

Three-jet Production in Neutral Current Deep Inelastic Scattering with ZEUS Detector at HERA

Preliminary Examination

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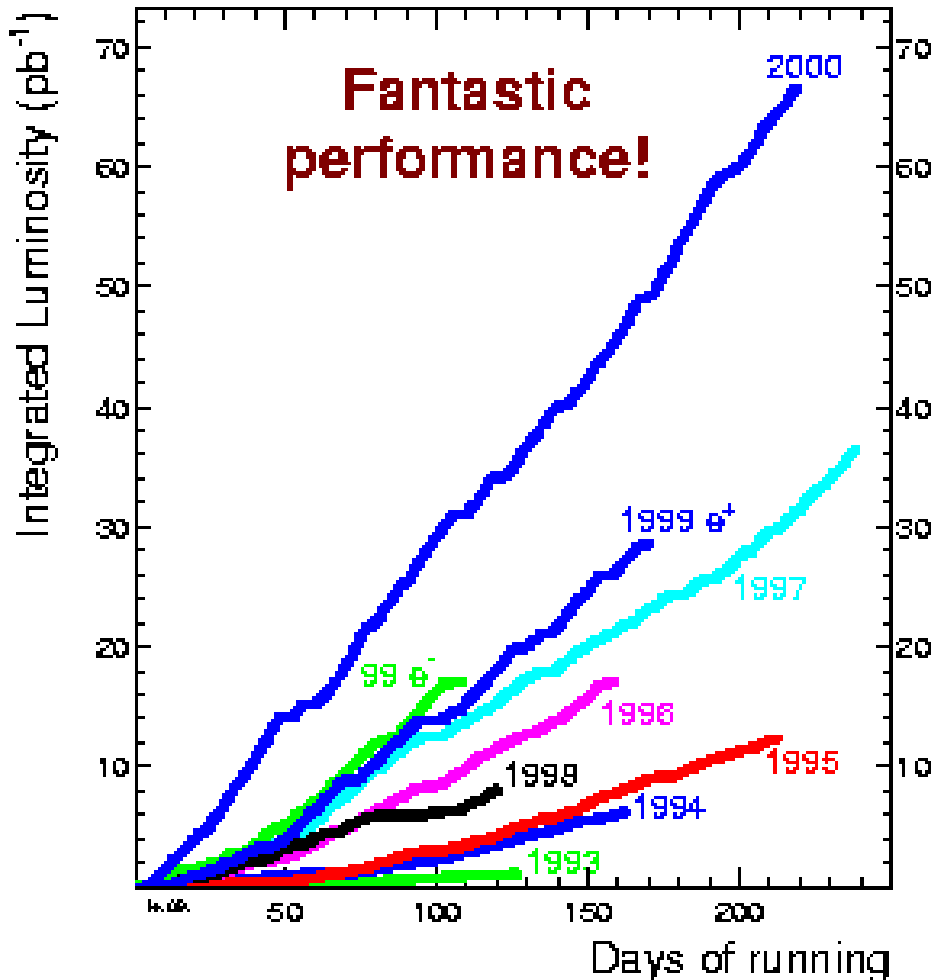
HERA



- 820/920 GeV proton
- 27.5 GeV electrons or positrons
- 300/318 GeV center of mass energy
- 220 bunches
96 ns crossing time
- Instantaneous luminosity
 $1.8 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$
- Currents:
~90mA protons
~40mA positrons

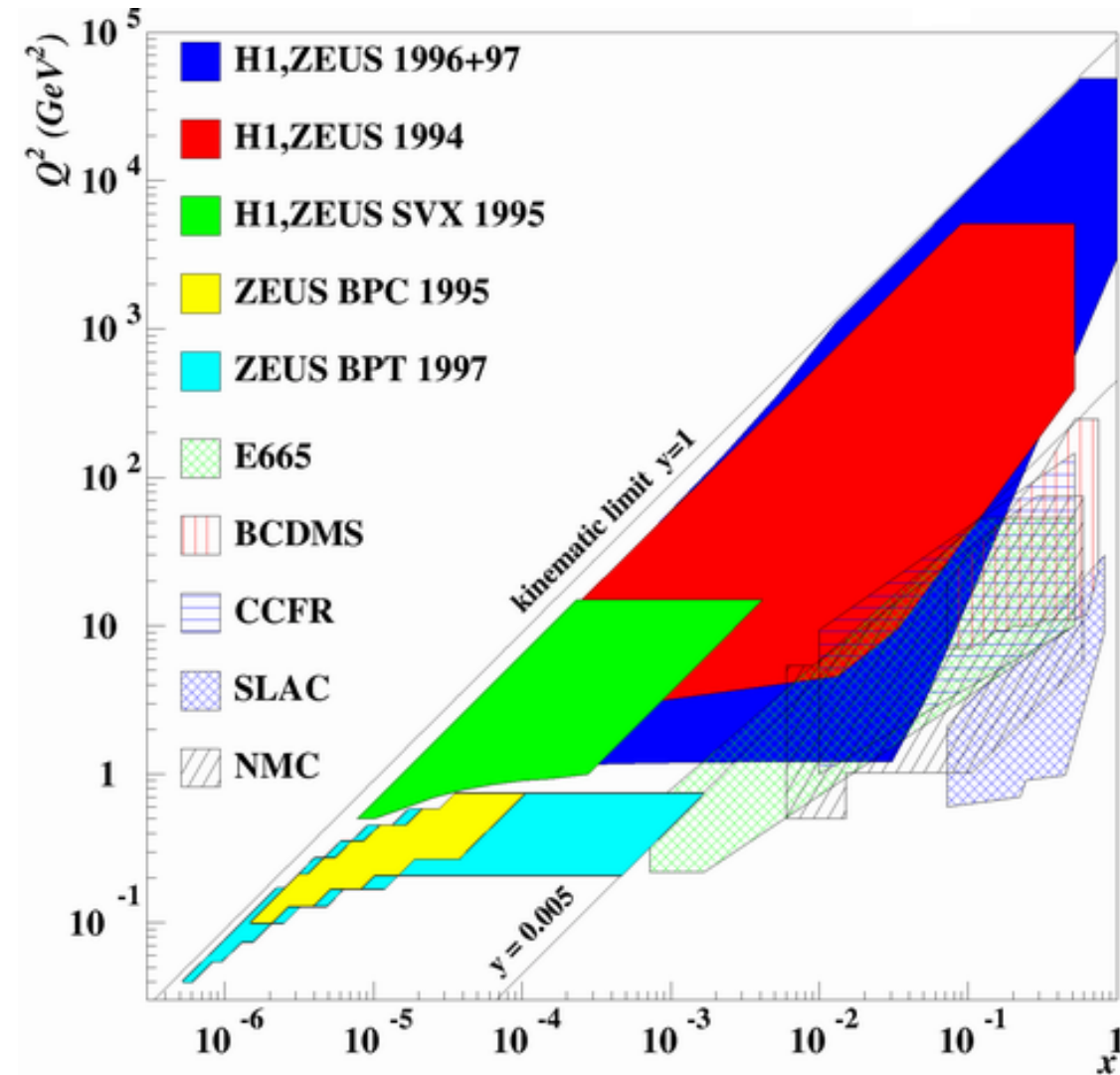
HERA Luminosity

HERA luminosity 1992 – 2000



- 820 GeV protons through 1997
- 920 GeV since 1998
- ZEUS integrated luminosity since 1992: $\sim 185 \text{ pb}^{-1}$
- Expect 1 fb^{-1} by end of 2005

HERA Kinematic Range

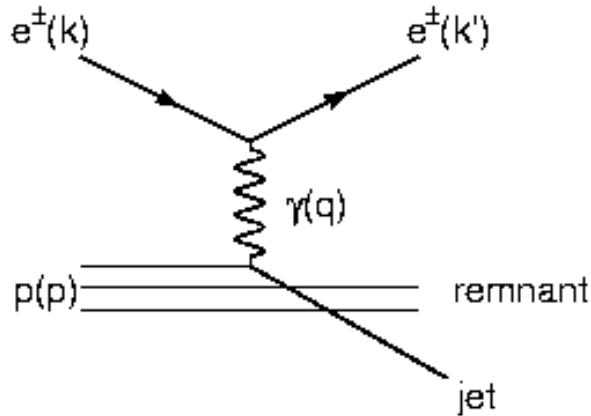


Extended kinematic region
not accessible by fixed target
experiments, with some
overlap.

H1 and ZEUS: DESY e-p
HERMES: DESY e-A
E665: Fermilab μ -A
BCDMS: CERN μ -A
CCFR: Fermilab ν -A
SLAC: many experiments e-A
NMC: CERN μ -A

Deep Inelastic Scattering

electron-proton scattering



Can also exchange Z, W^\pm

Any lepton - hadron pair

$e-p$ (HERA)

$e-A$ (SLAC)

$\nu-Fe$ (CCFR)

$\mu-A$ (E665, NMC, BCDMS)

DIS kinematic variables

$$Q^2 = -q^2 = -(k - k')^2 \quad \text{Momentum transfer}$$

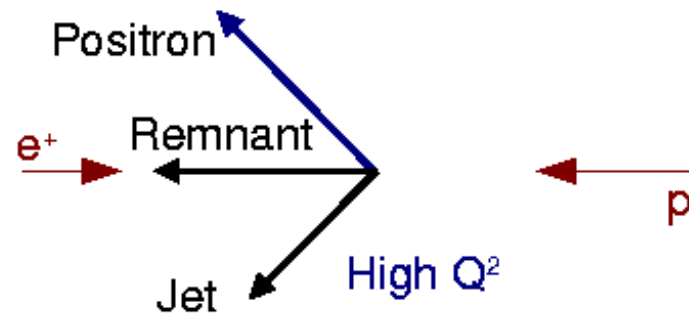
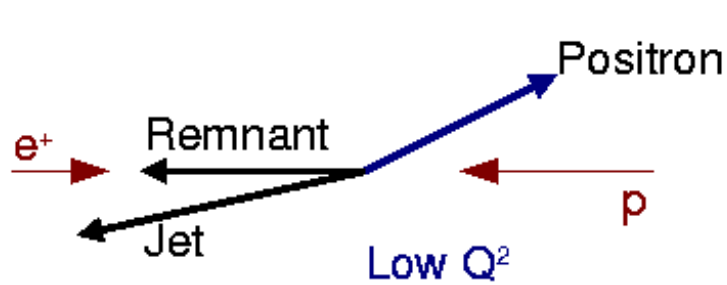
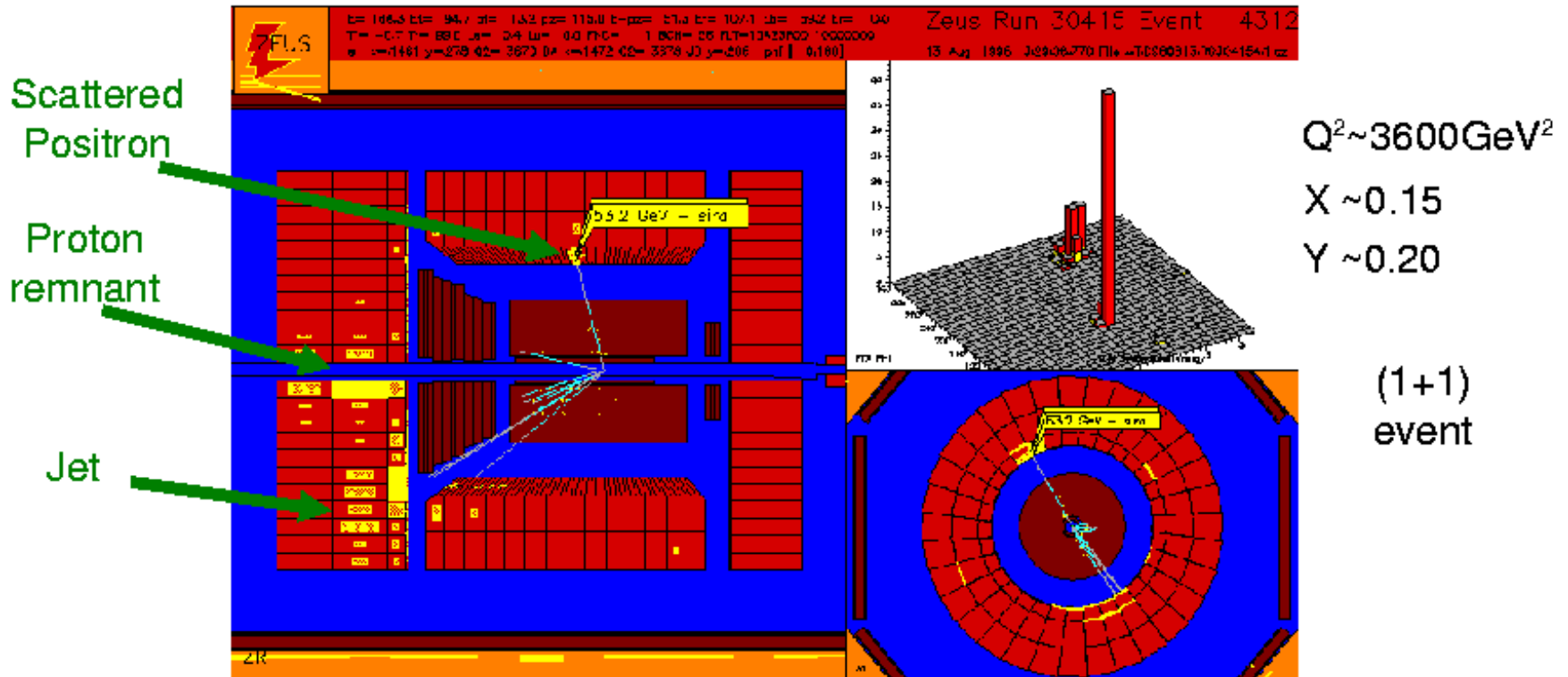
$$x = \frac{Q^2}{2p \cdot q} \quad \text{Fraction of the proton's momentum that participates in the hard scatter}$$

$$y = \frac{p \cdot q}{p \cdot k} \quad \text{Fraction of the electron's energy available in the proton's rest frame}$$

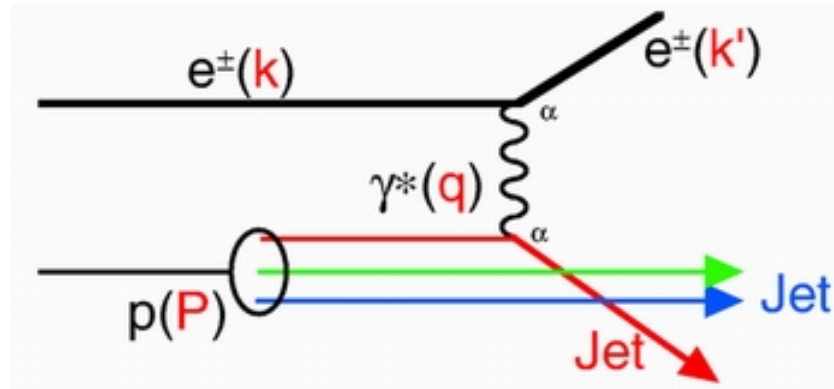
$$Q^2 = sxy \quad \text{s=center of mass energy squared}$$

$$W^2 = (p + q)^2 \quad \text{W=invariant mass of the final hadronic system}$$

e^+p DIS Event at HERA



DIS Cross Section



$$\frac{d^2\sigma(e^+p)}{dx dQ^2}(x, Q^2) = \frac{2\pi\alpha^2}{xQ^4} \left[\left(1 + (1-y)^2\right) \underline{F_2(x, Q^2)} - y^2 \underline{F_L(x, Q^2)} - (1 - (1-y)^2) x \underline{F_3(x, Q^2)} \right]$$

$F_2(x, Q^2)$: Interaction between transversely polarized photons & spin 1/2 partons ; Charge weighted sum of the quark distributions

$F_L(x, Q^2)$: Interaction between longitudinally polarized photons & the partons with transverse momentum.

$F_3(x, Q^2)$: Parity-violating structure function from Z^0 exchange.

Naïve Quark Parton Model

- Partons are point-like objects
- No interaction between the partons
- Structure function independent of Q^2

Bjorken Scaling (x) dependence:

$$F_2(x, Q^2) = \sum_{\text{quarks}} e_q^2(Q^2) \cdot (xq(x, Q^2) + x\bar{q}(x, Q^2))$$

$$F_2(x, Q^2) \rightarrow F_2(x), F_L = 0$$

QCD

Physics picture: presence of gluons

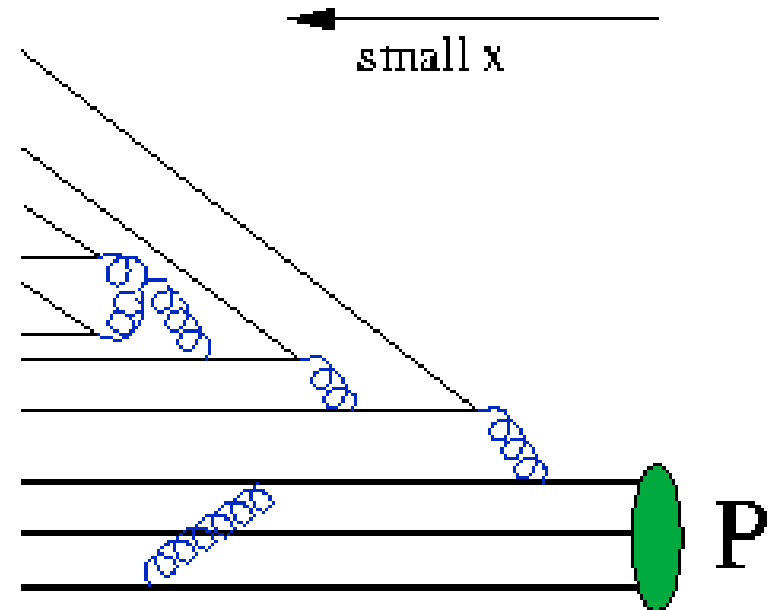
Parton Parton interactions,
mediated by **gluons**

Generate parton transverse
momentum

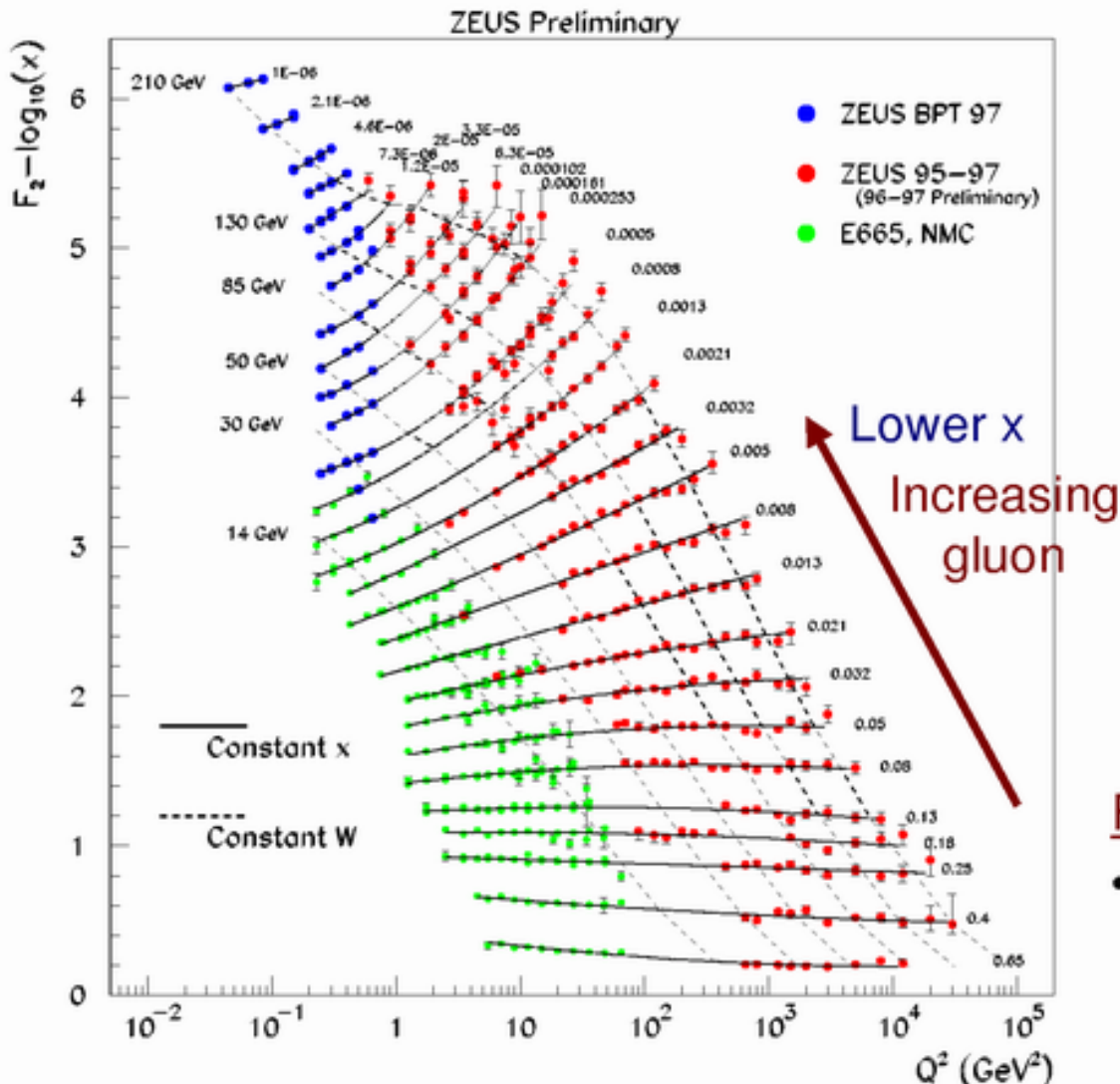
Non zero F_L

The structure functions gain
a Q^2 dependence

Scaling Violation



Scaling Violation



Scaling Violation

- F_2 has a Q^2 dependence due to gluon

$$F_2(x) \rightarrow F_2(x, Q^2)$$

- More significant at smaller x

F_2 scaling violation \rightarrow gluon density

- QCD evolution equations (Altarelli-Parisi) predict

$$g(x, Q^2) \sim dF_2(x, Q^2) / d \log Q^2$$

Parton Distribution Functions

Several different groups make global fits to DIS structure function data

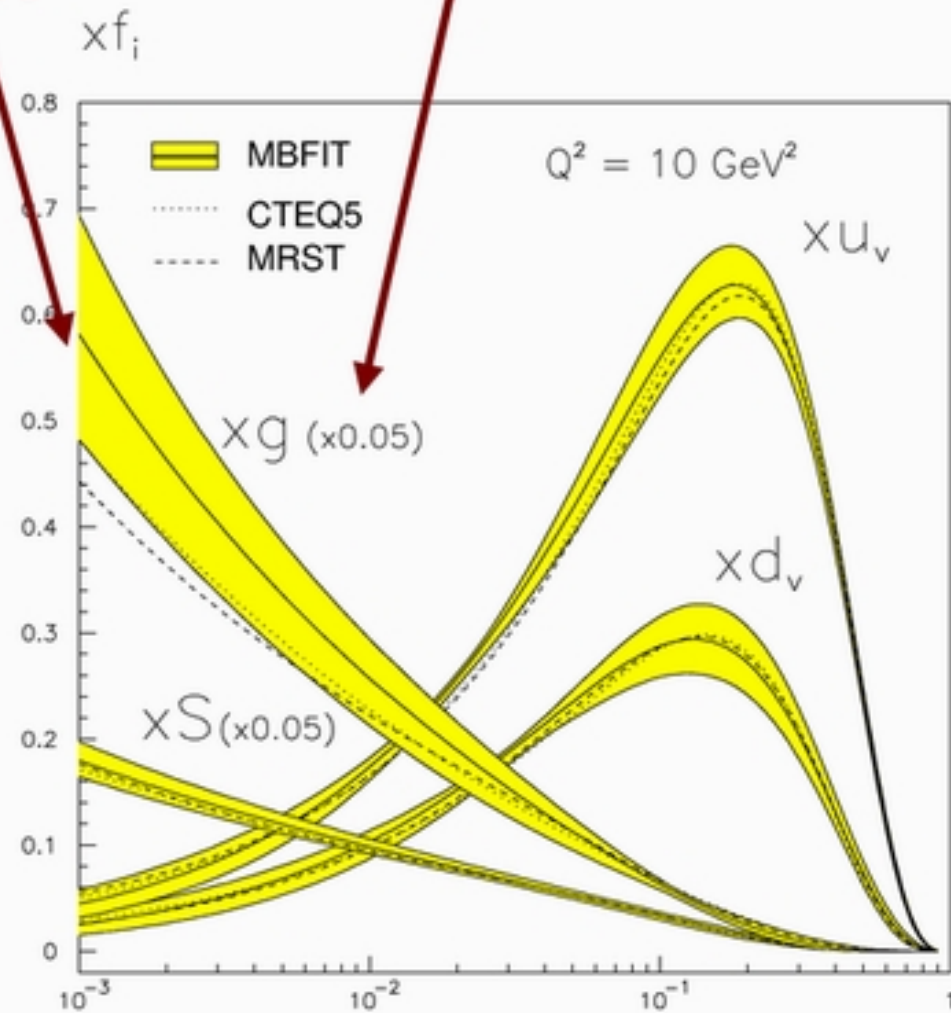
- **Parton Density Functions (PDFs)**
- CTEQ, MRST, GRV, MBFIT
- needed by Tevatron and LHC

Gluon density extracted **indirectly** from scaling violation of F_2

- $g(x, Q^2) \sim dF_2(x, Q^2) / d \log Q^2$
- relatively large uncertainty on $g(x, Q^2)$
- A measurement with **direct** sensitivity to the gluon would be nice...

16% uncertainty at $x \sim 10^{-3}$

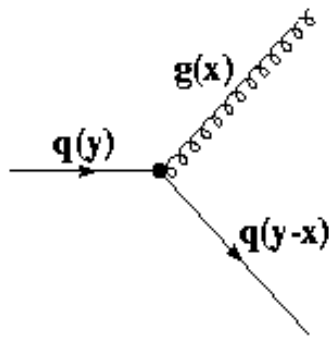
Factor of 20



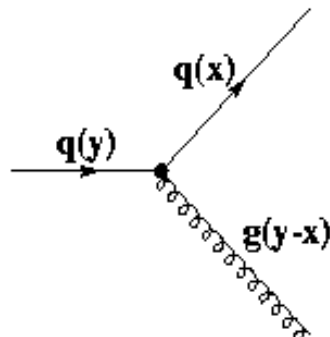
Derived from the experiment

DGLAP Evolution

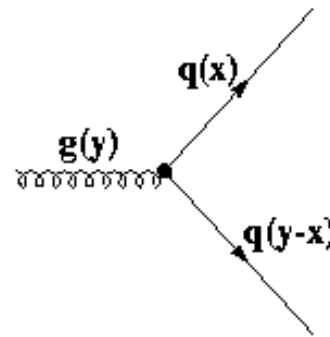
- Uses QCD splitting functions



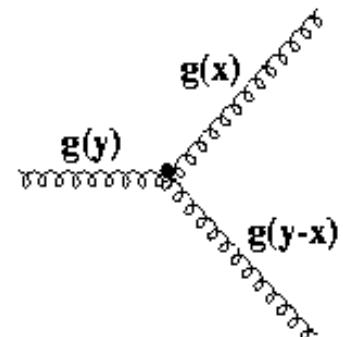
$$P_{gq}(z=x/y)$$



$$P_{qq}(z=x/y)$$



$$P_{qg}(z=x/y)$$



$$P_{gg}(z=x/y)$$

Dokshitzer-Gribov-Lipatov-Altarelli-Parisi equations describe evolution of parton densities to higher Q^2 — QCD prediction

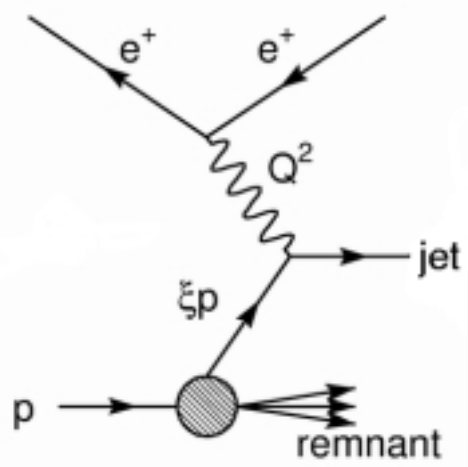
$$\frac{dq_i(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} \left[q_i(y, Q^2) P_{qq}\left(\frac{x}{y}\right) + g(y, Q^2) P_{qg}\left(\frac{x}{y}\right) \right]$$

$$\frac{dg(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} \left[\sum_i q_i(y, Q^2) P_{gq}\left(\frac{x}{y}\right) + g(y, Q^2) P_{gg}\left(\frac{x}{y}\right) \right]$$

α_s is the strong coupling constant.

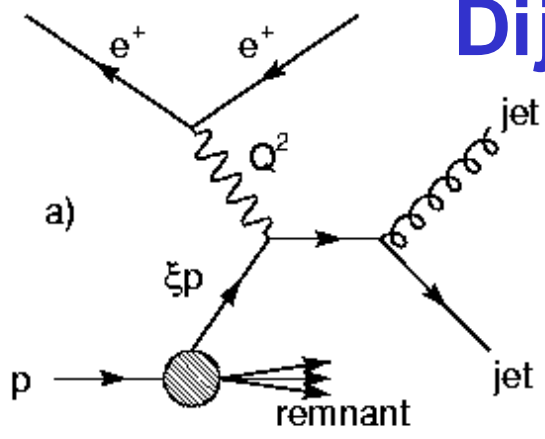
Single Jet Dijet Trijet

Single Jet

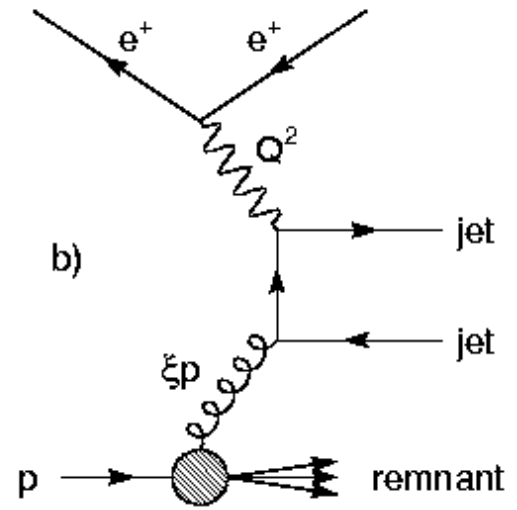


a) QCD Compton

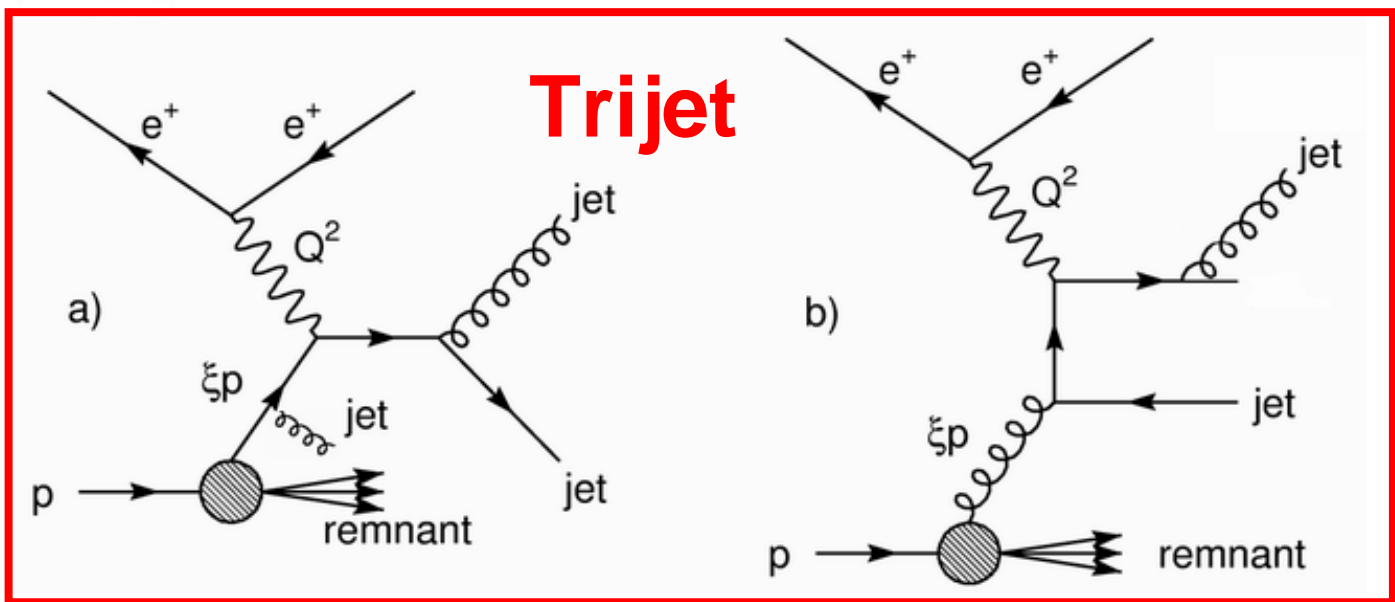
Dijet



b) Boson-Gluon Fusion



Trijet



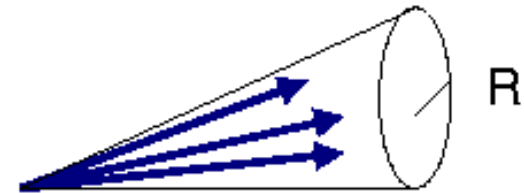
Why Trijet?

- Adding a gluon radiation to dijet
→ A direct test of QCD.
- Recent important advance in understanding QCD in dijets makes it an ideal laboratory for studying added gluon radiation.
- Measure α_s at a wide range of Q^2
- In the the ratio of $R_{3/2} = \sigma_{\text{trijet}}/\sigma_{\text{dijet}} = O(\alpha_s)$,
there is a cancellation of some experimental and theoretical uncertainties

Jet Algorithms

Cone algorithms

- conceptually simple
- theoretical/implementation issues
 - seed requirements
 - infrared unsafe at NNLO

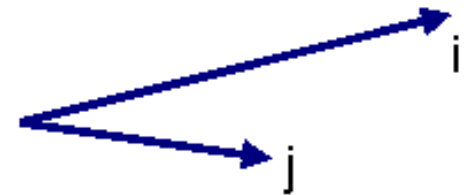


Maximize E_T within
a cone of radius R

Inclusive-mode k_T -algorithm (KTCLUS)

- No known theoretical or implementation issues
 - infrared safe
 - seeding not necessary
- Smaller hadronization corrections in some regions

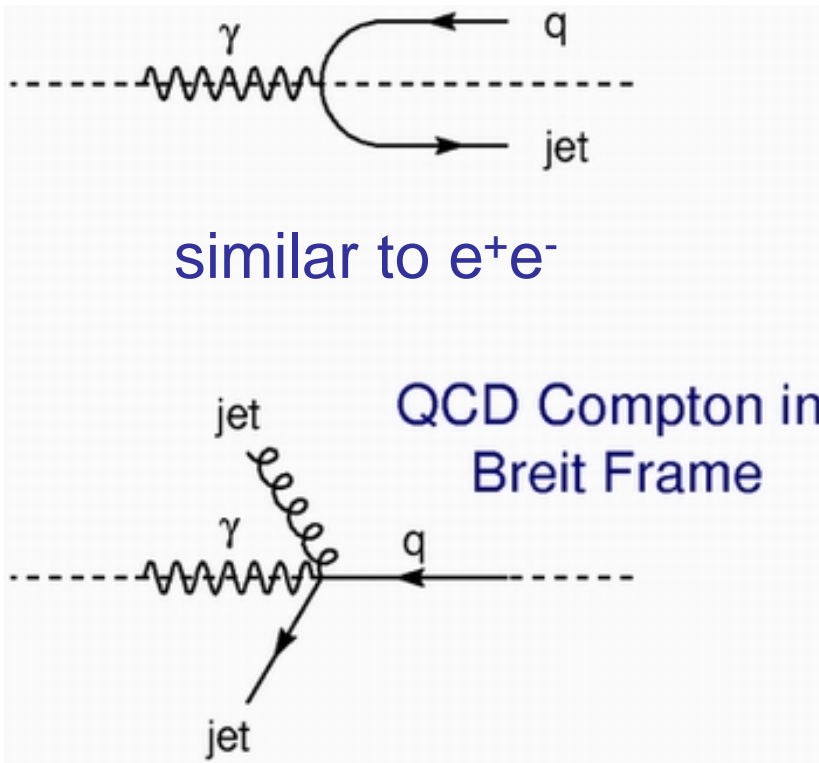
$$d_i = E_{T,i}^2$$
$$d_{ij} = \min\{E_{T,i}^2, E_{T,j}^2\}(\Delta\eta^2 + \Delta\phi^2)/R^2$$



Combine i and j if d_{ij}
is smallest of $\{d_i, d_{ij}\}$

Breit Frame

Single jet event in Breit Frame



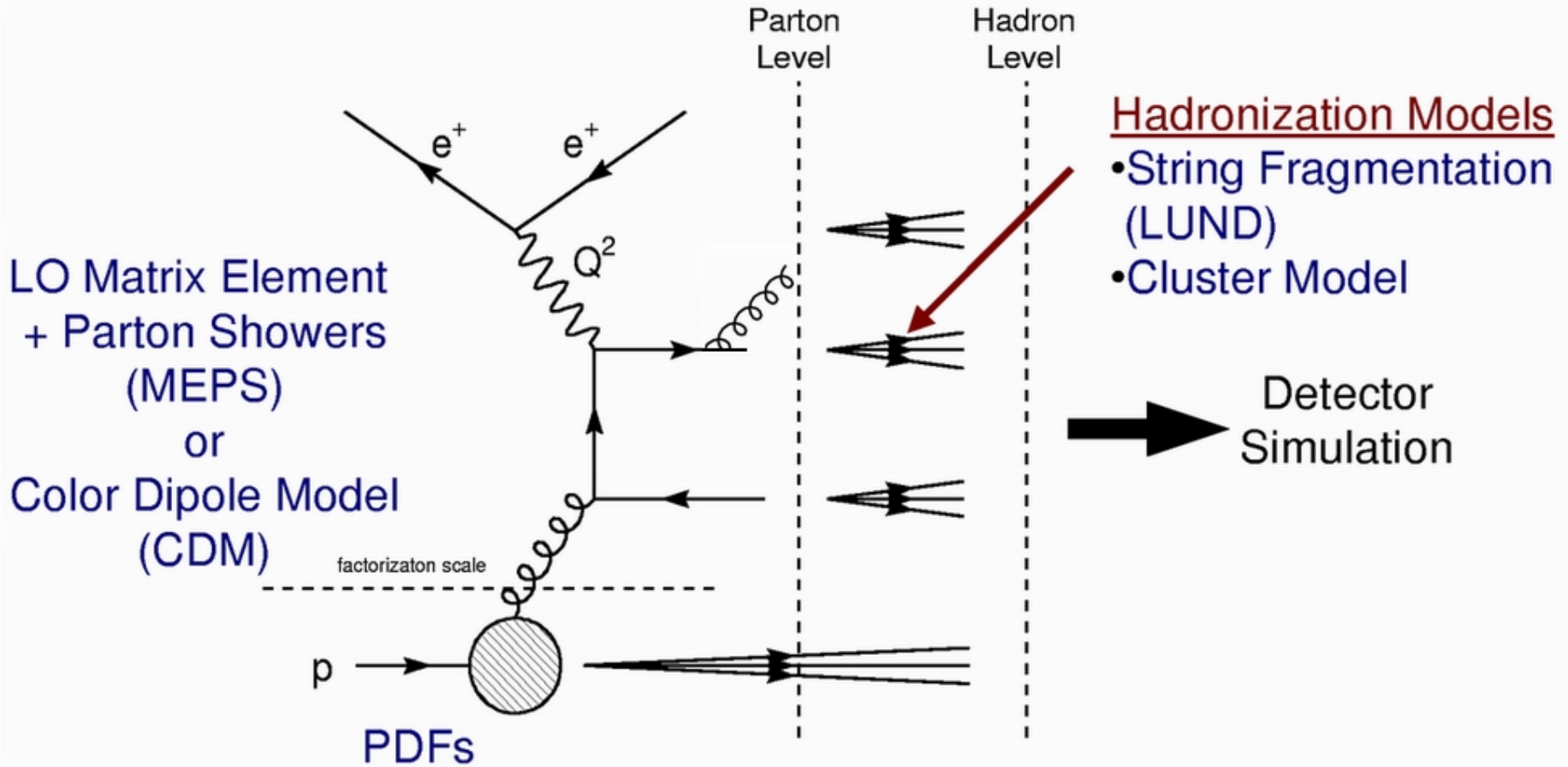
$$\mathbf{q} + 2\mathbf{p} = 0$$

In single jet events, struck quark rebounds with equal and opposite momentum, the resulting jet has zero E_T (transverse energy)

In multi-jet events, the outgoing jets are balanced in E_T

The proton and exchanged photon collide on a common axis, with the z-direction chosen to be the proton direction.

Leading Order MC



Hadronization Models

- String Fragmentation (LUND)
- Cluster Model

Detector Simulation

Choose

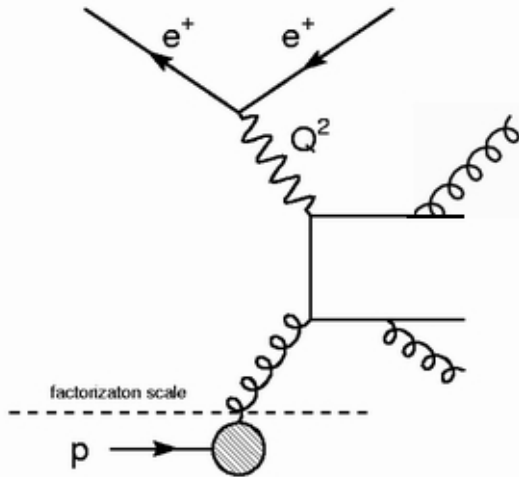
Programs

- LEPTO (MEPS+LUND)
- ARIADNE (CDM+LUND)
- HERWIG (MEPS+CLUSTER)

LO models used only for

- detector corrections
- hadronization corrections

Next Leading Order MC

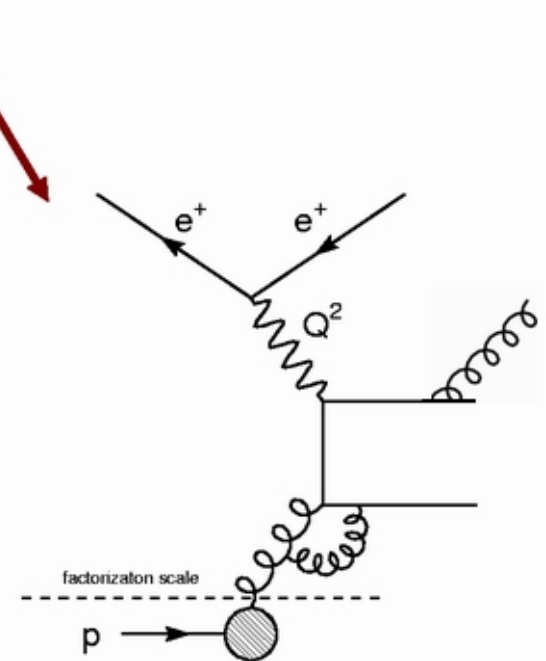


NLO Matrix Elements

- Four parton final states
- Large improvement over LO
- Soft/collinear and virtual loop divergences cancel

Programs for DIS

- DISINT (subtraction method)
- DISASTER++ (subtraction method)
- MEPJET (phase space splicing method)
- JETVIP (phase space splicing method)



NLO calculations do not include hadronization models

MC Scales

Factorization scale

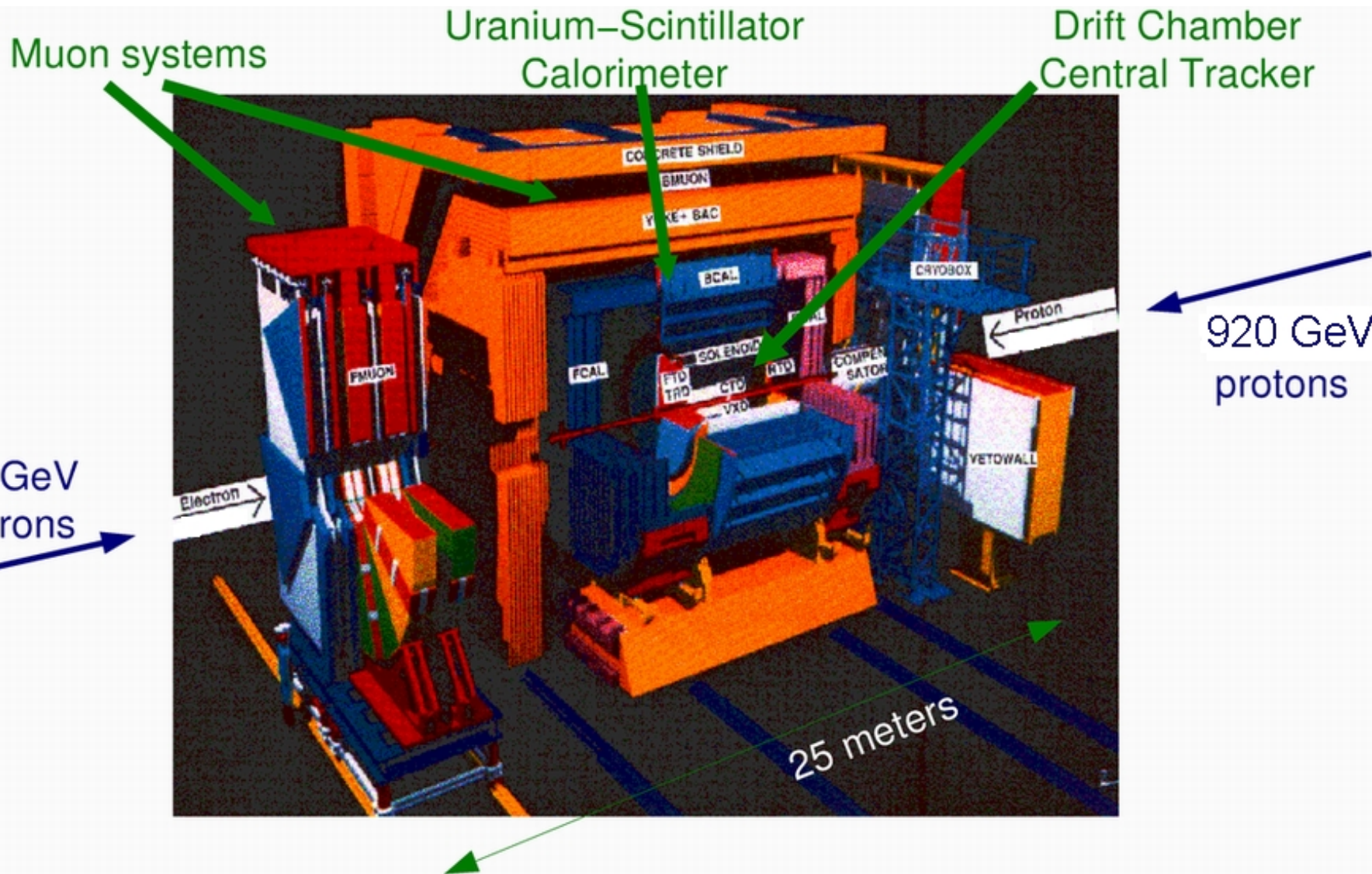
$f_q(x, \mu_f^2)$, scale μ_f at which the parton densities are evaluated and where the hadronization begins

Renormalization scale

$\alpha_s(\mu_r^2)$, scale μ_r at which the constant α_s is evaluated

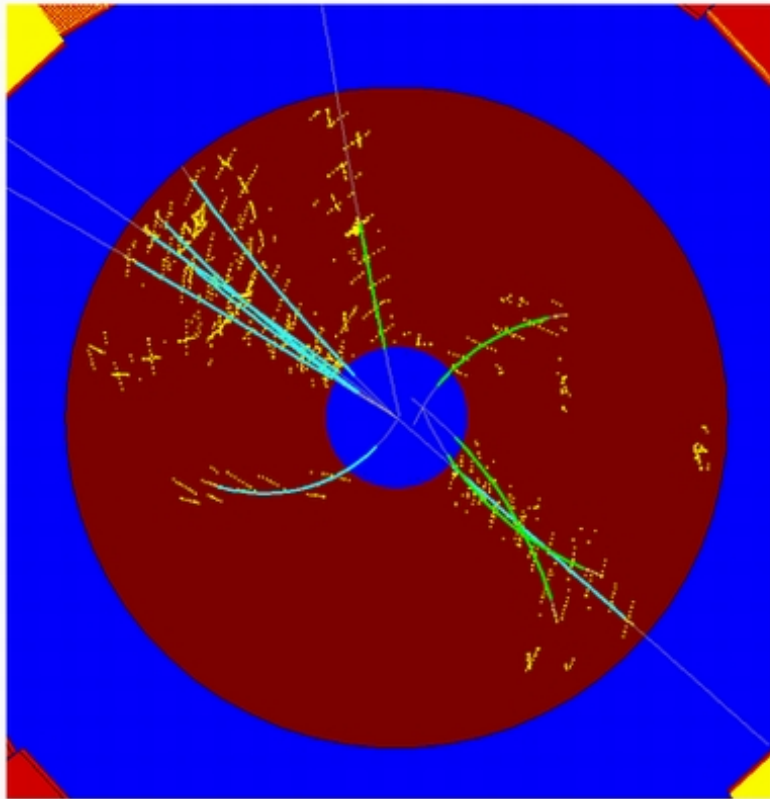
- NLO reduces renormalization scale dependence with respect to LO MC
- Renormalization scale uncertainty is the largest contribution to NLO theoretical uncertainty

ZEUS Detector



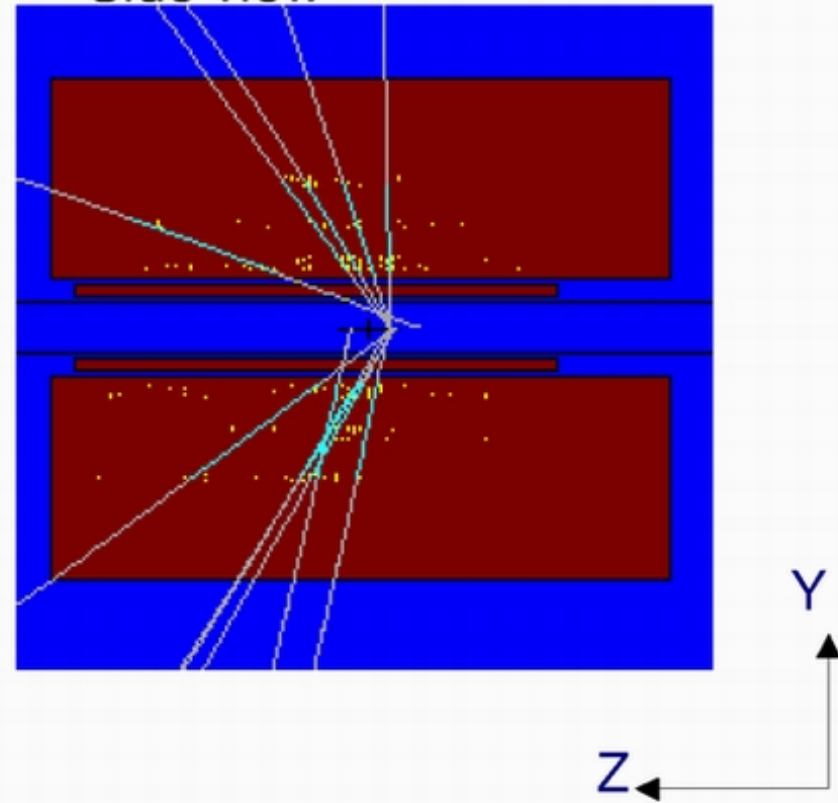
Central Tracking Detector

View along beamline



Drift Chamber
inside 1.43T solenoid

Side view

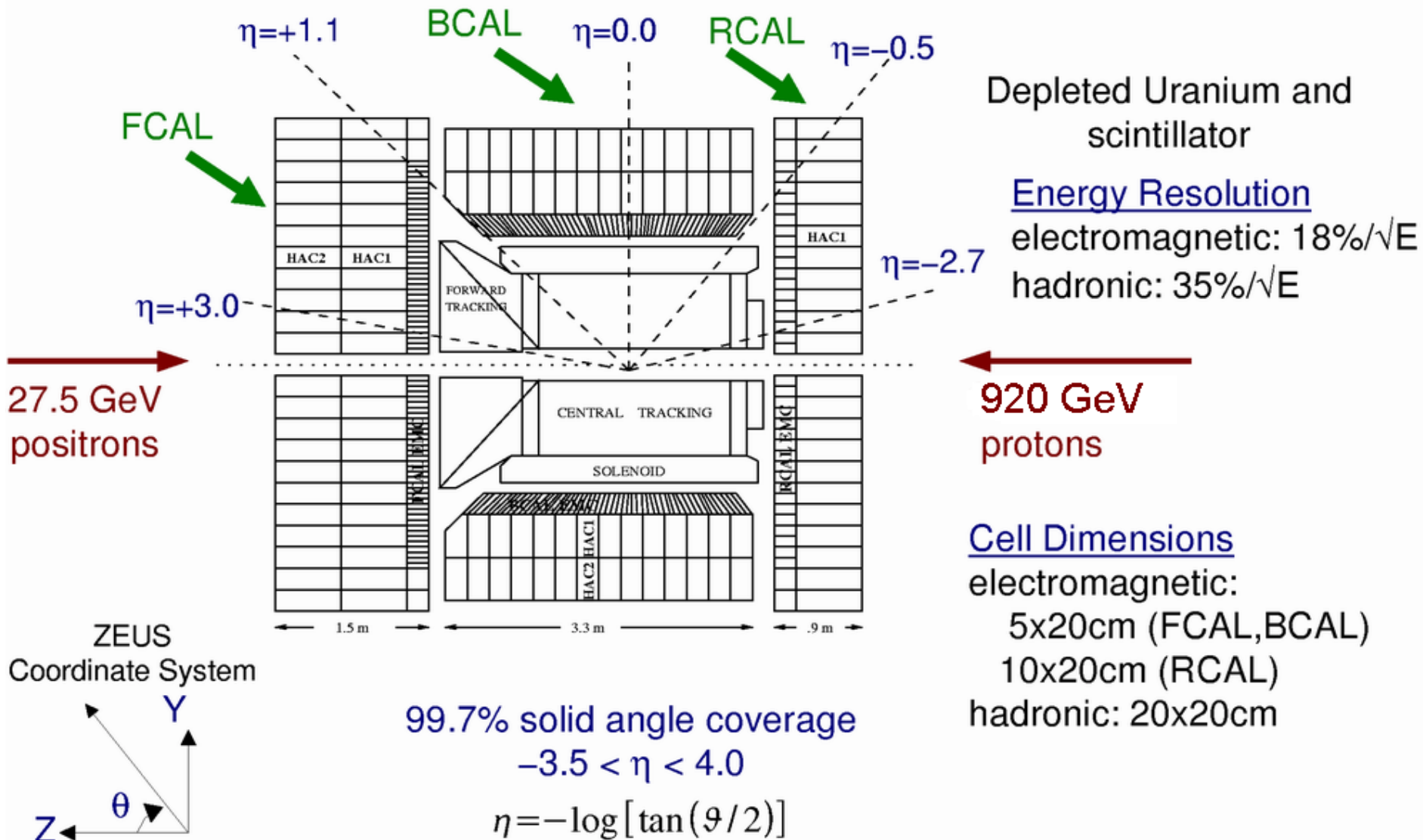


Vertex Resolution

longitudinal (z): 4mm

transverse (x-y): 1mm

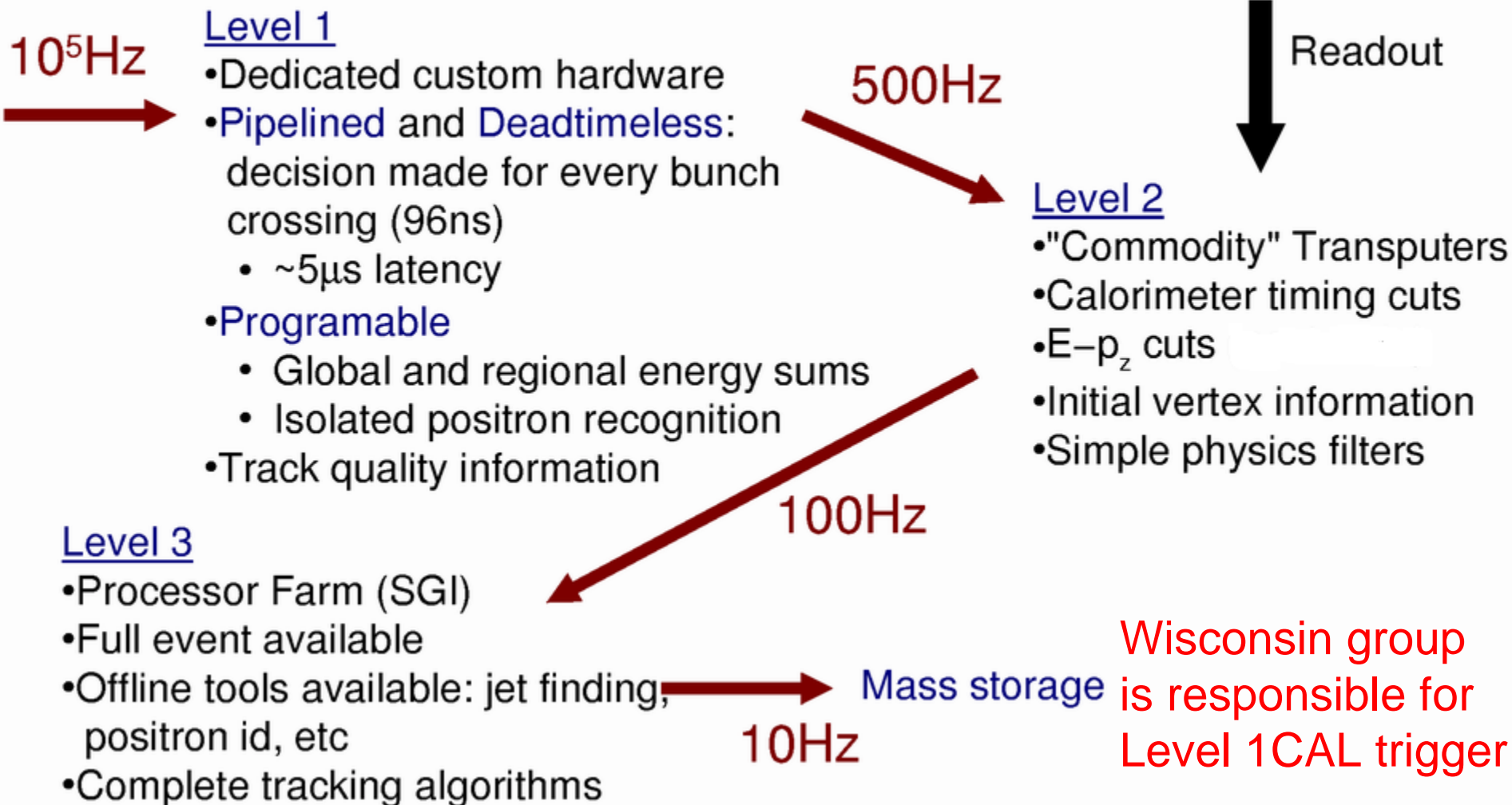
ZEUS Uranium Calorimeter



US groups, including Wisconsin are responsible for barrel calorimeter

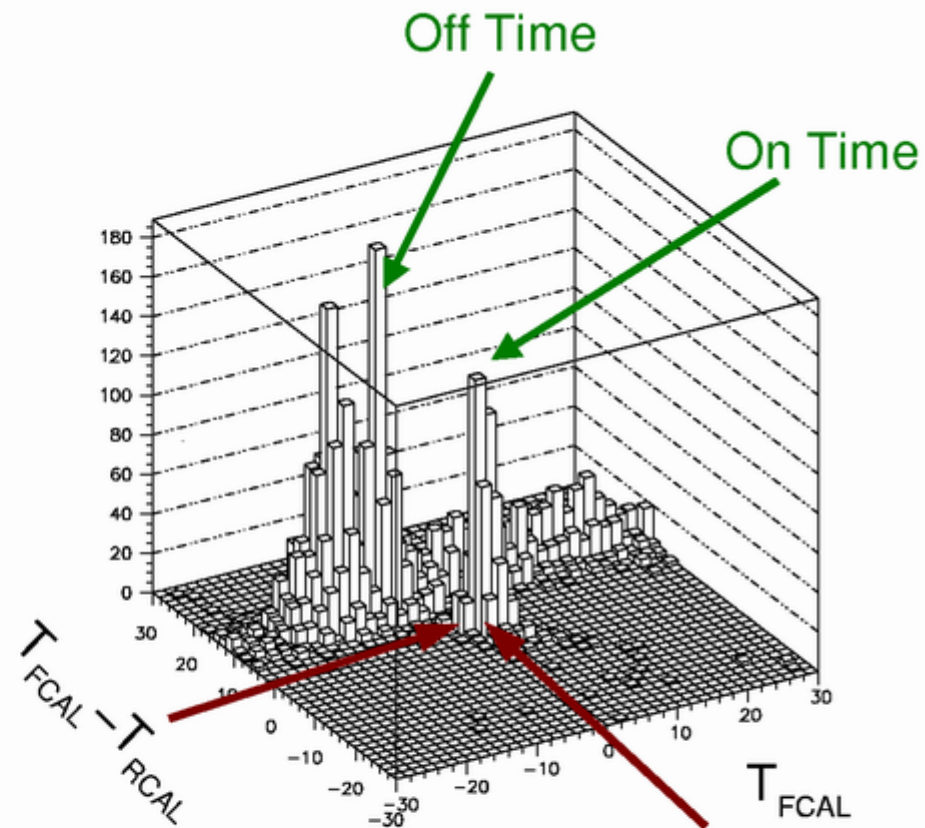
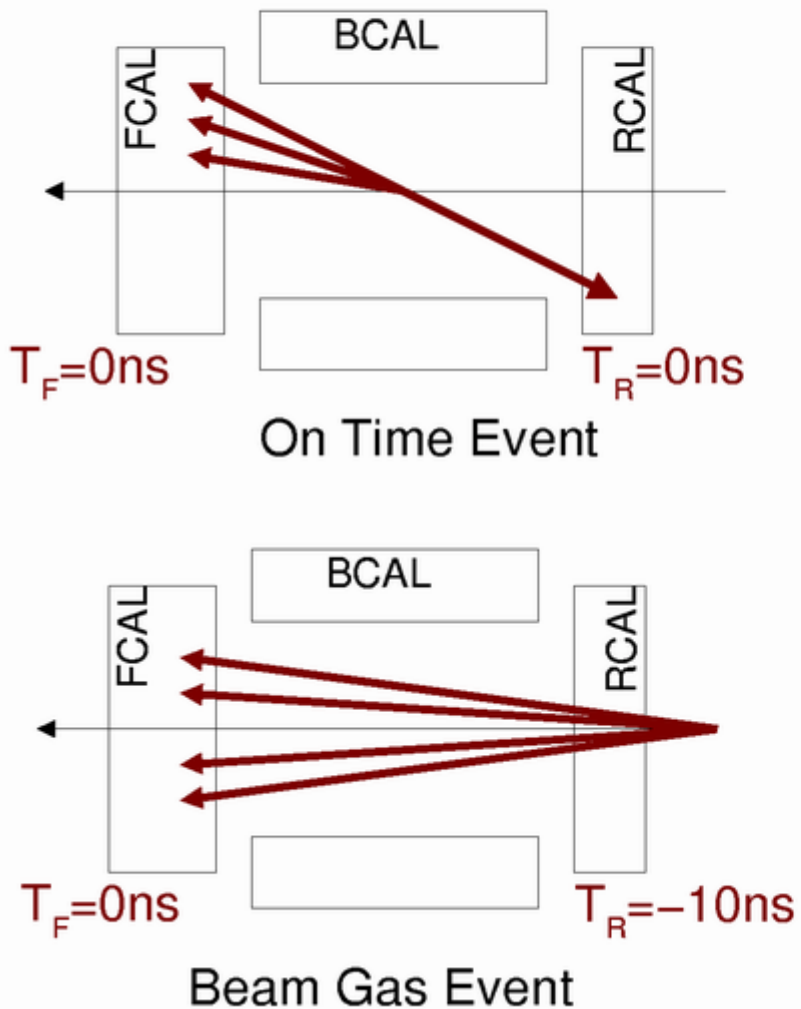
ZEUS Trigger

10 MHz bunch crossing rate
Extract 10Hz Physics from 100kHz background



Background Rejection: Timing

"Distance" between FCAL and RCAL is $\sim 10\text{ns}$



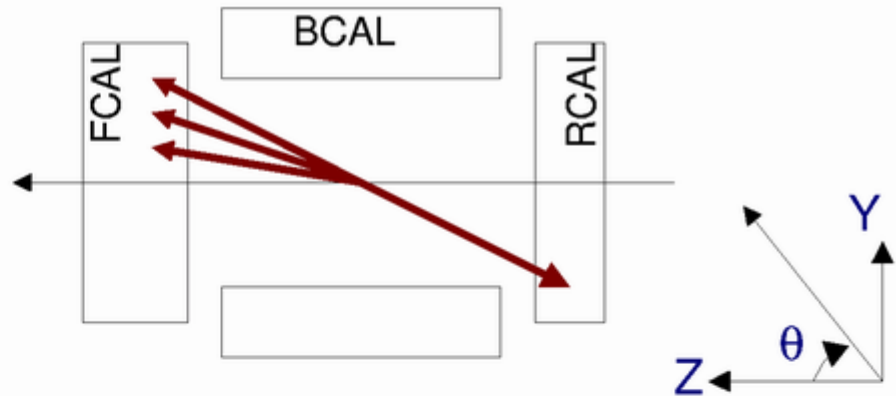
Calorimeter timing at Level 2

ZEUS Calorimeter timing resolution $< 1\text{ns}$

Background Rejection: E-P_z

$$E - p_z = \sum_i E_i (1 - \cos \theta_i)$$

Sum runs over calorimeter cells

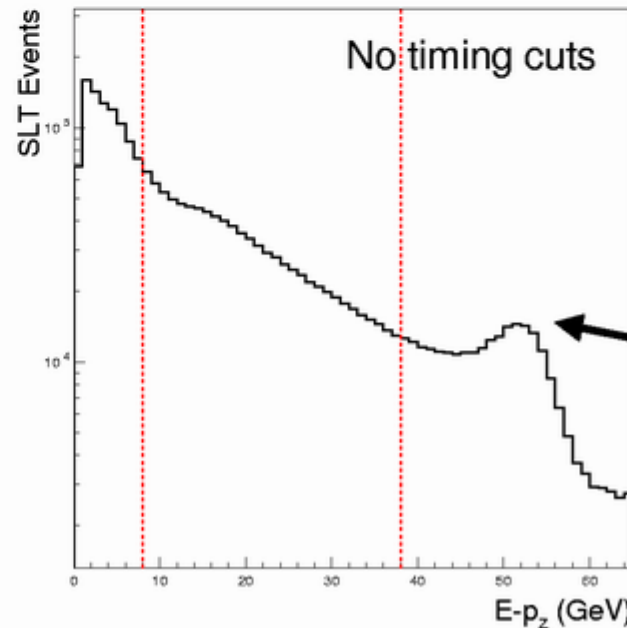


In a given frame, E-p_z is conserved



Before: $E - p_z = 2E_{\text{beam}} = 55\text{GeV}$

Unless energy escapes down rear beam pipe, E-p_z after collision will be near $2E_{\text{beam}}$ for NC DIS



Contained
NC DIS
 $E - p_z \approx 55$

Event Reconstruction

- Track finding and the event vertex
 - Require good track.
 - Use good track to find event vertex
- Electron finding
 - Find electrons by looking for isolated EM energy deposits in Calorimeter cells.
 - The ZEUS primary electron finder is $> 95\%$ efficient for electron energy > 10 GeV.

Kinematic Reconstruction

Two Kinematic Variables: x Q^2

Four Measured Quantities: E_e' θ_e E_h γ_h

Electron Method

Use scattered electron energy, electron angle

Good resolution in x and Q^2 , best at low Q^2

Sensitive to miscalibrations (energy scale uncertainties)

Double Angle Method

Use electron angle and hadronic jet angle

Depends only on ratios of energies

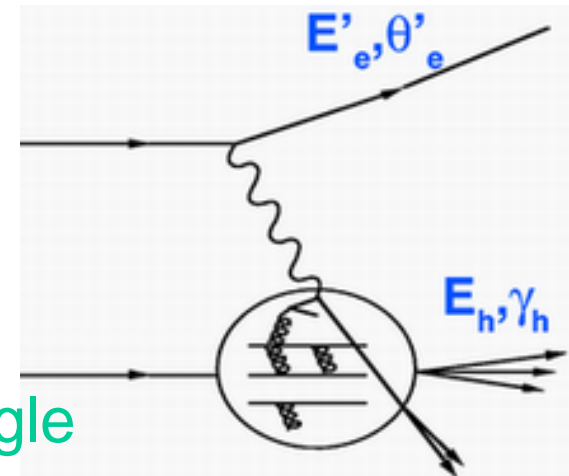
Better mean resolution in x and Q^2

Weakly affected by miscalibrations

Jacquet-Blondel Method

Use hadronic energy and hadronic jet angle

Gives an accurate determination of y for small values of y



Data Offline Cuts

- **Detector Acceptance and Efficiency**

$y_{JB} > 0.04$

positron found with $E > 10$ GeV

vertex cut: $-50\text{cm} < z < 50\text{cm}$

positron position cut: $|X| > 14$ cm or $|Y| > 14$ cm

- **Signal Selection**

$38 \text{ GeV} < E - P_z < 65 \text{ GeV}$

A well found track with $p > 5$ GeV & DCA < 15 cm

$y_{EL} < 0.95$

- **Physics and Kinematic Requirement**

$100 \text{ GeV}^2 < Q^2 < 10000 \text{ GeV}^2$

Three jets found using KTCLUS algorithm

$|\eta| < 2.0$

$E_T^{\text{LAB}} > 5 \text{ GeV}$

Asymmetric cut, MC calculation requirement

$E_{T,1}^{\text{BRT}} > 8 \text{ GeV}, E_{T,2}^{\text{BRT}} > 5 \text{ GeV}, E_{T,3}^{\text{BRT}} > 5 \text{ GeV}$

First Comparison of Data and MC DIS Trijets at ZEUS

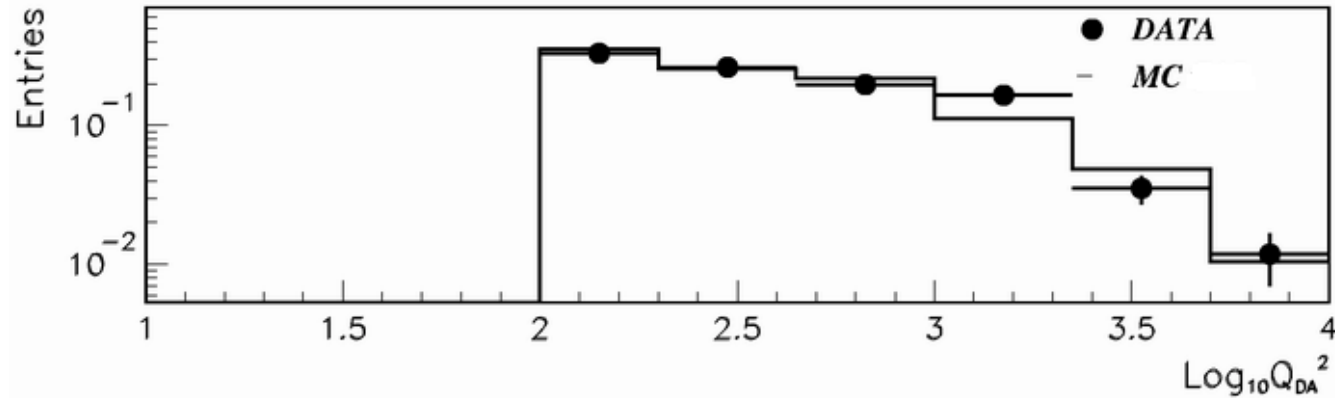
Data:

- HERA 1996-1997 running period
- ZEUS data sample: integrated luminosity 38.4pb^{-1}
- 510 DIS trijets passed offline cuts in a sample of 39576 DIS events

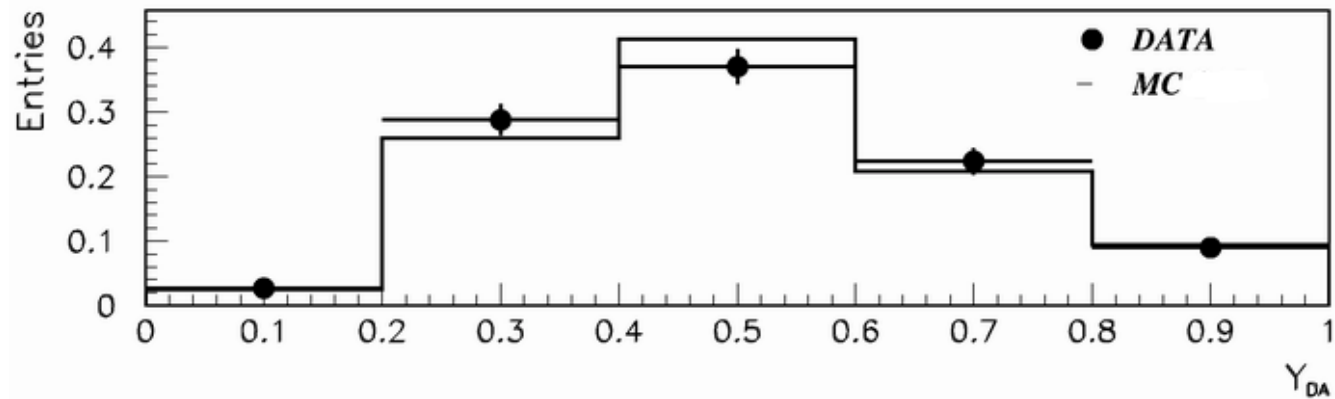
MC:

- Use leading order MC program **Ariadne**
- 568 DIS trijets selected from MC
- No Adjustment, run with default ZEUS settings

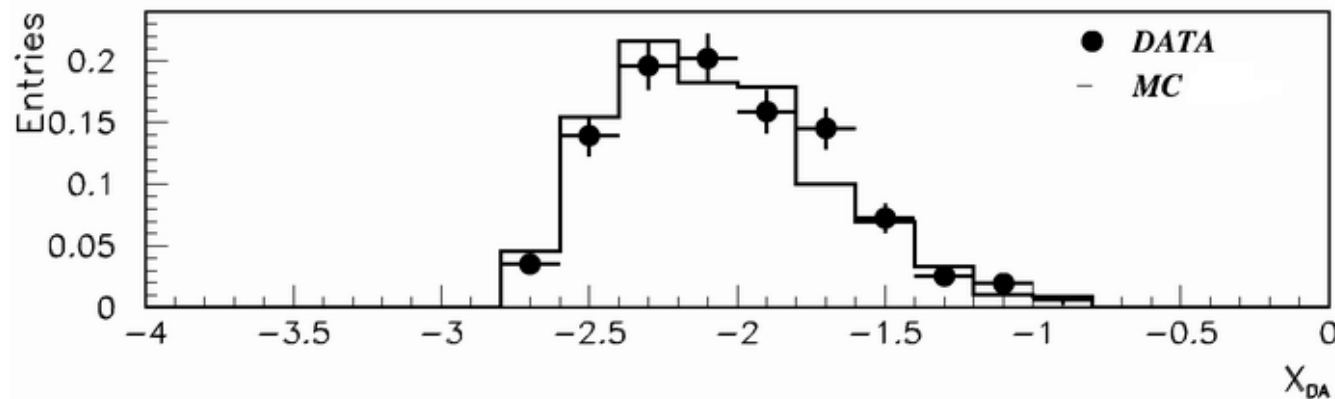
DIS Trijets — A First Look



Reconstructed
Q², y, x by
Double Angle
method



Use Ariadne
for all MC



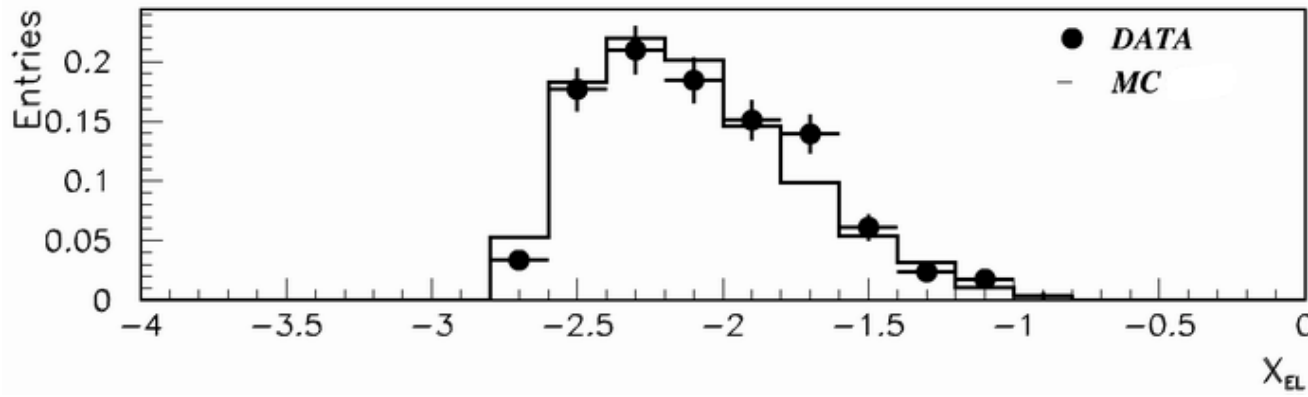
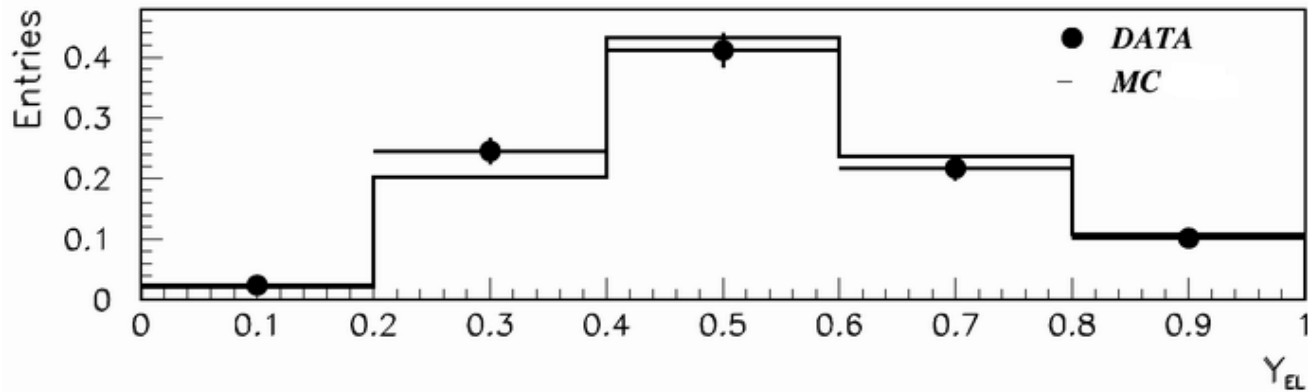
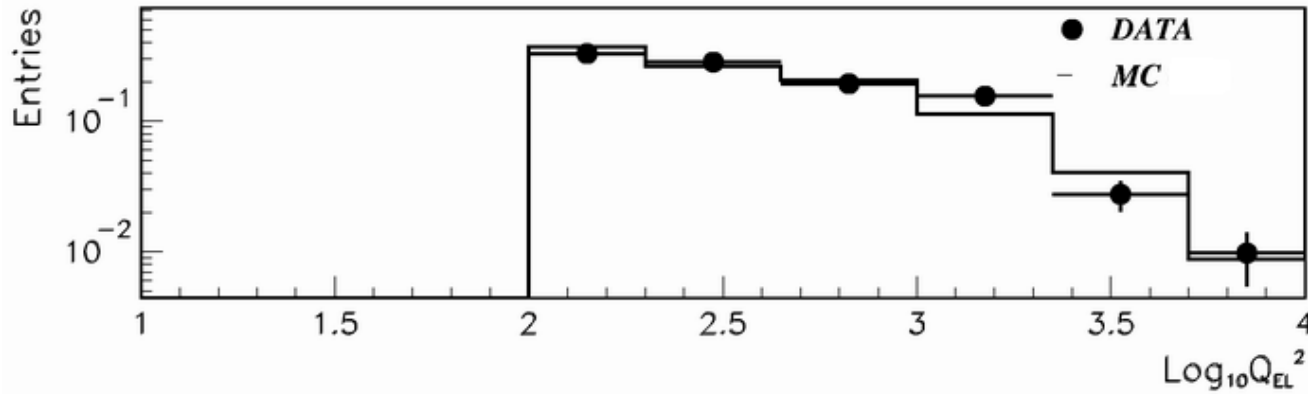
OK agreement
between
data and MC

DIS Trijets — A First Look

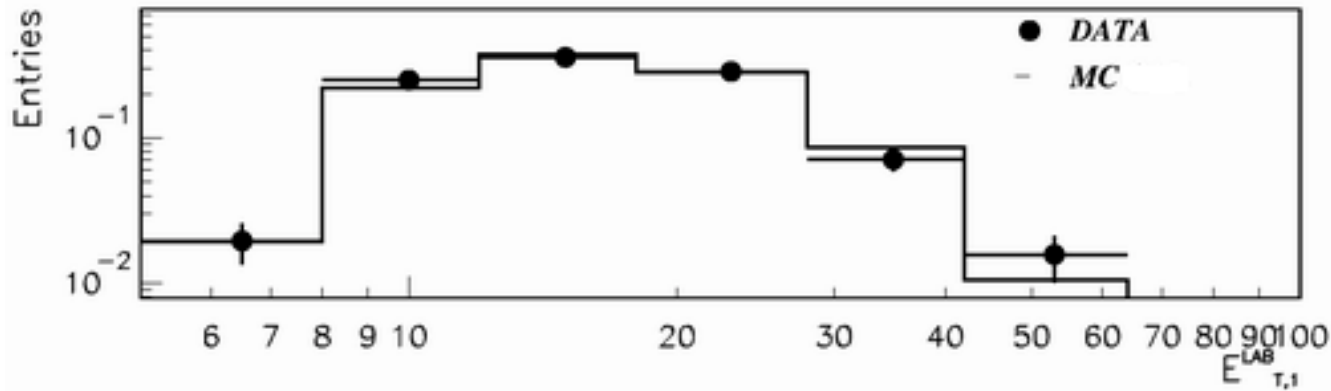
Reconstructed
 Q^2 , y , x by
Electron method

Better agreement
between
data and MC

→ Pick Electron
method

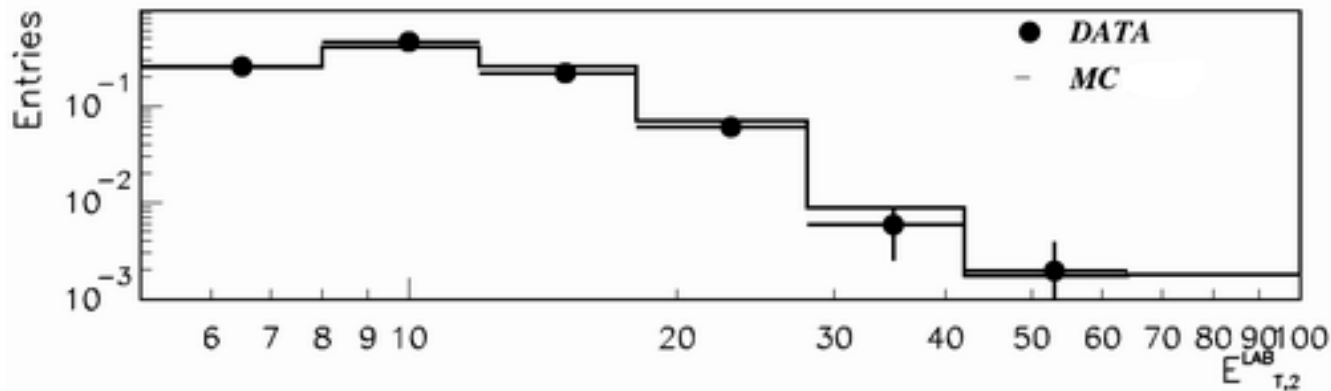


DIS jet E_T in the Lab Frame

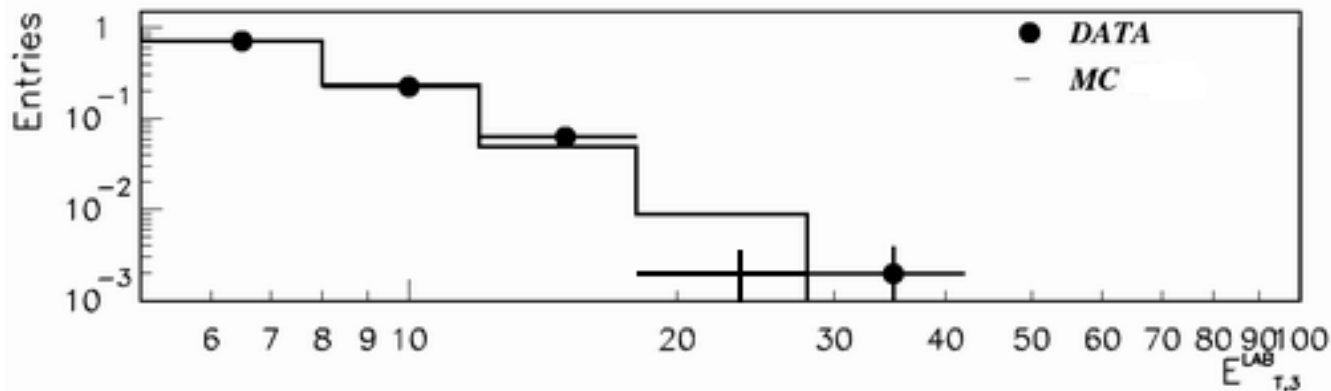


E_T Ordered:

First Jet



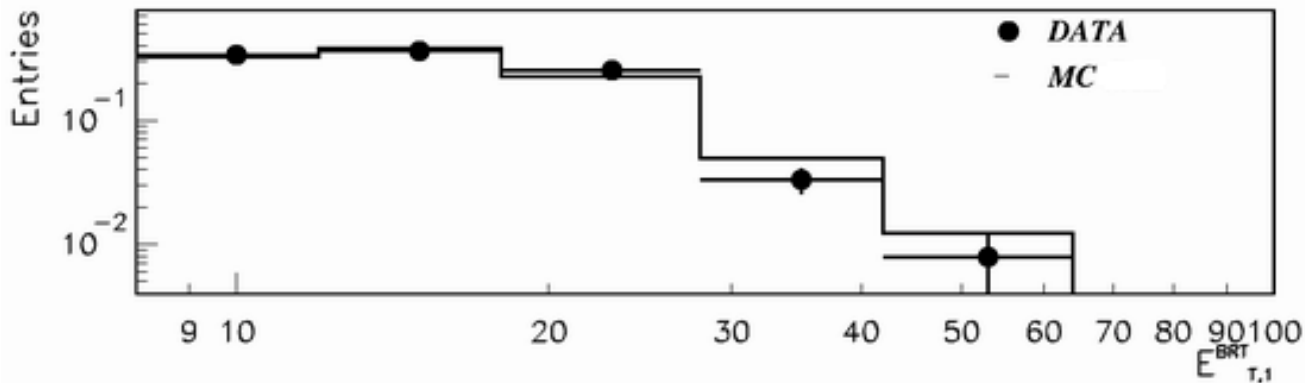
Second Jet



Third Jet

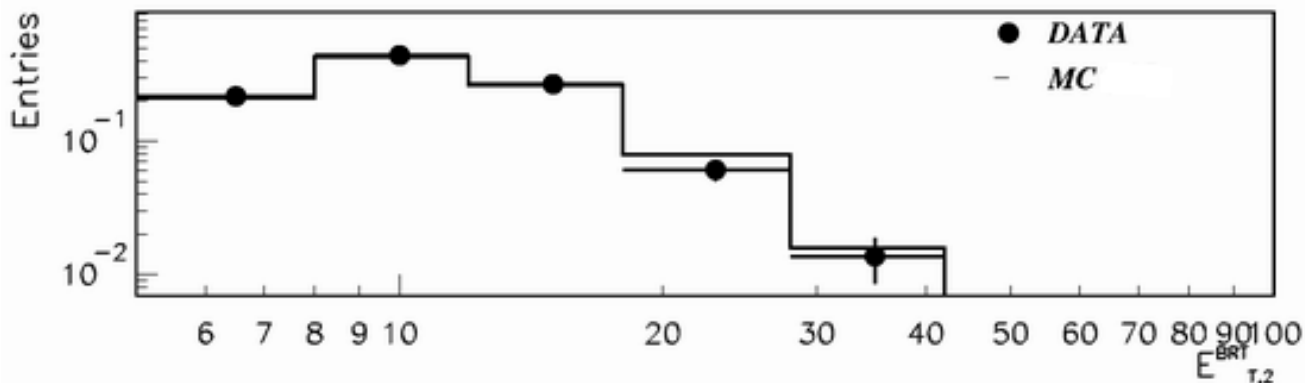
Reasonable agreement

DIS jet E_T in the Breit Frame

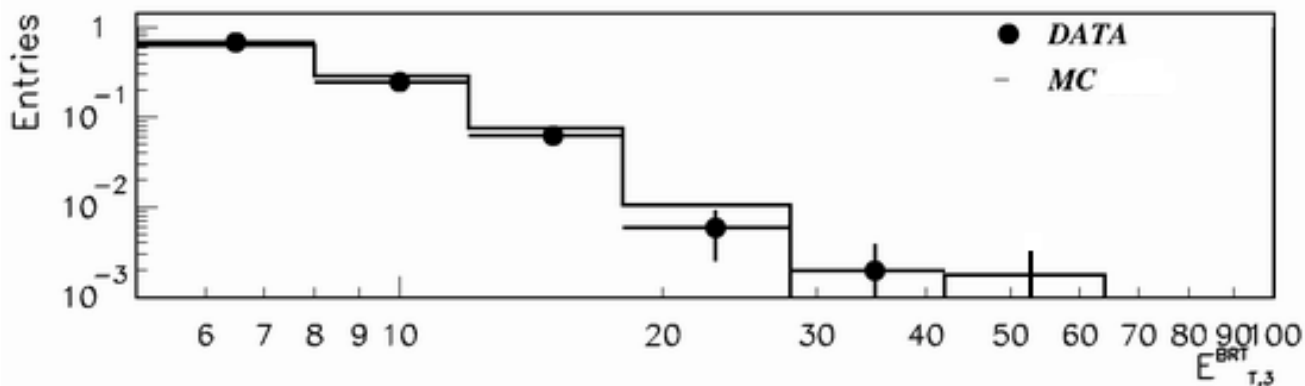


E_T Ordered:

First Jet



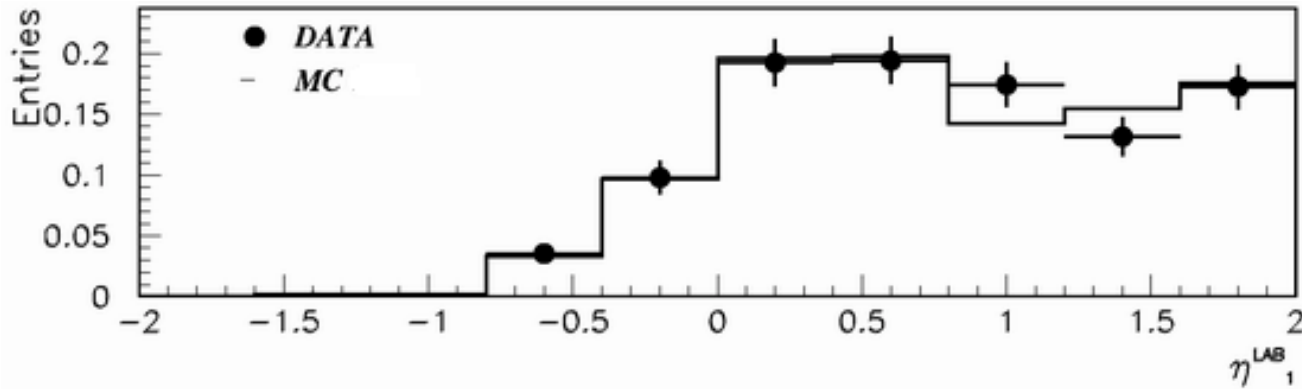
Second Jet



Third Jet

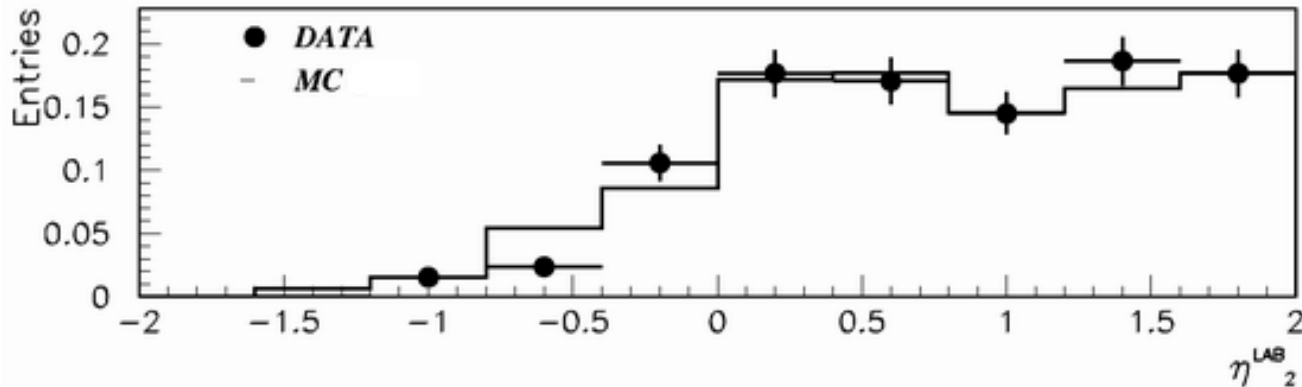
Reasonable agreement

DIS jet η in the Lab Frame

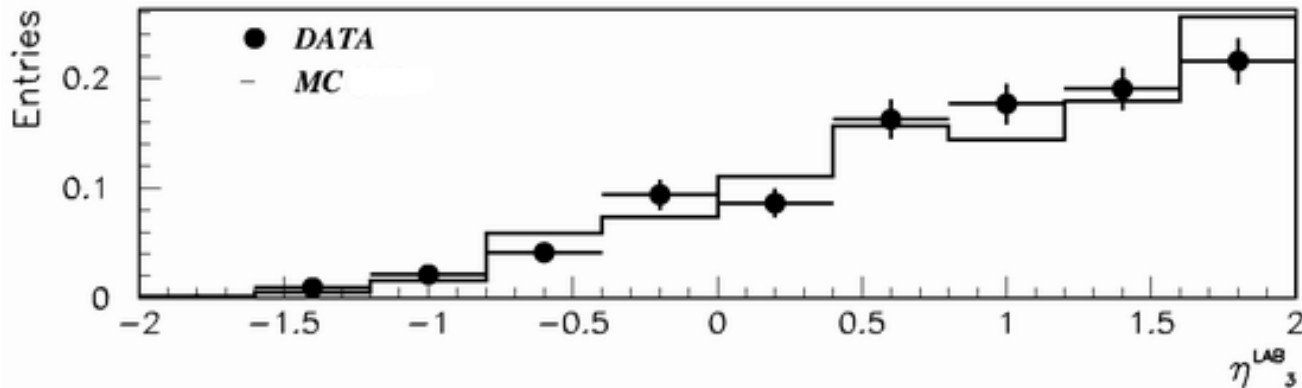


E_T Ordered:

First Jet



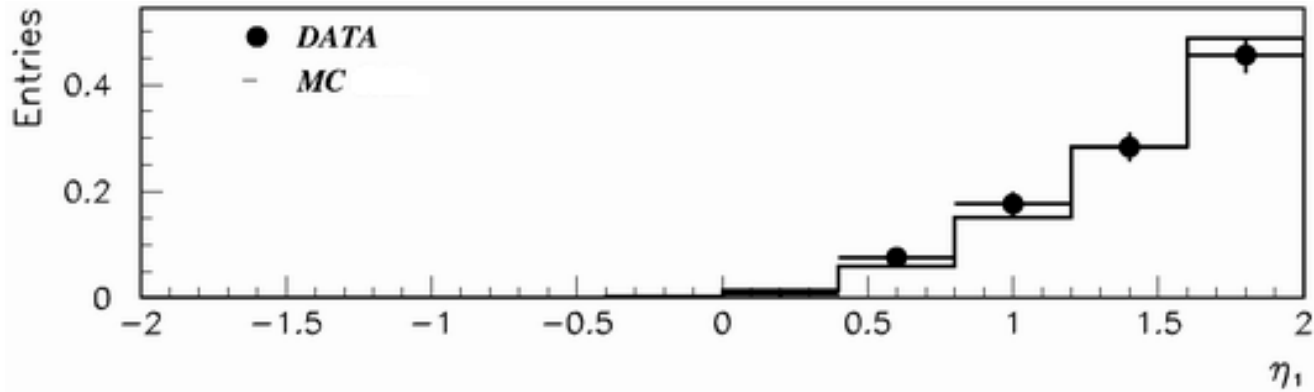
Second Jet



Third Jet

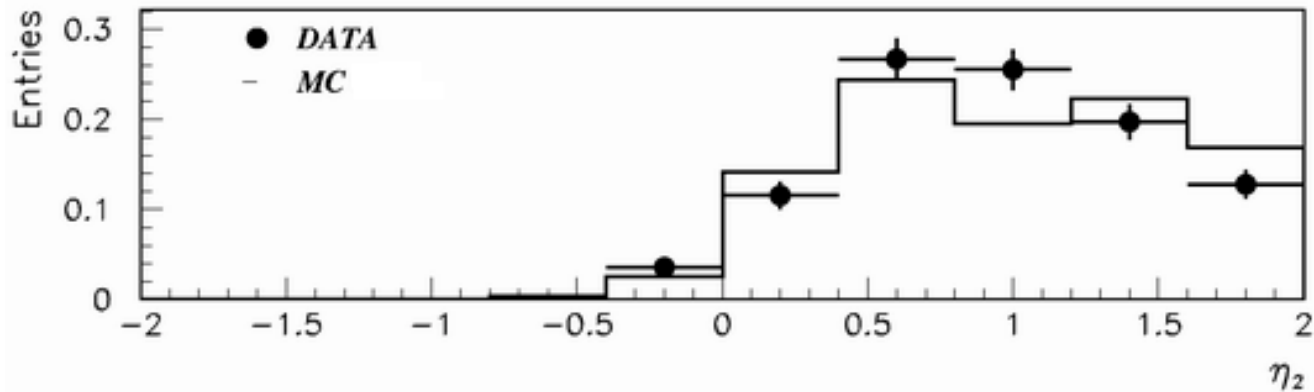
Reasonable agreement

DIS jet η in the Lab Frame

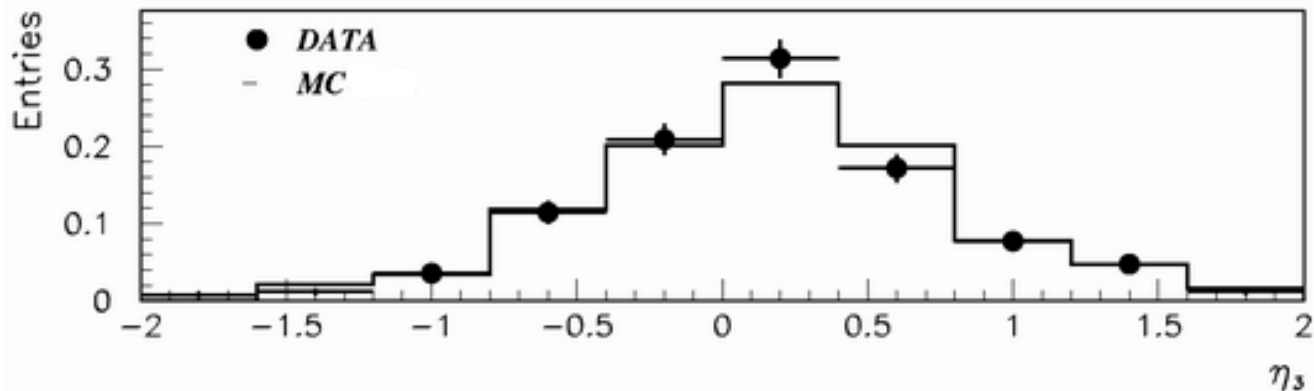


η Ordered:

First Jet



Second Jet



Third Jet

Reasonable agreement

Conclusions and Expectations

- First look at DIS Trijets at ZEUS, reasonable agreement
- Ariadne used as it is, a good starting point
- Need to calculate cross sections and compare in detail with other QCD calculations and different MC programs
- Add new data (99-00) to have more statistics
- Explore systematic uncertainties