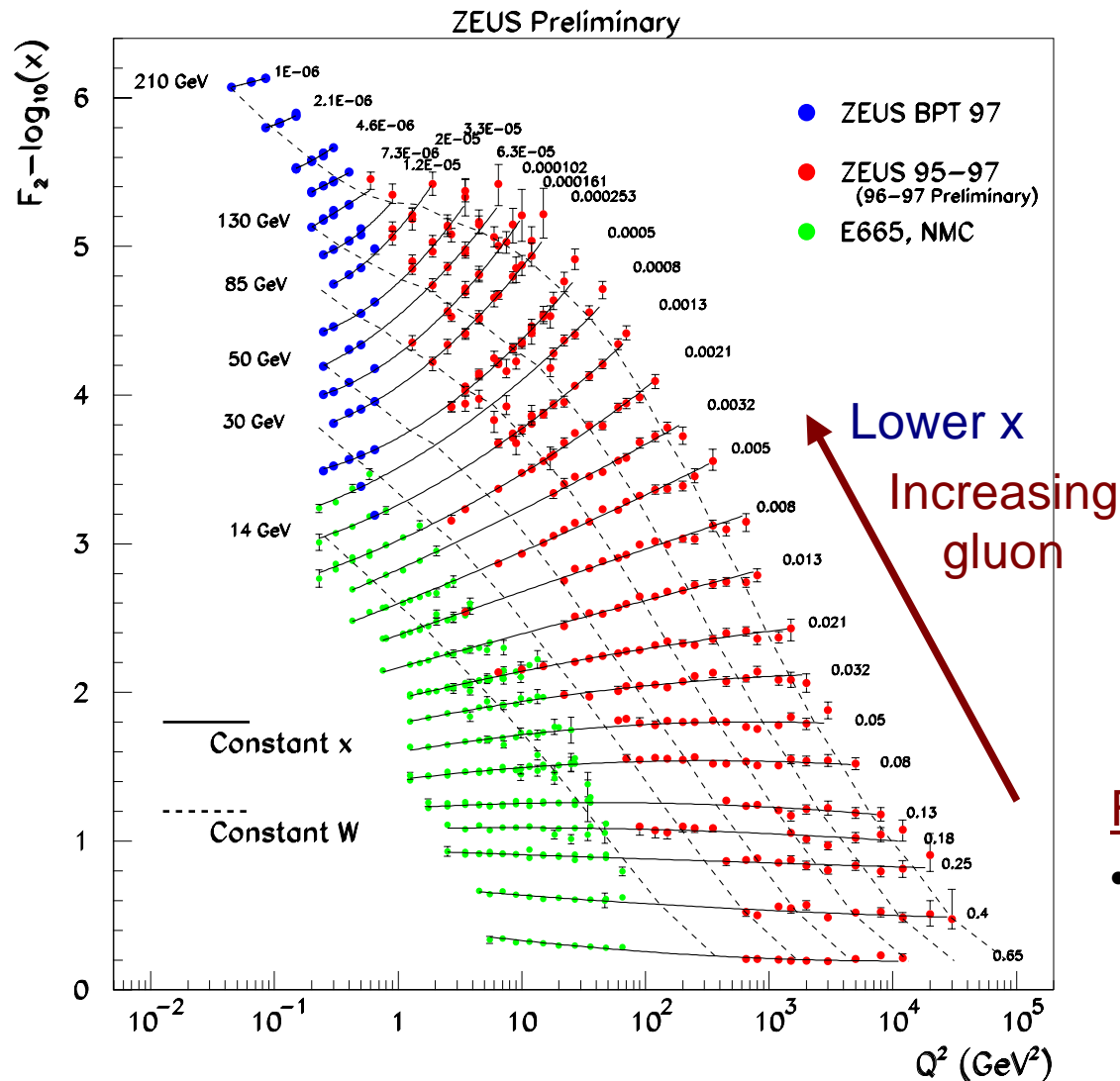


Gluon-driven increase of small x quarks is reflected in F_2



HERA discovery!
Strong rise of F_2 at low x !

Scaling Violation and the Gluon



Scaling Violation

- F_2 has a Q^2 dependence due to gluon

$$F_2(x) \rightarrow F_2(x, Q^2)$$

- More significant at smaller x

F_2 scaling violation \rightarrow gluon density

- QCD evolution equations (Altarelli-Parisi) predict

$$g(x, Q^2) \sim dF_2(x, Q^2) / d \log Q^2$$

Parton Density Functions

Several different groups make global fits to DIS structure function data

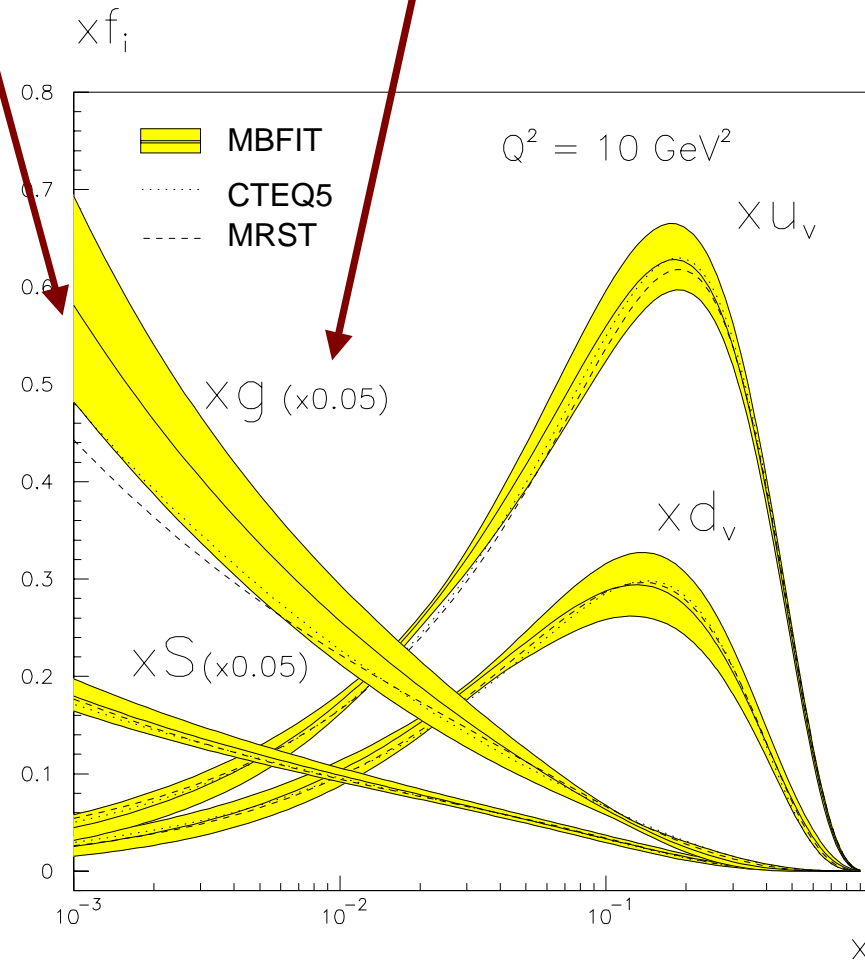
- Parton Density Functions (PDFs)
- CTEQ, MRST, GRV, MBFIT
- needed by Tevatron and LHC

Gluon density extracted indirectly from scaling violation of F_2

- $g(x, Q^2) \sim dF_2(x, Q^2) / d \log Q^2$
- relatively large uncertainty on $g(x, Q^2)$
- A measurement with **direct** sensitivity to the gluon would be nice...

16% uncertainty at $x \sim 10^{-3}$

Factor of 20!



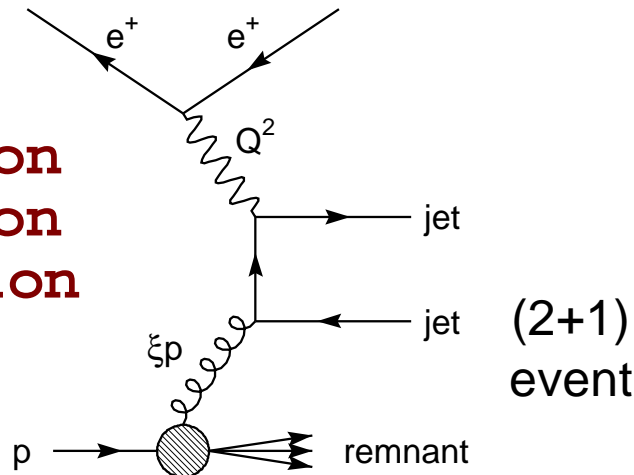
Dijet Production at HERA

Why settle for one jet,
when you can have two!

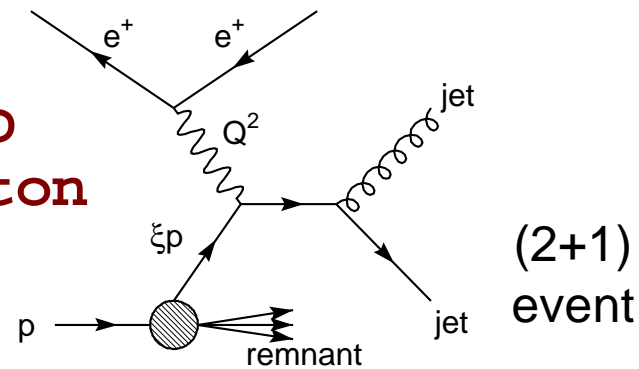


HERA \sqrt{s} provides
high E_T final
hadronic state

**Boson
Gluon
Fusion**



**QCD
Compton**



DIS variables still apply

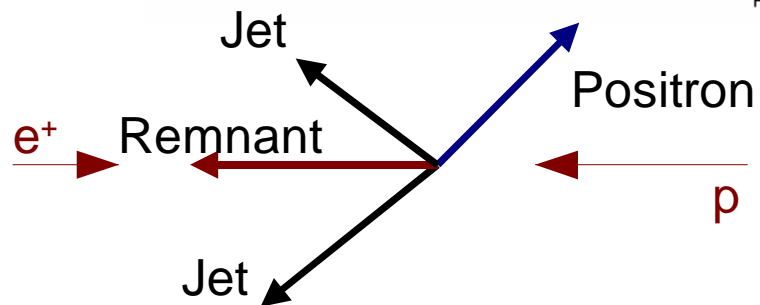
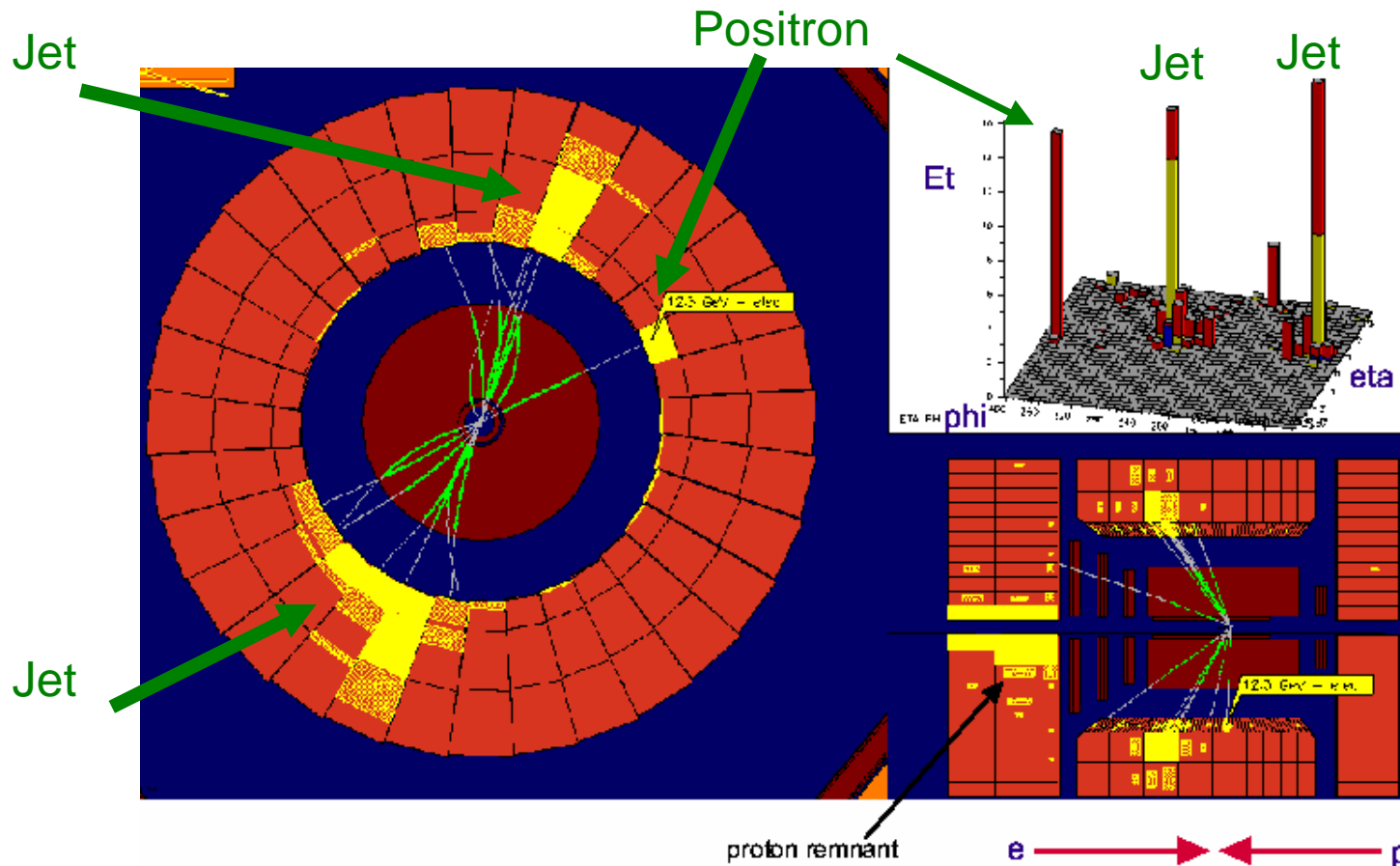
$$Q^2 = -q^2 = -(k - k')^2$$

$$y = \frac{p \cdot q}{p \cdot k} \quad x = \frac{Q^2}{2p \cdot q}$$

But now the momentum fraction of
the incident parton (at LO) is

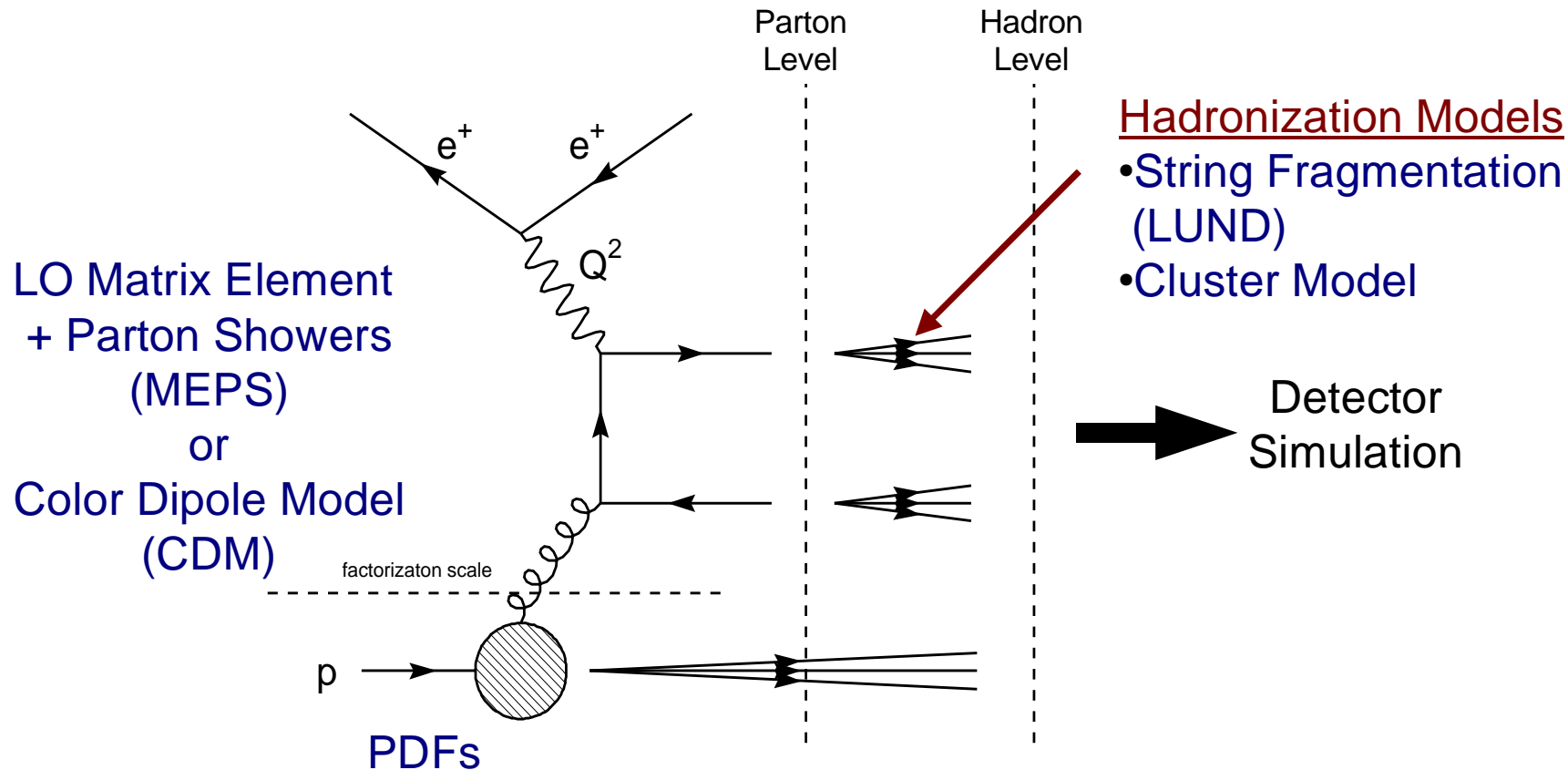
$$\xi = x \left(1 + \frac{M_{jj}^2}{Q^2} \right) \quad M_{jj} = \text{dijet mass}$$

Dijet Event at ZEUS



How is this type of event understood within the context of pQCD? →

Leading Order Monte Carlo Models



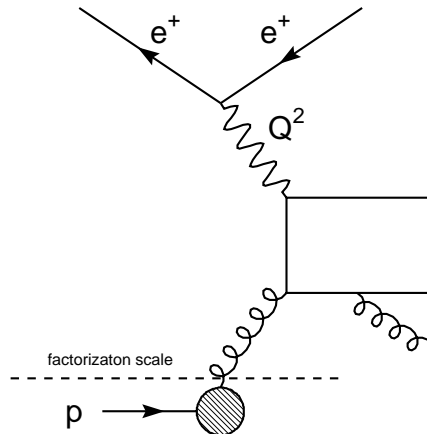
Programs

- [LEPTO](#) (MEPS+LUND)
- [ARIADNE](#) (CDM+LUND)
- [HERWIG](#) (MEPS+CLUSTER)

LO models used only for

- detector corrections
- hadronization corrections

NLO Calculations



NLO Matrix Elements

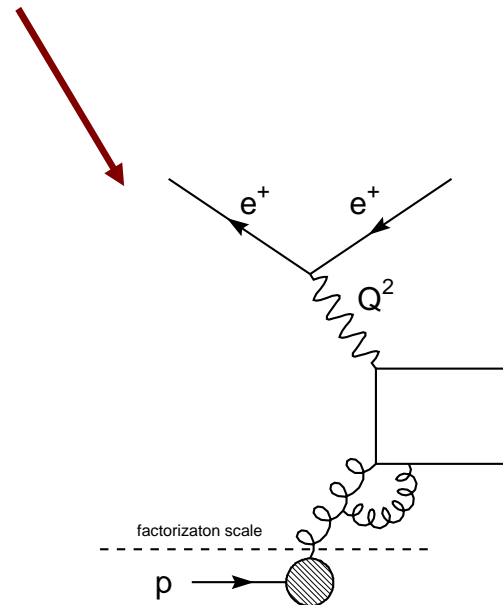
- Three parton final states
- Large improvement over LO
- Soft/collinear and virtual loop divergences cancel

Programs for DIS

- DISINT (subtraction method)
- DISASTER++ (subtraction method)
- MEPJET (phase space splicing method)
- JETVIP (phase space splicing method)

Issues

- Renormalization scale uncertainty (next slide)
- Hadronization effects (discussed later)
 - non-perturbative: Partons \rightarrow Hadrons
 - NLO calculations provide only 3-parton final states



Renormalization Scale

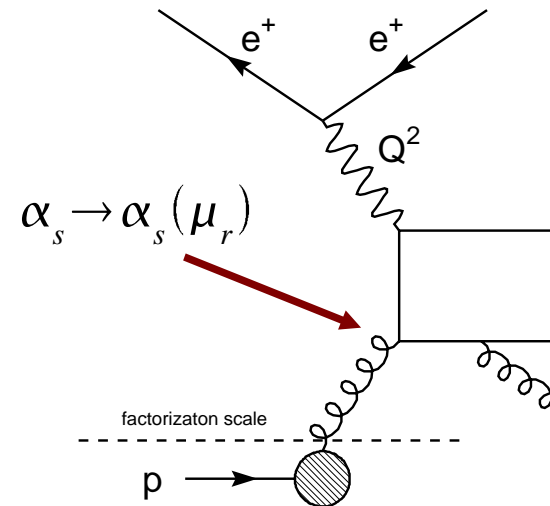
Renormalization scale (μ_r): Scale at which the strong coupling constant is evaluated

Factorization scale (μ_f): Scale at which the parton densities are evaluated

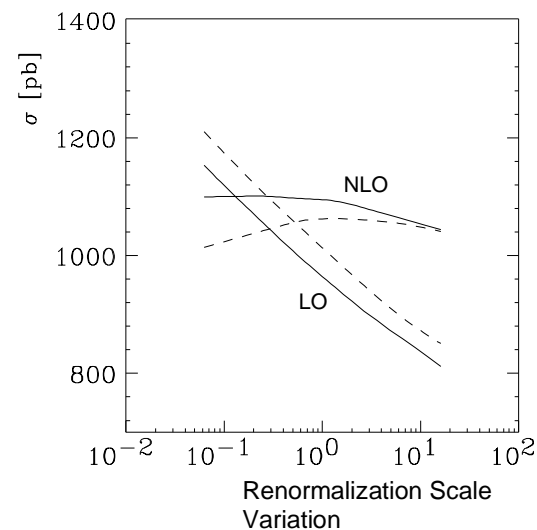
$d\sigma/d\mu_r=0$ only for all-order perturbation

Otherwise uncertainty in final cross section

- Uncertainty due to **factorization scale** is typically small (<5% for this analysis)
- Uncertainty due to **renormalization scale** can be large (>50%) even at NLO



Example: Inclusive DIS Cross Section



Renormalization Scale Choices and Resulting Uncertainty

Choices for renormalization scale

- Typically choose the hardest scale available
- In single jet DIS: $\mu_r^2=Q^2$
- In dijet events, jet E_T^2 can be larger than Q^2
 - $\mu_r^2=E_T^2/4$ also reasonable
 - E_T =sum of jet E_t s
 - $E_T^2/4 \sim$ square of mean dijet E_T

Estimate of renormalization scale uncertainty

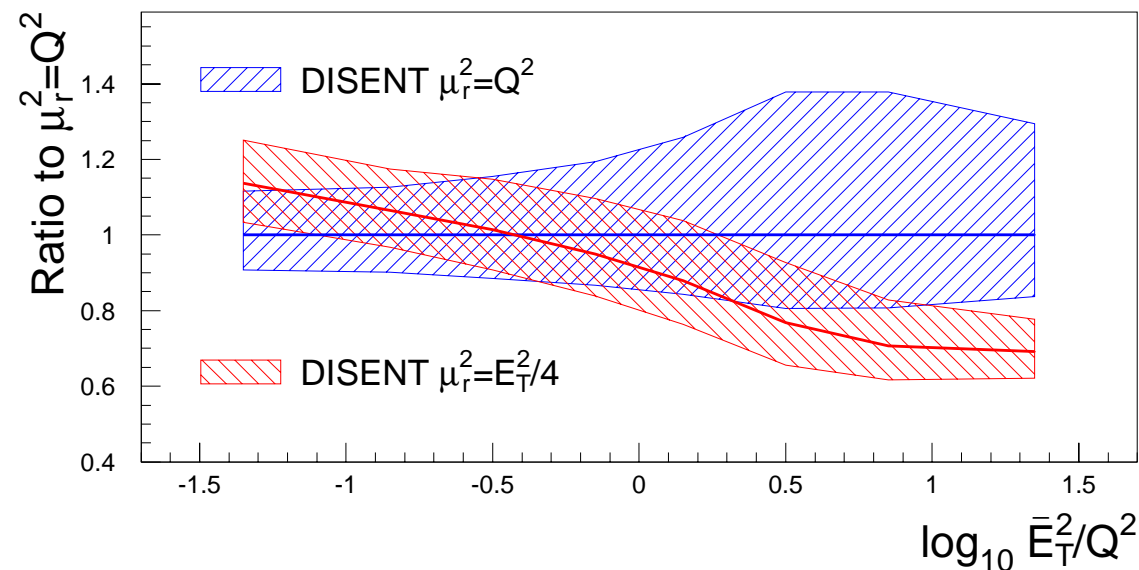
- vary μ_r by factor of 2 (conventional)

NLO Calculation of Inclusive Dijet Cross Section

- mean jet $E_T > 6.5\text{GeV}$
- $Q^2 > 10\text{GeV}^2$

Uncertainty due to renormalization scale can be large

- at least 40%

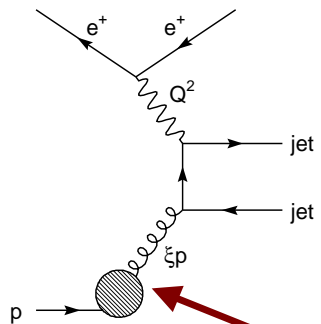


Dijets \rightarrow Gluon Density

BGF process dominates at low Q^2

BGF contribution directly proportional to the gluon density

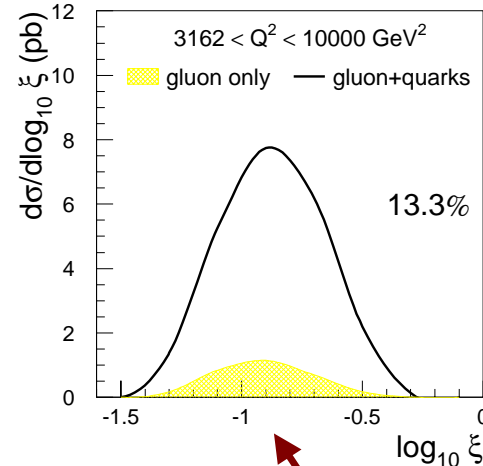
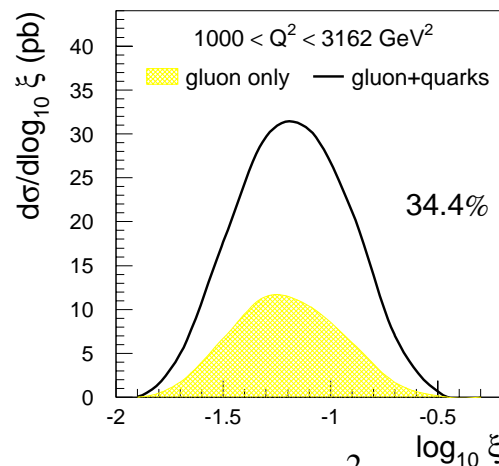
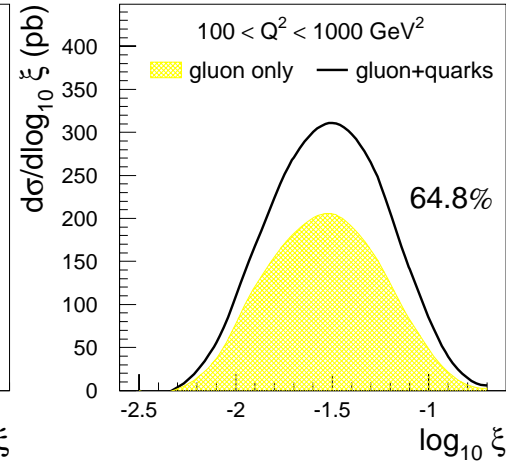
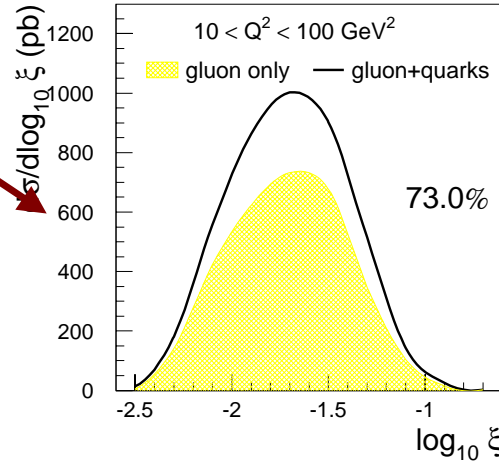
$$\sigma^{2+1} \sim \hat{\sigma}_{BGF} \cdot g(x, Q^2)$$



Use matrix elements from pQCD

Probe gluon in proton

Opportunity for a direct extraction of the gluon density.

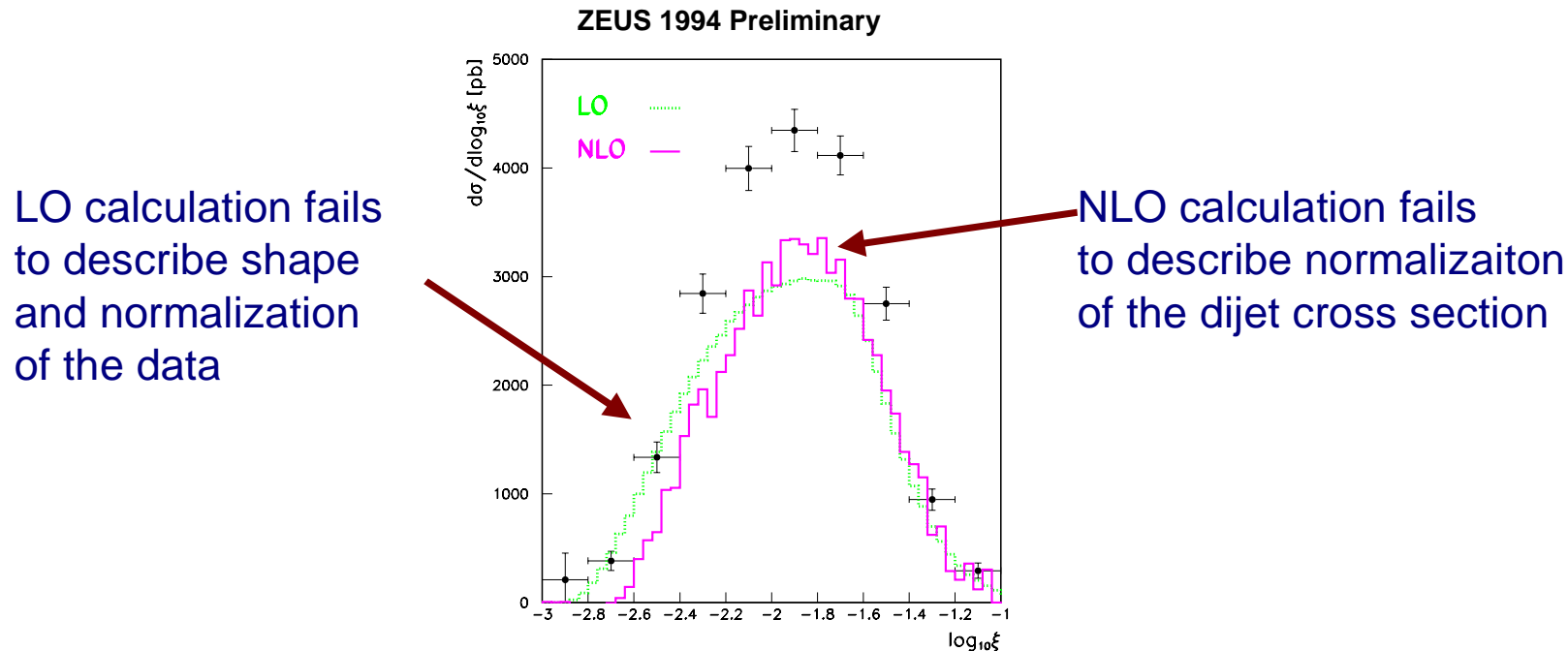


$$\xi = x \left(1 + \frac{M_{jj}^2}{Q^2} \right)$$

Momentum fraction of incident parton

At high Q^2 sensitive to high x quarks

Previous ZEUS Dijet Results



ZEUS 1994 Dijet Preliminary Result

- Lower Q^2 region $7 < Q^2 < 100$
 - maximum sensitivity to gluon
- Compared dijet cross section with NLO pQCD calculations
 - Normalization difference of ~40%
 - Shape of distributions accurately described

Result inspired new set of goals

- Measure cross section over large Q^2 range
- Understand normalization discrepancy with NLO calculations
- Quantify measurement and theoretical uncertainties

Analysis: Data Selection

HERA 1996 and 1997 running period

- 820 GeV protons — 27.5 GeV positrons

ZEUS data sample

- integrated luminosity: **38.4 pb⁻¹**

Cross Section Definition

$10 < Q^2 < 10000 \text{ GeV}^2$

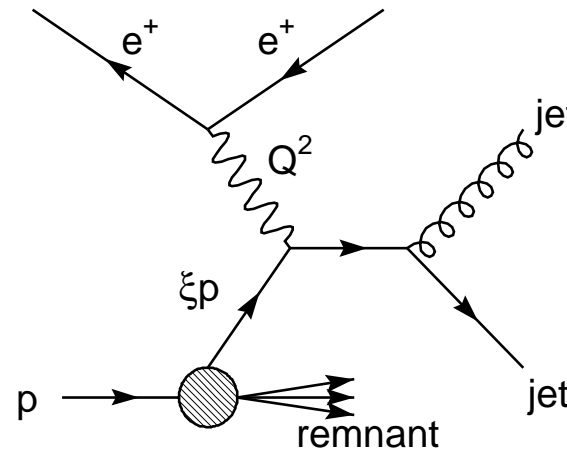
- acceptance understood above 10 GeV^2
- statistics limited above 10000 GeV^2

$y > 0.04$ and $E_e > 10 \text{ GeV}$

- detector efficiency and background reduction

Jets defined by inclusive-mode k_T -algorithm run in Breit Frame (next slides)

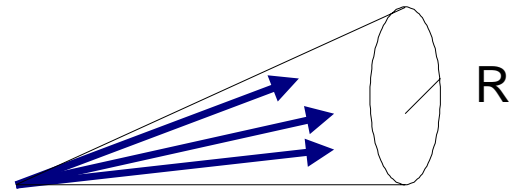
- Lab Frame: Jet $E_T > 5 \text{ GeV}$, $|\eta| < 2.0$
 - well understood acceptance region
- Breit Frame: Jet $E_{T,1} > 8 \text{ GeV}$ and Jet $E_{T,2} > 5 \text{ GeV}$
 - **asymmetric cut**: controlled region of NLO calculations (explained later)



Jet Algorithms

Cone algorithms

- conceptually simple
- theoretical/implementation issues
 - seed requirements
 - infrared unsafe at NNLO

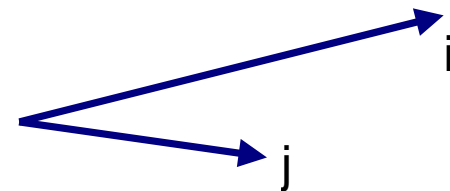


Maximize E_T within
a cone of radius R

Inclusive-mode k_T -algorithm (KTCLUS)

- No known theoretical or implementation issues
 - infrared safe
 - seeding not necessary
- Smaller hadronization corrections in some regions

$$d_i = E_{T,i}^2$$
$$d_{i,j} = \min\{E_{T,i}^2, E_{T,j}^2\}(\Delta\eta^2 + \Delta\phi^2)/R^2$$



Combine i and j if d_{ij} is
smallest of $\{d_i, d_{ij}\}$

Breit Frame

Breit Frame axis is γ -proton axis

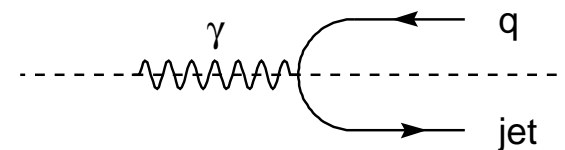
Exchanged γ 4-mom has only a z-component

Experimentally: constructed from measured 4-mom of scattered positron

Example: single jet DIS (Quark Parton Model event)

- Struck quark rebounds with equal and opposite momentum
 - "Brick Wall" Frame
- Jet has no E_T

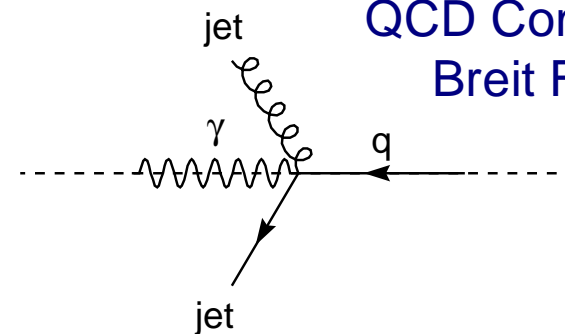
QPM event in Breit Frame



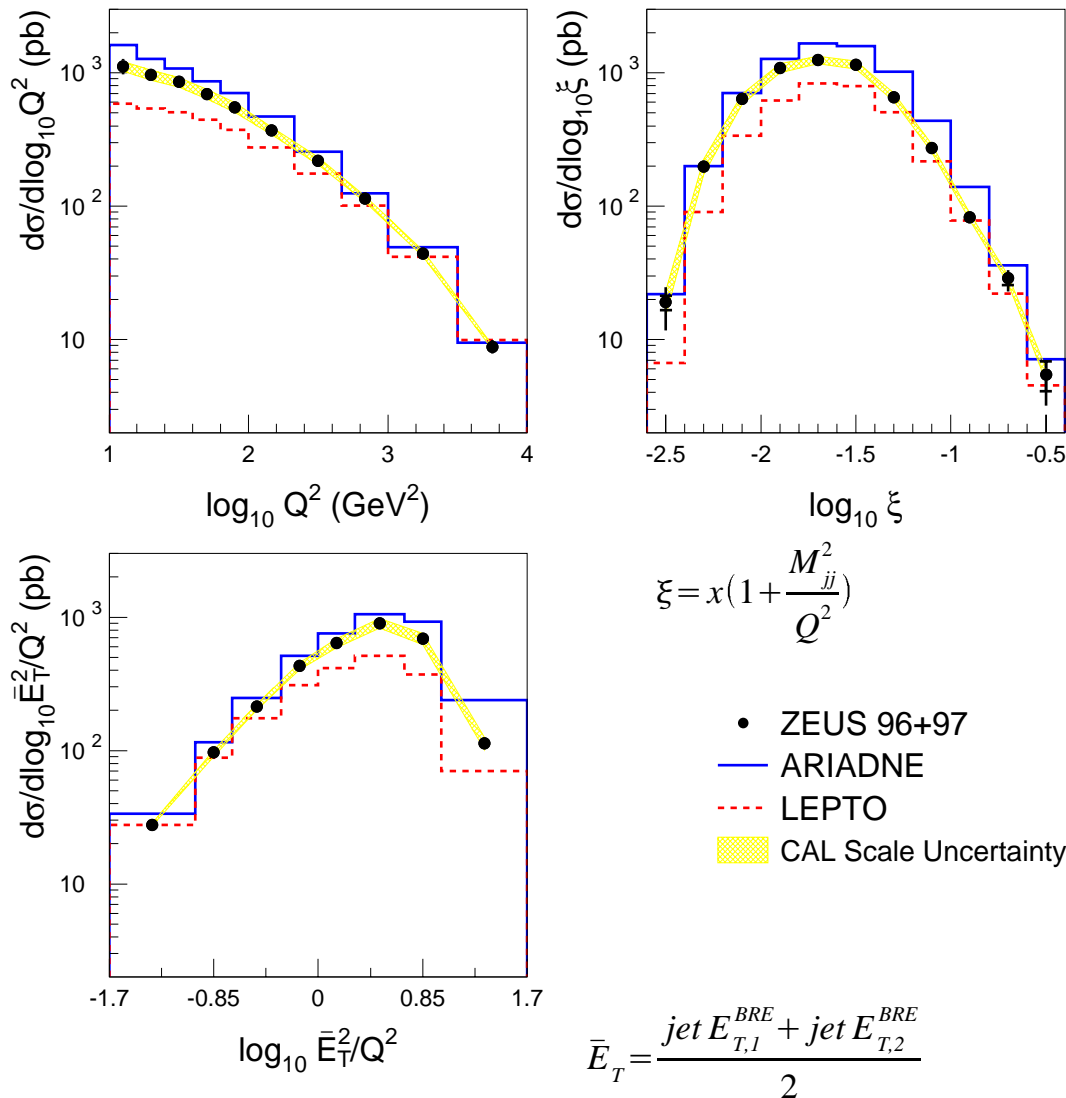
Example: Dijet in DIS (QCDC or BGF)

- Jets have balanced E_T in Breit frame

QCD Compton in Breit Frame



Comparison with LO Models



Inclusive Dijet Cross Section

$$10 < Q^2 < 10000 \text{ GeV}^2$$

$$y > 0.04, E_e > 10 \text{ GeV}$$

asymmetric jet E_T cut: 5 GeV, 8 GeV

$$-2.0 < \text{jet } \eta < 2.0$$

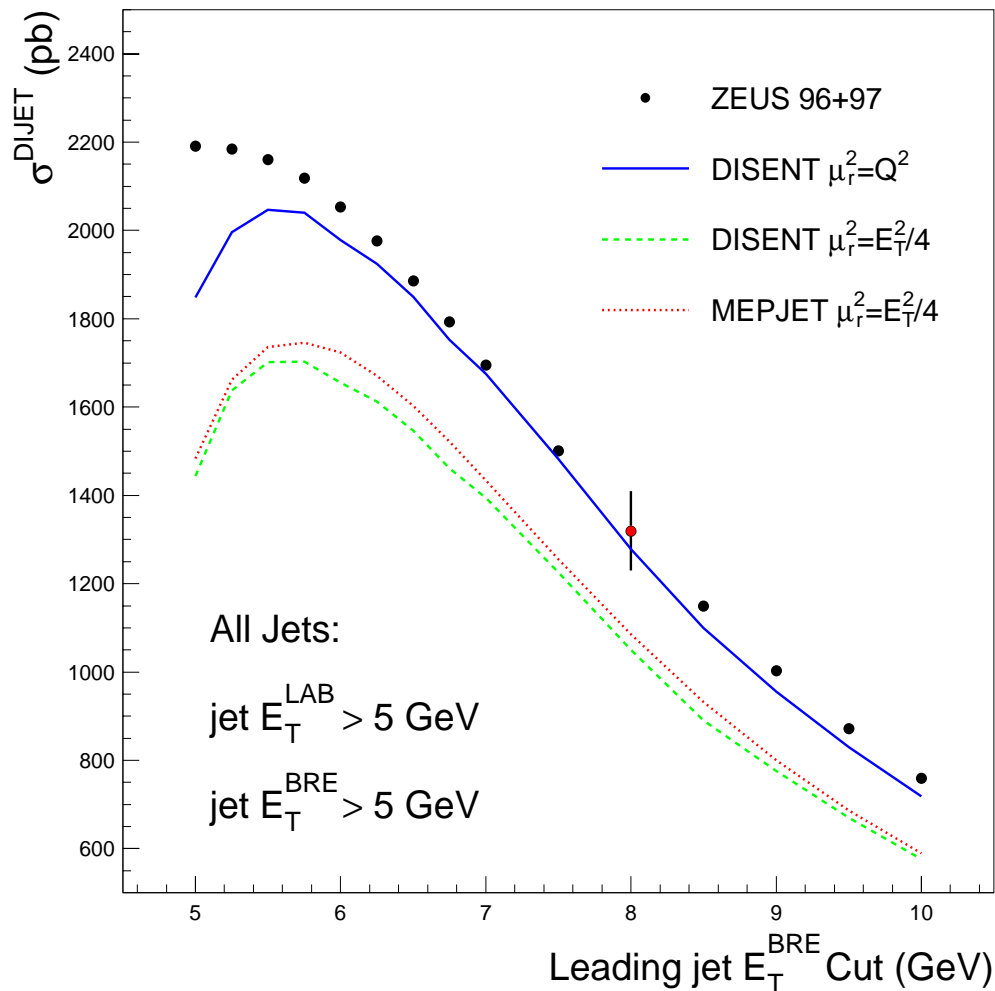
- Data corrected for detector effects to hadron level

LO MC models fail to describe normalization

- Attempts to model higher order effects inadequate
 - parton showers and CDM
- Large renormalization scale dependence

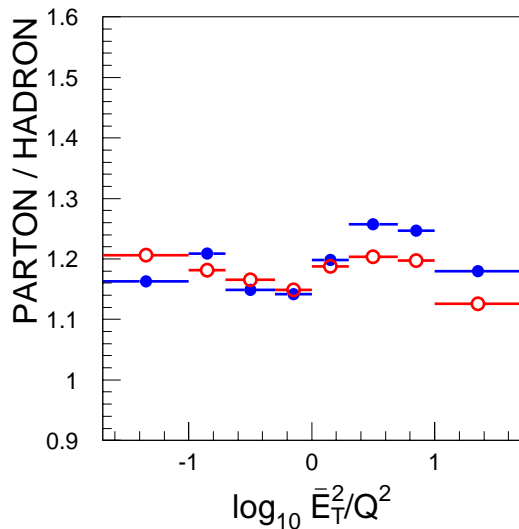
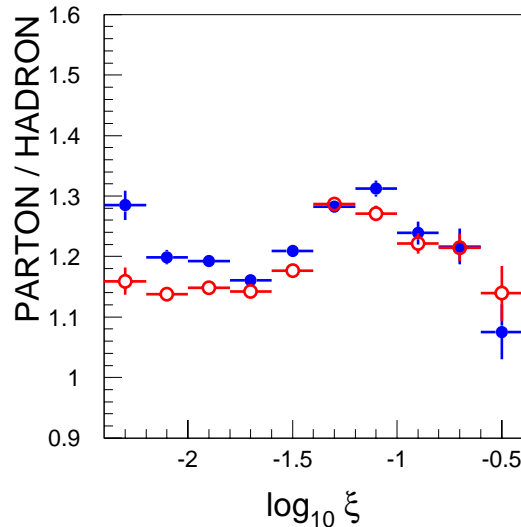
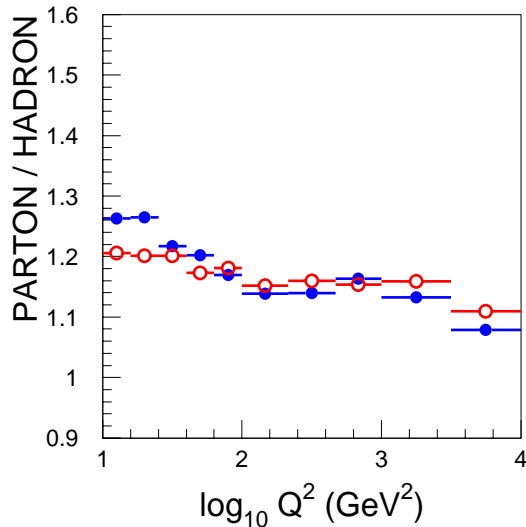
Comparing with NLO Calculations

Total Inclusive Dijet Cross Section



- NLO calculations MEPJET and DISENT behave unphysically near the symmetric cut
 - Reduced phase space for 3-parton final states that cancel negative-weight 2-parton final states
 - Asymmetric jet cut of $E_{T,1} > 8 \text{ GeV}$, $E_{T,2} > 5 \text{ GeV}$ avoids sensitive region
- Large difference due to choice of renormalization scale
- 40% normalization difference between NLO and ZEUS 1994 results understood!

Hadronization Effects



$$\xi = x \left(1 + \frac{M_{ij}^2}{Q^2} \right)$$

• ARIADNE: CDM + LUND

• LEPTO: MEPS + LUND

$$\bar{E}_T = \frac{\text{jet } E_{T,1}^{BRE} + \text{jet } E_{T,2}^{BRE}}{2}$$

- NLO calculations do not include hadronization models
- Use LO MC models to **estimate** non-perturbative hadronization effects
- Apply corrections from LO MC model estimates to NLO calculations
 - Ariadne used by default

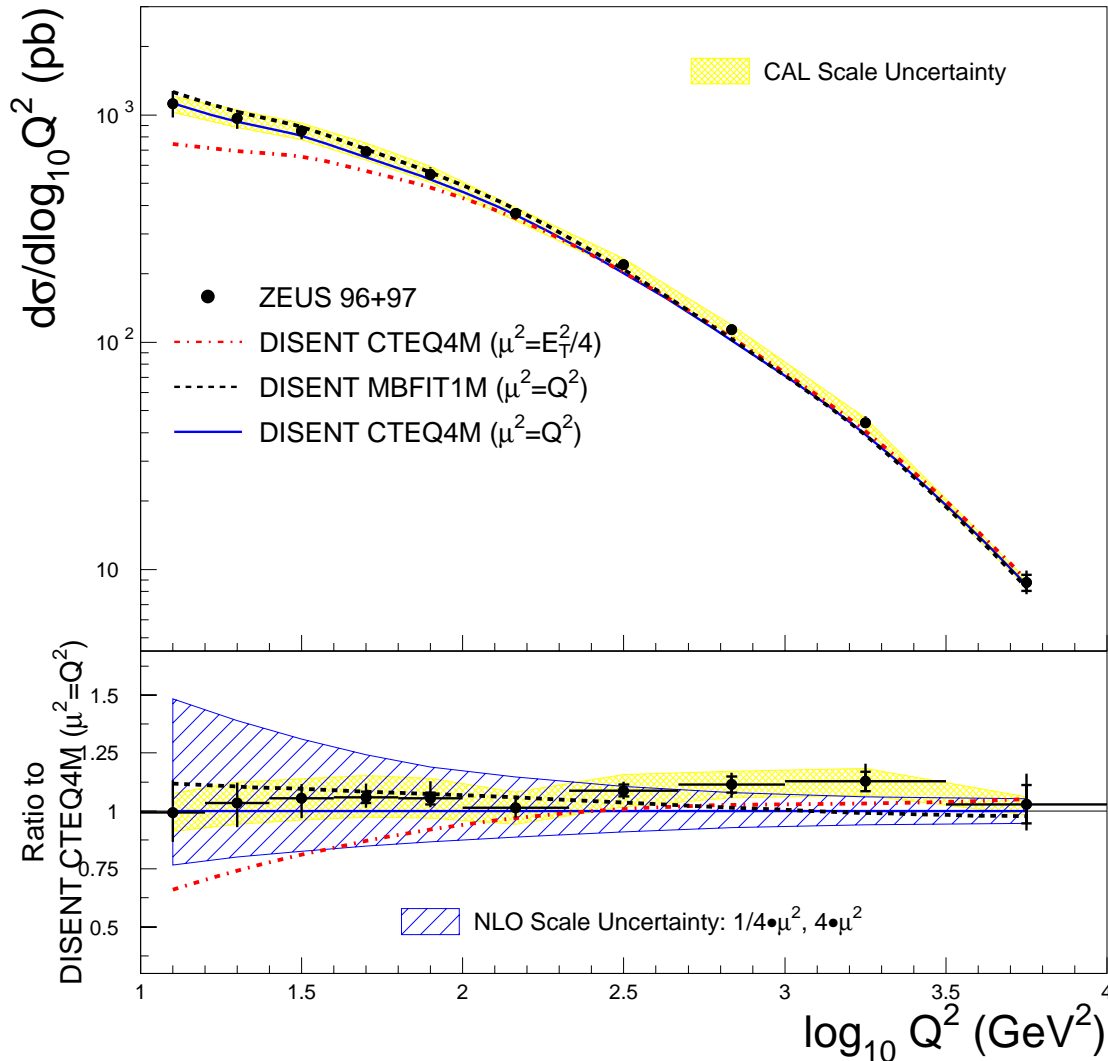
• Hadronization corrections vary between 10% and 30%

• Estimates from Ariadne and Lepto LO models vary typically by 5%, and no more than 10%

- additional theoretical uncertainty

Inclusive Dijet Cross Section vs Q^2

Asymmetric jet cut: 5, 8GeV; $E_e > 10$, $y > 0.04$



NLO Comparison

Success for pQCD!

Within NLO scale uncertainty estimate, NLO calculations reproduce measured cross section to within 10%

- over three orders of magnitude in Q^2
- over 2 orders of magnitude in value

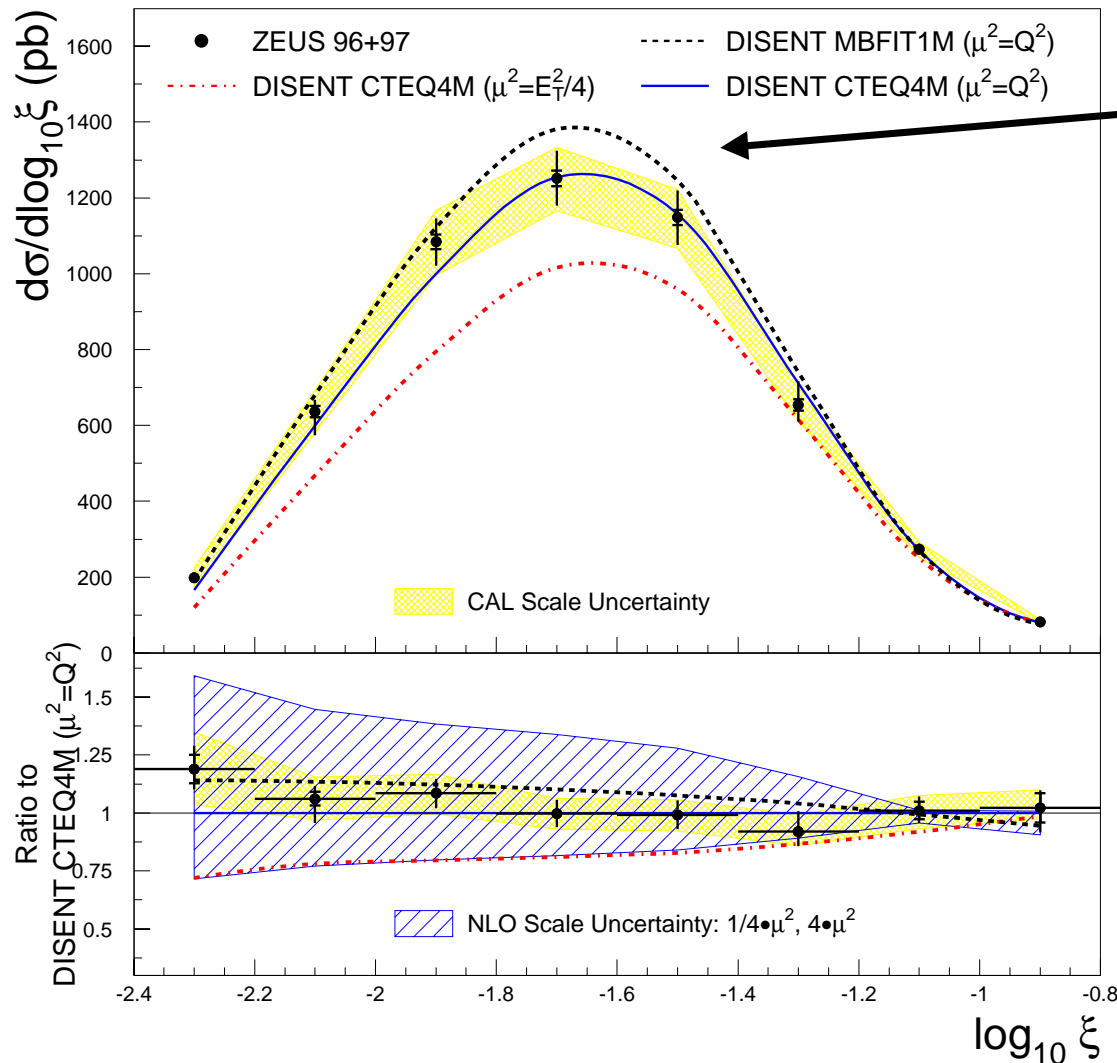
For $Q^2 < \sim 200$ measurement uncertainties less than renormalization scale uncertainty

- Need improved theoretical calculations with reduced renormalization scale dependence

Inclusive Dijet Cross Section vs ξ

Asymmetric jet cut: 5, 8GeV; $E_e > 10$, $y > 0.04$

NLO Comparison



NLO calculation shows dependence on input parton densities

- MBFIT has larger gluon

All NLO calculations are consistent with the data

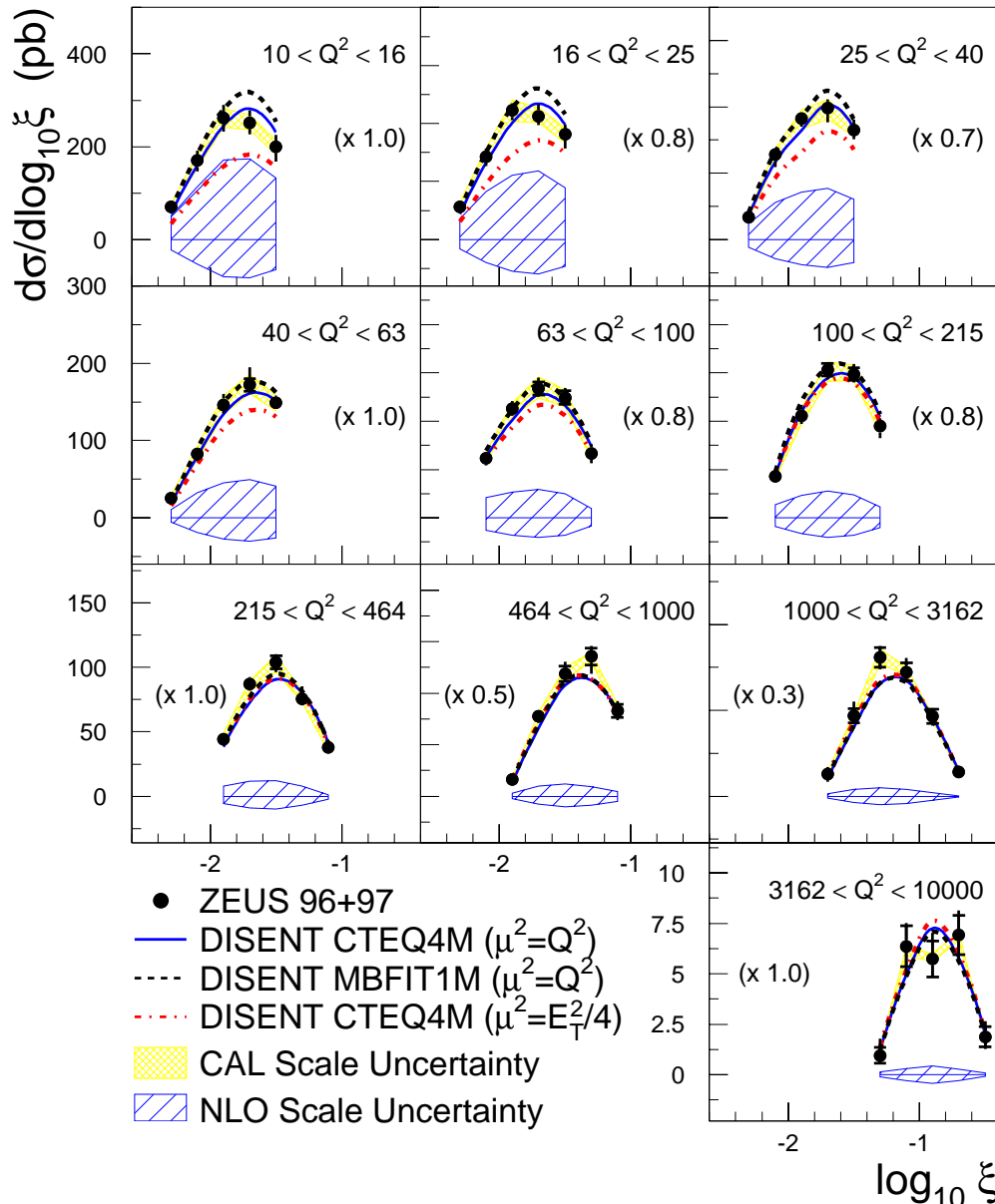
- within all uncertainties
- regardless of input parton densities

Gluon densities in current PDFs

- consistent with the data and pQCD calculations

$$\xi = x \left(1 + \frac{M_{ij}^2}{Q^2} \right)$$

Inclusive Dijet Cross Section vs ξ and Q^2



NLO Comparison

Gluon density sensitivity of NLO calculation seen for $Q^2 < \sim 200 \text{ GeV}^2$

NLO calculations converge at higher Q^2

- quark densities well constrained
- smaller renormalization scale uncertainty

For $Q^2 < \sim 200$ renormalization scale uncertainty larger than measurement uncertainty

Future improved pQCD calculations will enable use of the dijet cross section measurement for $Q^2 < \sim 200$

- used in global fits
- or used to extract gluon density directly

Conclusions on Dijet Cross Section and pQCD

Inclusive dijet cross section measured for $10 < Q^2 < 10000 \text{ GeV}^2$ and asymmetric jet E_T cuts of 5 and 8 GeV

NLO pQCD reproduces the dijet cross section within 10% over three orders of magnitude in Q^2 and over 2 orders of magnitude in value. Triumph for pQCD!

Universality of gluon density: Gluon extracted from scaling violation of F_2 can be used to describe the dijet cross section.

NLO calculations exhibit large renormalization scale dependence for $Q^2 < \sim 200 \text{ GeV}^2$; exactly where sensitivity to gluon density is largest.

- Extraction of gluon density using dijet cross section possible with future improved calculations