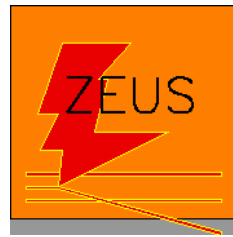

Measurement of the Proton F_2 Structure Function from 1996 and 1997 ZEUS data

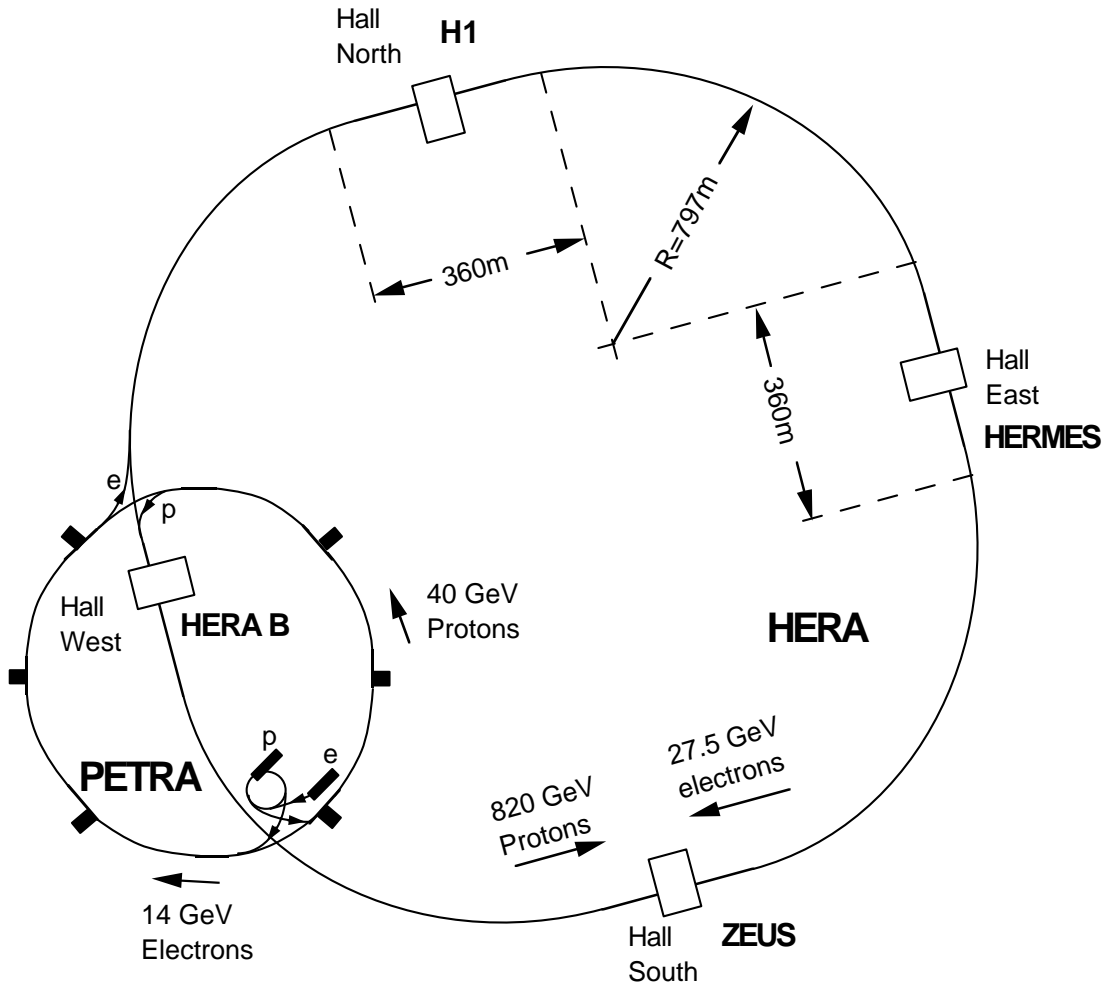
Mike Wodarczyk

November 30, 1999



- HERA and ZEUS
- Kinematic Variables and Structure Functions
- F_2 History
- ZEUS Events
- F_2 Analysis Challenges
 - Electron Energy Angle
 - Vertex
 - Hadronic System
- F_2 Binning & Unfolding
- F_2 Result
- Conclusions

HERA



- Fixed Target C.O.M. energy (Max ~35 GeV)

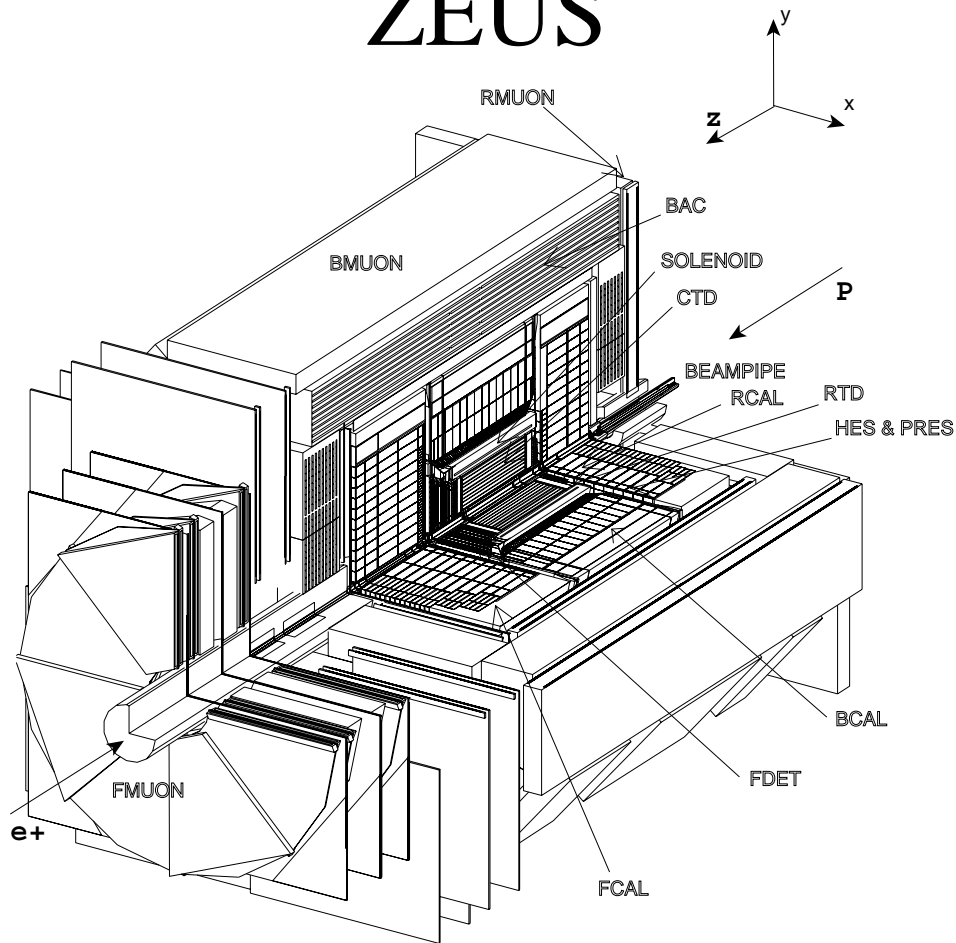
$$\sqrt{s} = \sqrt{2 * M_P * E_{beam}}$$

- HERA

$$\sqrt{s} = \sqrt{4 * E_{Beam} * P_{Beam}} = 300 \text{ GeV}$$

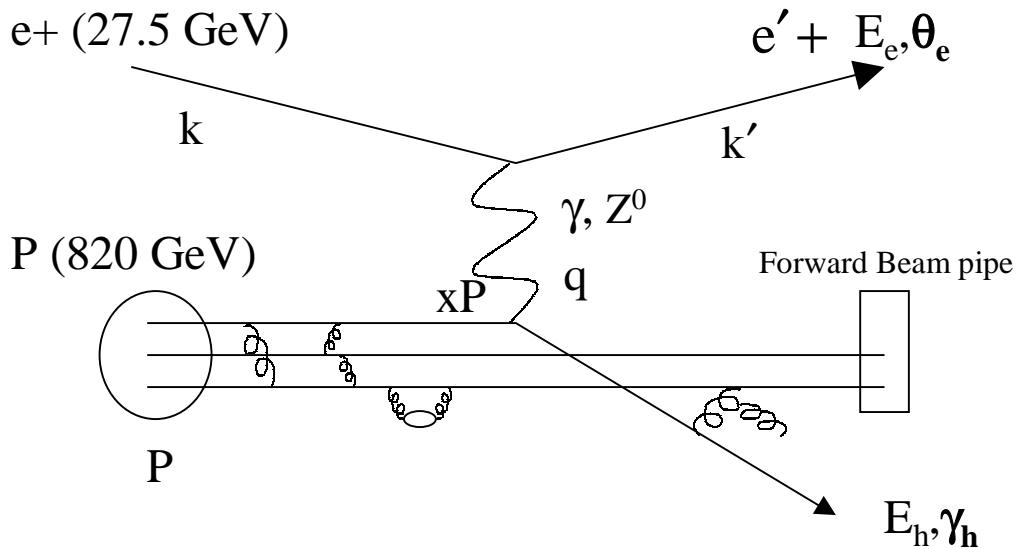
- Equivalent to a 48 TeV Fixed Target experiment
- Boosted C.O.M. in the proton direction

ZEUS



- Calorimeter
 - Rear, Barrel, Forward Calorimeters
 - Longitudinal segmentation (Electro-magnetic, Hadronic separation)
 - Transverse segmentation (Particle location)
- Central Tracking Detector (CTD)
 - Drift chamber (9 cylindrical layers, 8 sense wires per layer)
 - Inside a 1.4 T magnetic field
- Presampler / SRTD
 - Scintillator strips at the face of the rear calorimeter counting charged particles.

Kinematic Variables



$Q^2 = -(k - k')^2$ Inverse scale at which the proton is probed

$\lambda \approx \frac{2xM_P}{Q^2}$ Virtual photon wavelength λ

$x = \frac{Q^2}{2P \cdot q}$ Fraction of proton's momentum carried by struck quark

$y = \frac{P \cdot q}{P \cdot k}$ Fraction of electron's momentum transferred by the photon in the proton's rest frame

- Only 2 independent variables

$sxy = Q^2$ Center of mass energy $= \sqrt{s}$

Structure Functions

$$\frac{d^2\sigma(e^+p)}{dx dQ^2}(x, Q^2) = \frac{2\pi\alpha^2}{xQ^4} \left[(1 + (1-y)^2)F_2(x, Q^2) - y^2F_L(x, Q^2) - (1 - (1-y)^2)xF_3(x, Q^2) \right]$$

- e^+p cross section, σ , can be written in terms of unit-less structure functions F_2 , F_L , xF_3

- **Quark Parton Model**

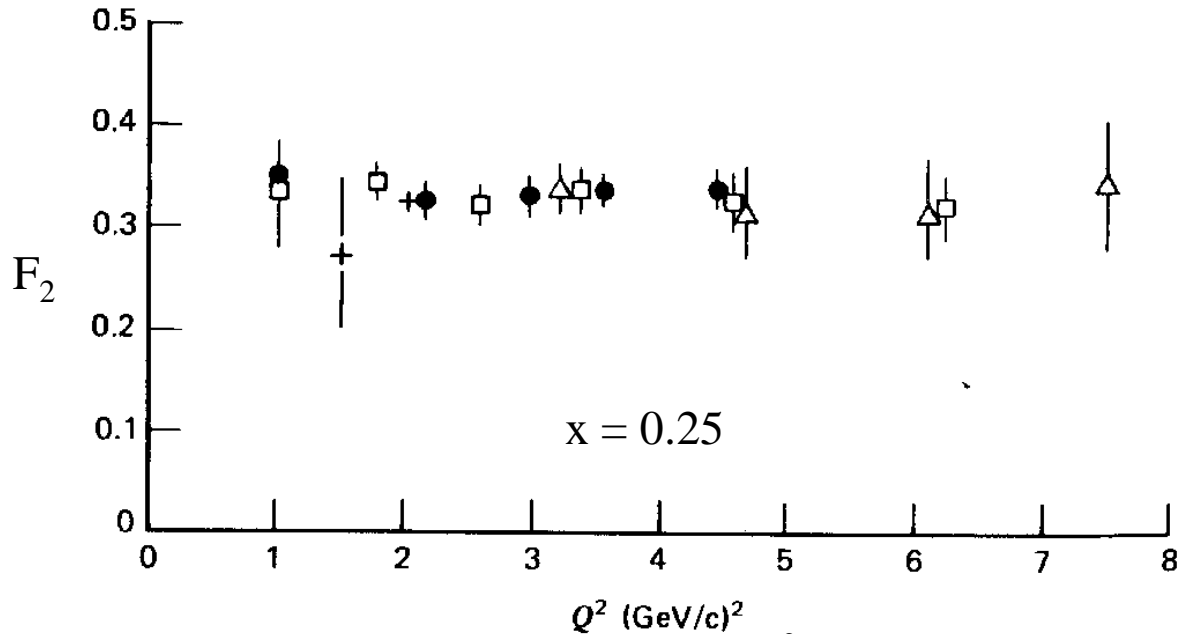
- Proton is made up of quarks (uud..)and gluons

$$F_2(x, Q^2) = \sum_q^{quarks} e_q^2 x [q(x, Q^2) + \bar{q}(x, Q^2)]$$

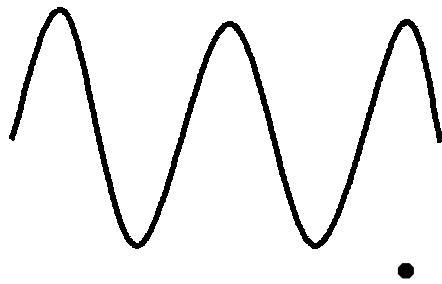
- F_2 is the contributions of photon-quark interactions.
- F_L is the contributions from photons with longitudinal polarization
 - F_L is 0 in leading order QCD theory
 - Contribution small at low y
- xF_3 is the electro-weak contribution from Z^0 exchange.
 - Contribution small for $Q^2 \ll M_Z^2 = 8100 \text{ GeV}^2$

Historical Results

- SLAC (1972) e beam on fixed target p



- Scaling; F_2 does not depend on Q^2
- Number of quarks does not depend on $Q^2=1/\lambda$



Low Q^2



High Q^2

- Quarks are point-like objects in the proton

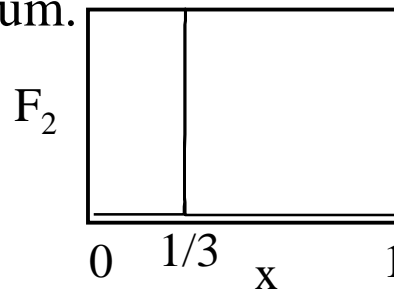
F₂

$$F_2(x, Q^2) = \sum_q^{quarks} e_q^2 x [q(x, Q^2) + \bar{q}(x, Q^2)]$$

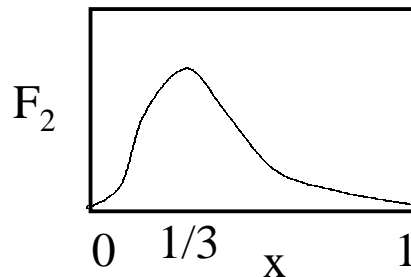
- What might one expect?

- Proton is only 3 quarks (uud) which equally share the protons momentum.

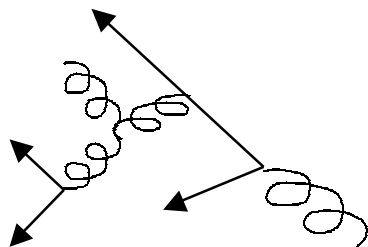
- F₂=0 except at x=1/3



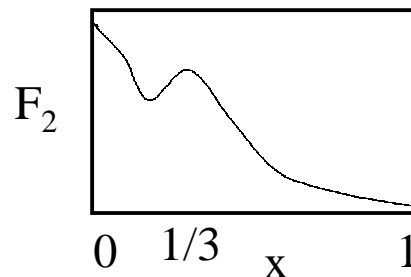
- 3 quarks exchange momentum via gluons but still carry the full protons momentum



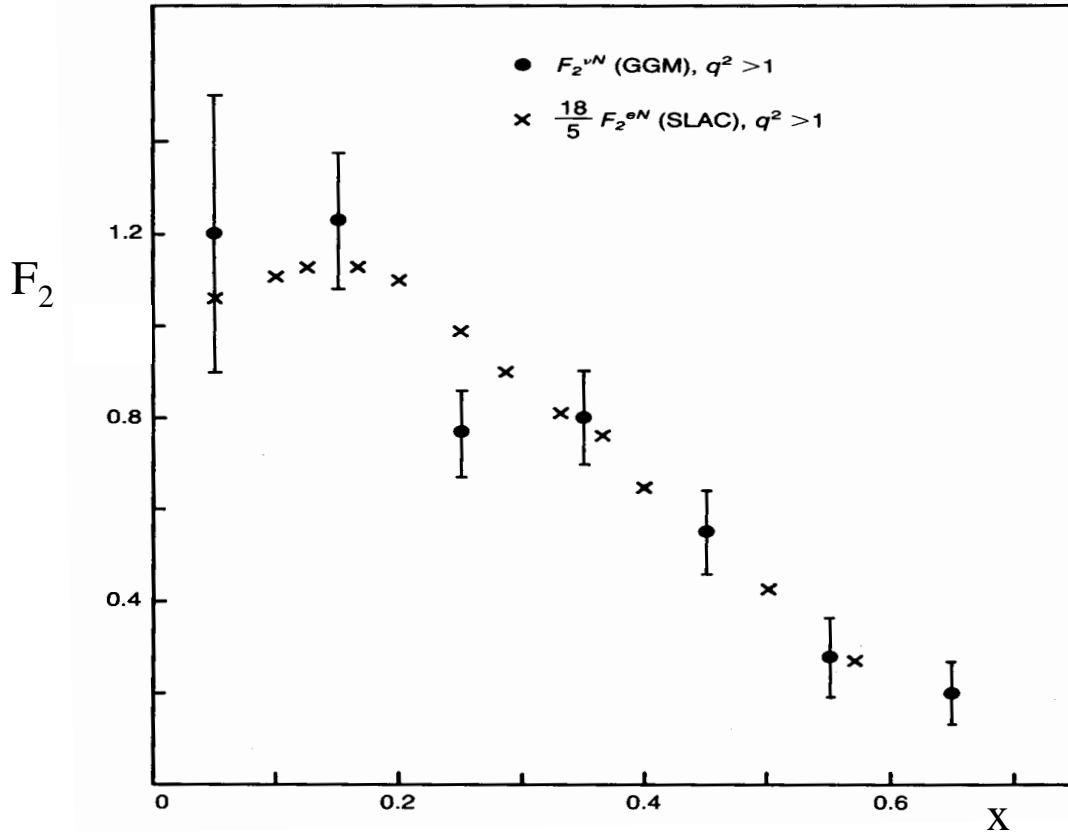
- 3 quarks + a sea of quarks from pair production from gluons



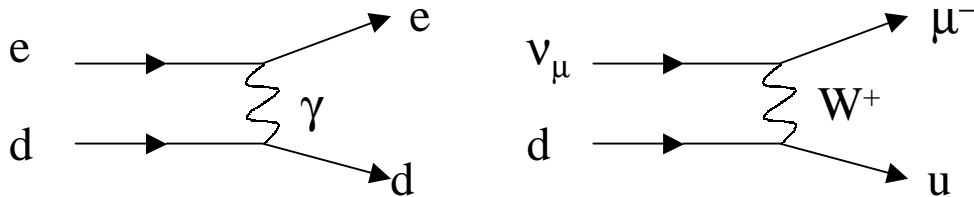
Decreasing x ←



CERN + SLAC



– Comparison of neutrino data and electron data



- Confirms fractional charges of quarks

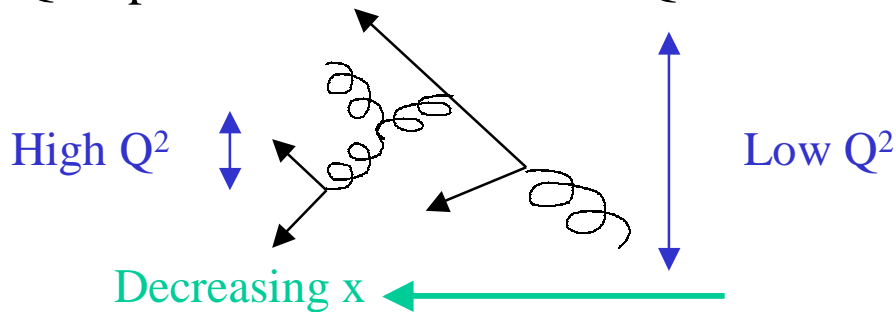
– Integral of quark momenta

$$\int_x F_2^{\nu N}(x) dx = \int_x x[u(x) + d(x) + \bar{u}(x) + \bar{d}(x) + \dots] dx \approx 0.5$$

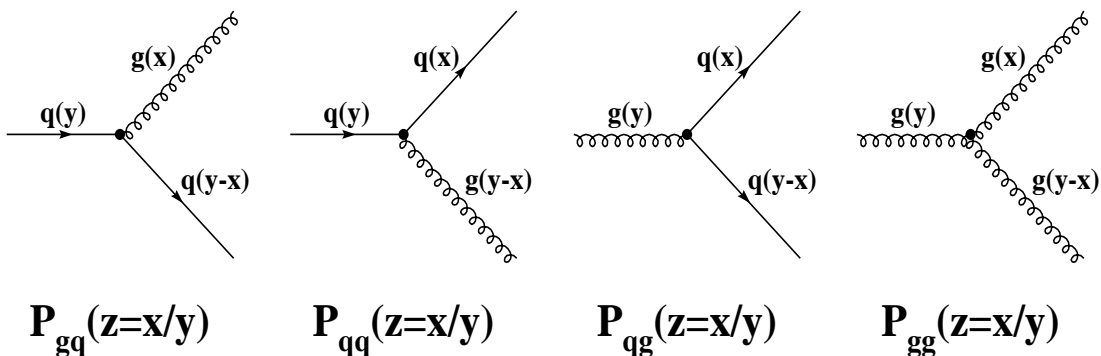
- Only half of proton's momentum!
- Remainder carried by gluons

DGLAP Evolution

- Dokshitzer-Gribov-Lipatov-Altarelli-Parisi
- Prediction for the Q^2 evolution of F_2
 - From parton distribution functions at an initial Q^2 it predicts the PDFs at all Q^2



- Uses QCD splitting functions

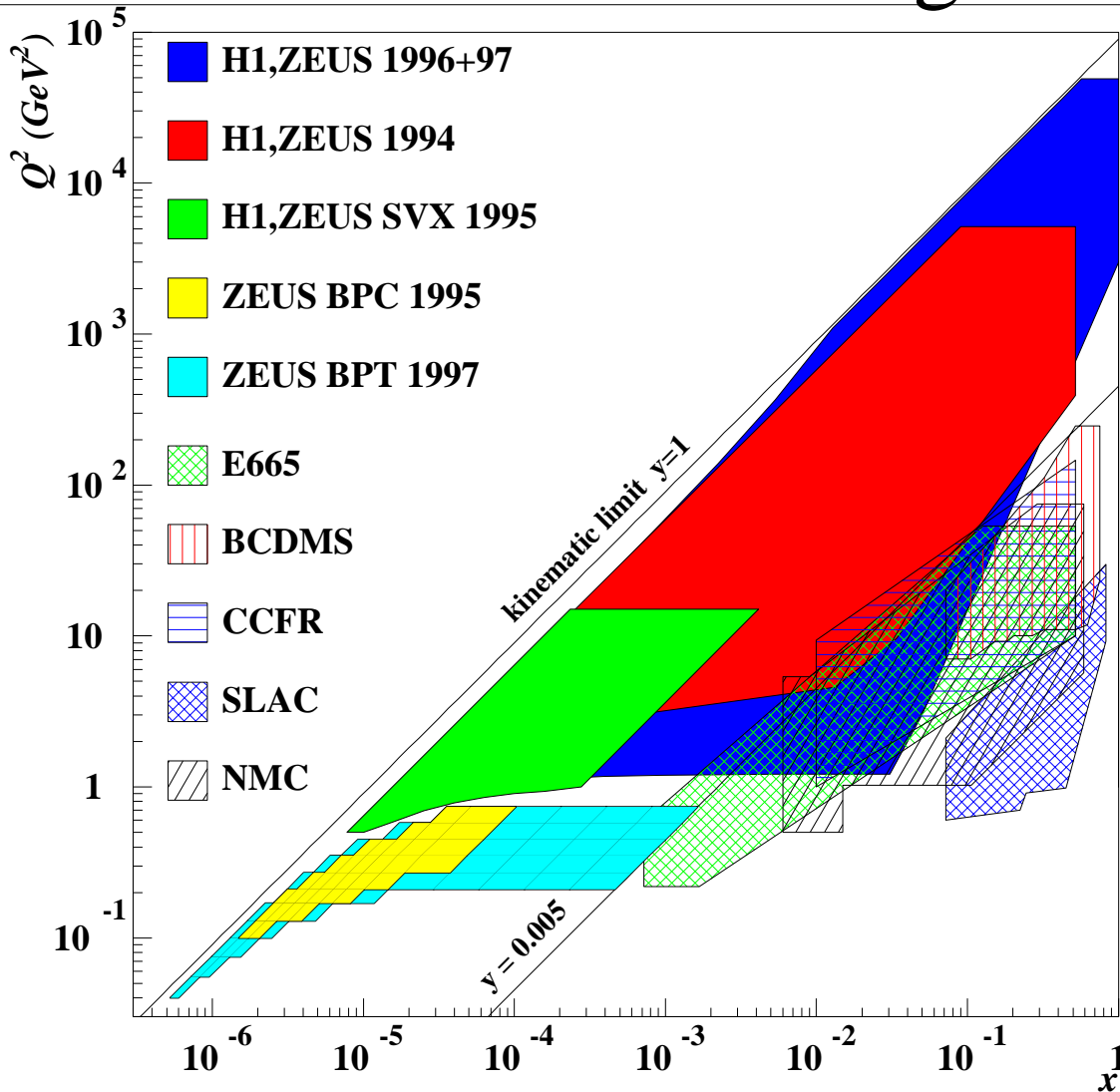


- As scale gets smaller (Q^2 increases) it is more likely for a parton to split into daughter partons at lower x

$$\frac{dq_i(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} \left[q_i(y, Q^2) P_{qq} \left(\frac{x}{y} \right) + g(y, Q^2) P_{qg} \left(\frac{x}{y} \right) \right]$$

$$\frac{dg(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} \left[\sum_i q_i(y, Q^2) P_{gq} \left(\frac{x}{y} \right) + g(y, Q^2) P_{gg} \left(\frac{x}{y} \right) \right]$$

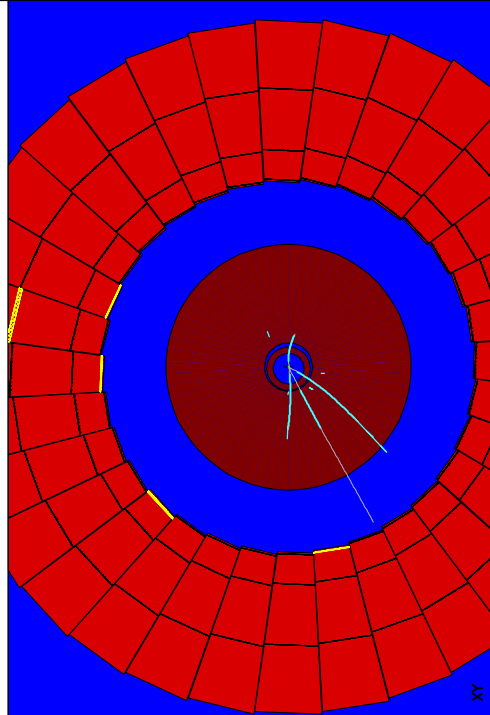
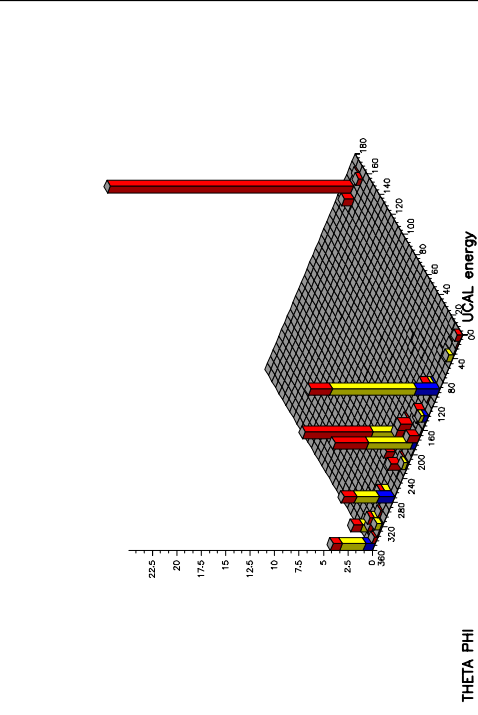
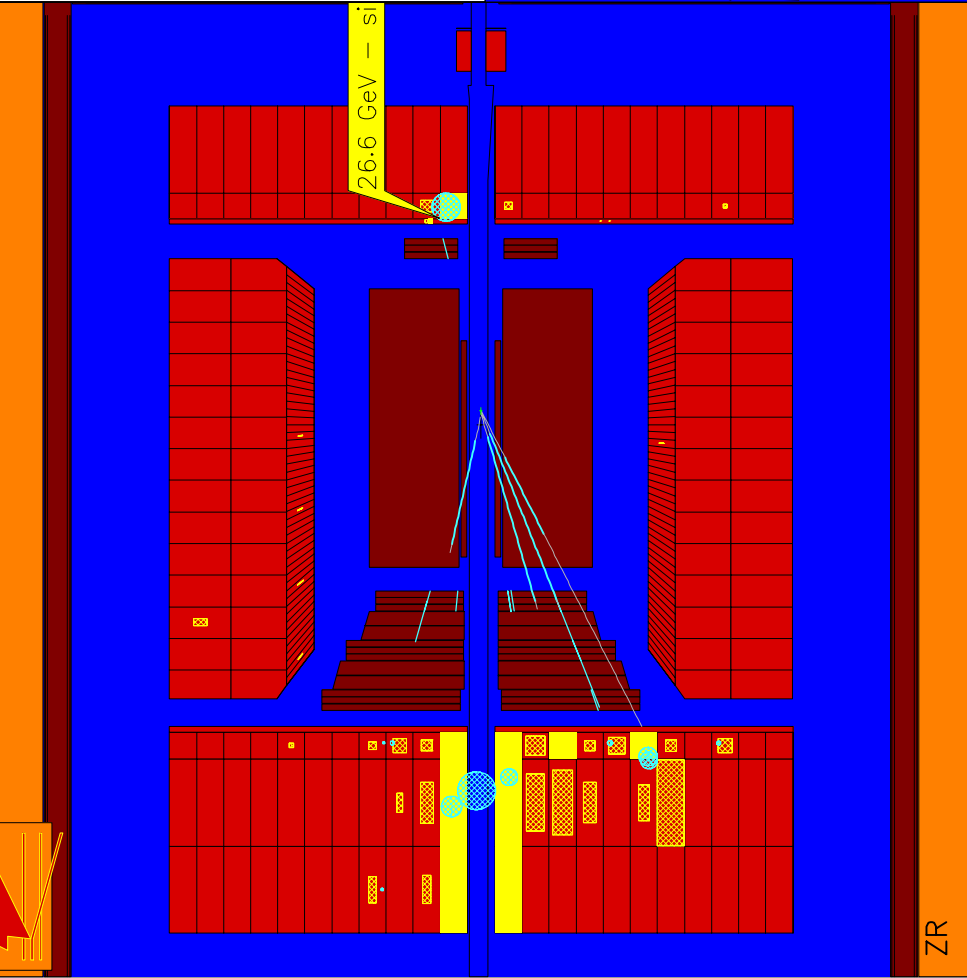
Kinematic Coverage



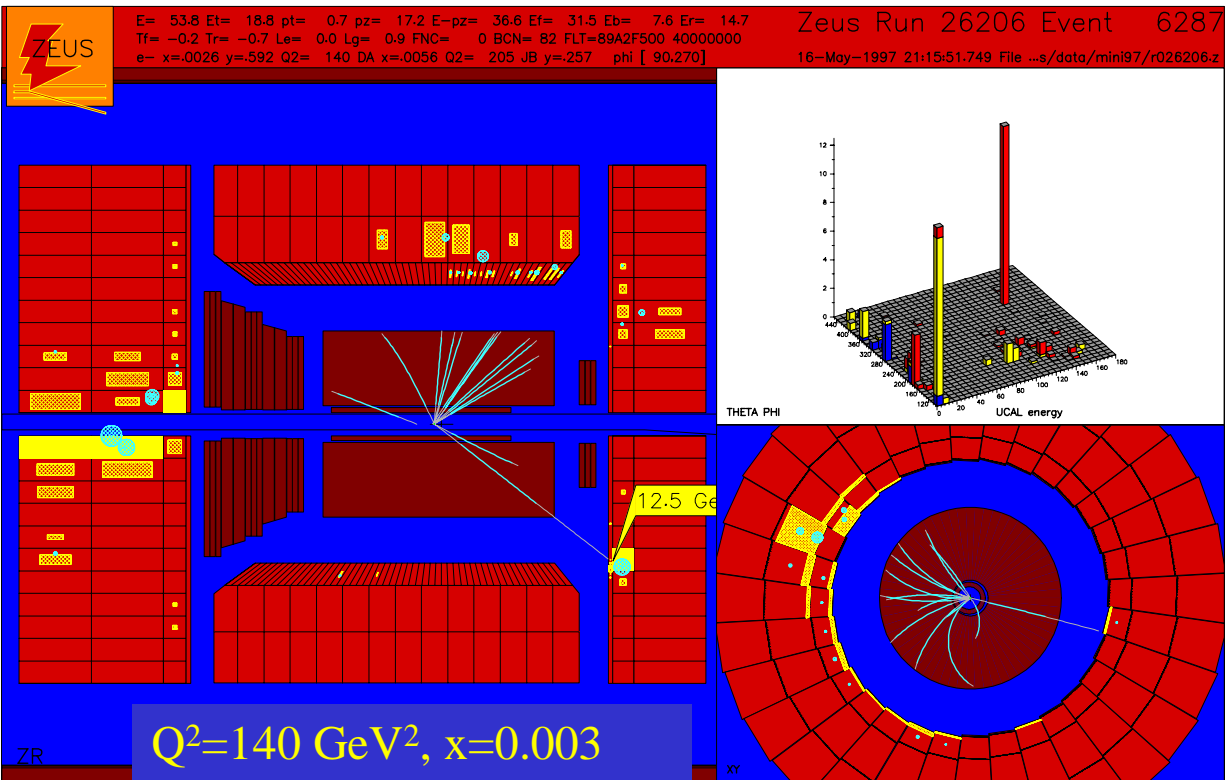
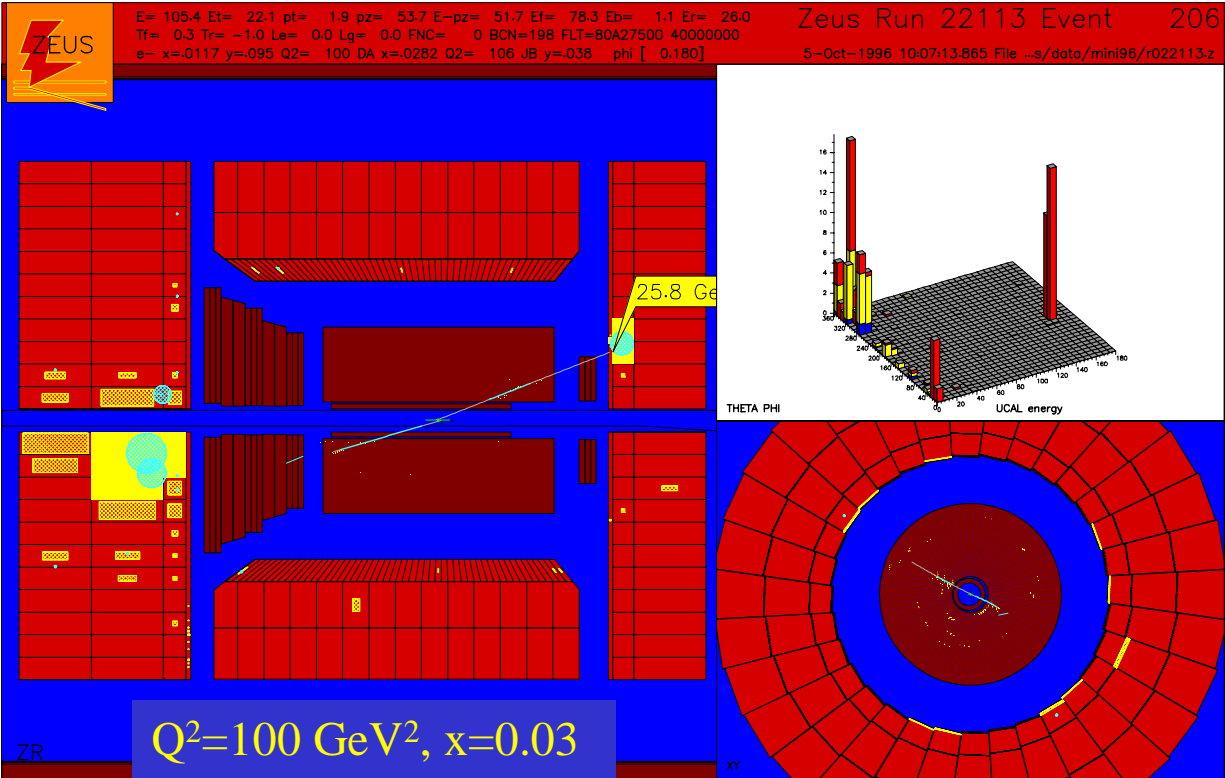
- 1997 and 1997 F_2 analysis
 - Covers 1994+1995 shifted vertex data
 - Reaches higher Q^2 and lower Q^2 in a single analysis
- BPC and BPT
 - electron detectors near the rear beam pipe to accept very low Q^2

Low Q^2 Event

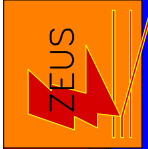

 Zeus Run 26882 Event 4911
 27-Jul-1997 9:00:07.608 File ...s/data/mini97/r026882.z
 E= 88.4 Et= 15.1 pt= 1.5 pz= 33.6 E-pz= 54.8 Ef= 61.2 Eb= 0.5 Er= 26.7
 Tf= 0.9 Tr= 1.5 Le= 0.3 Lg= 0.4 FNC= 0 BGN= 58 FLT=01826000 40000000
 e- x=-0.082 y=-0.43 Q2= 32 DA x=-0.098 Q2= 32 JB y=-0.34 phi [0.180]



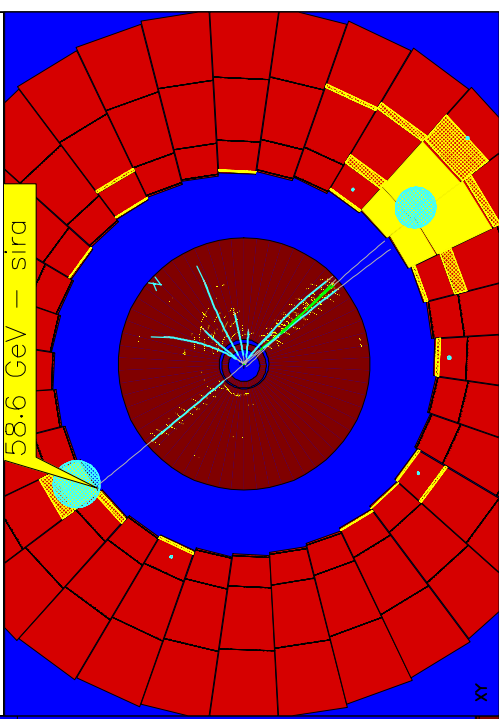
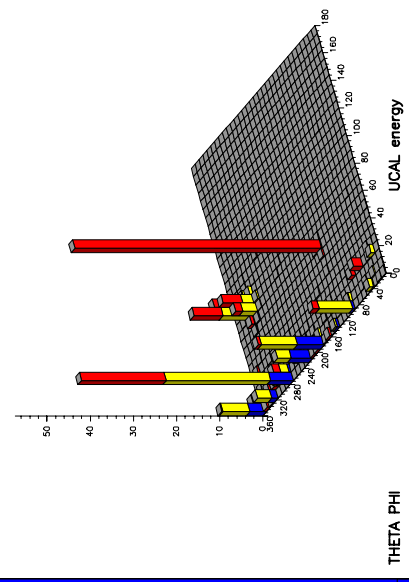
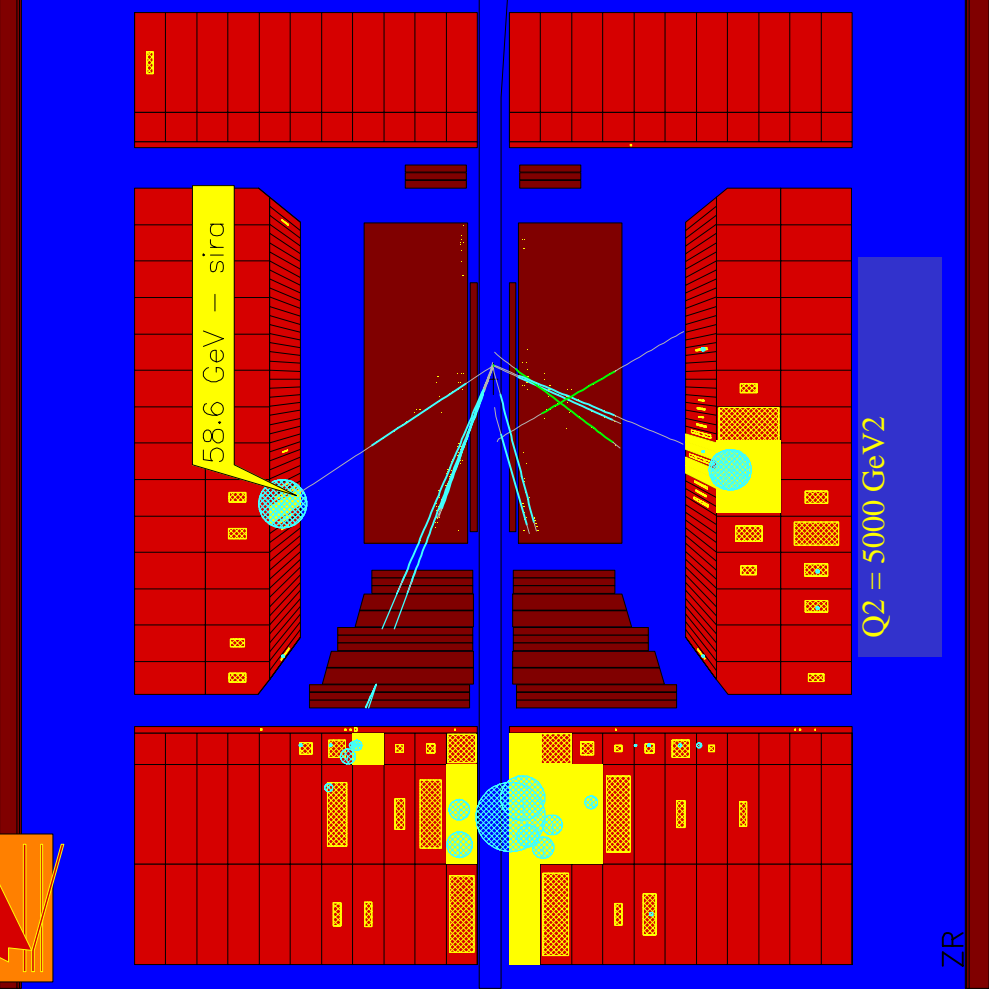
Medium Q^2



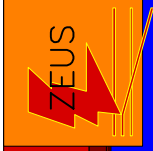
High Q^2 Event



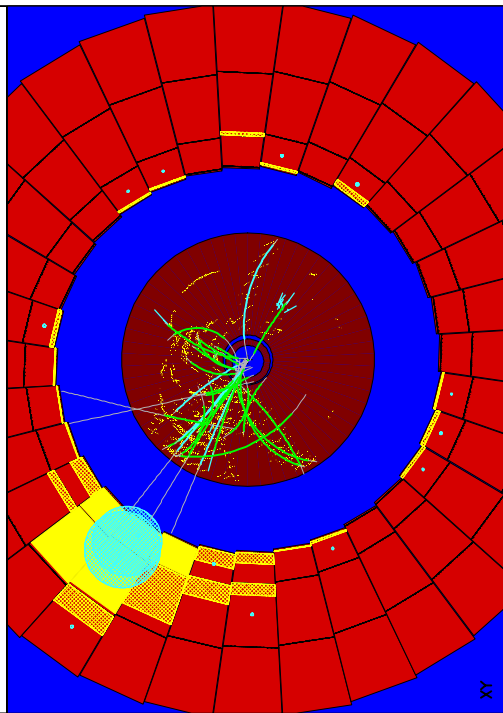
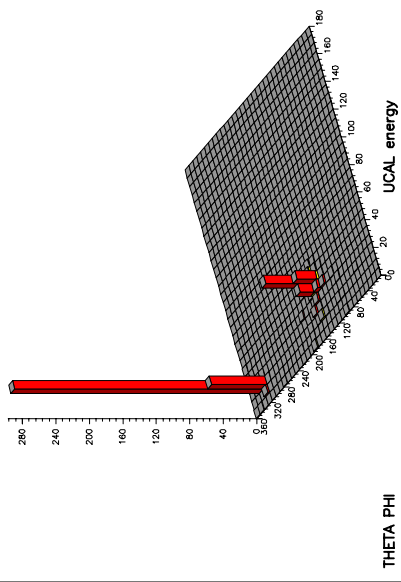
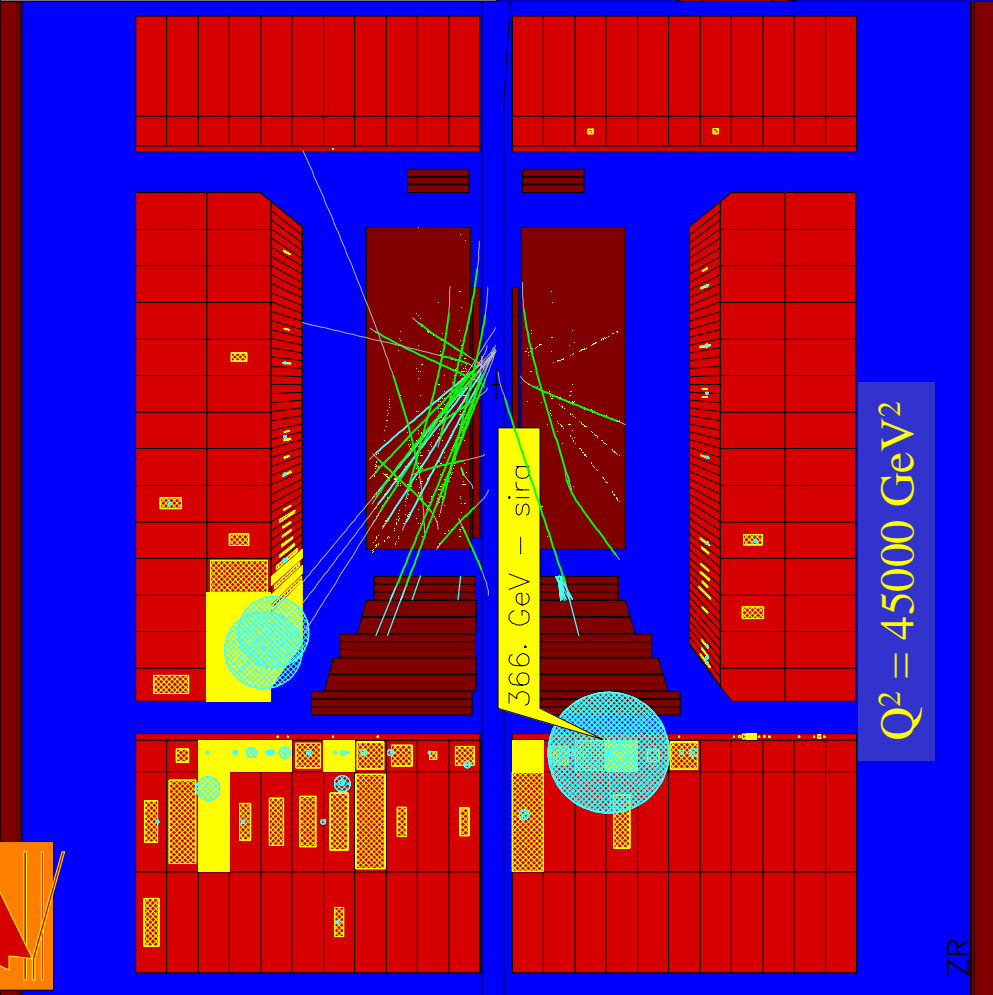
$E = 233.2$ Et= 101.7 pt= 8.1 pz= 179.6 E-pz= 53.6 Ef= 131.5 Eb= 101.6 Er= 0.1
 $Tf = -0.3$ Tr= 99.0 Le= 0.0 Lg= 0.1 FNC= 0 BCN=108 FLT=91A23F00 100000000
 e- x=-1062 y=-522 Q2= 5001 DA x=0960 Q2= 4739 JB y=493 phi [0.180]
 Zeus Run 22171 Event 16220
 10-Oct-1996 18:55:31.571 File ...s/data/mini96/r022171.z



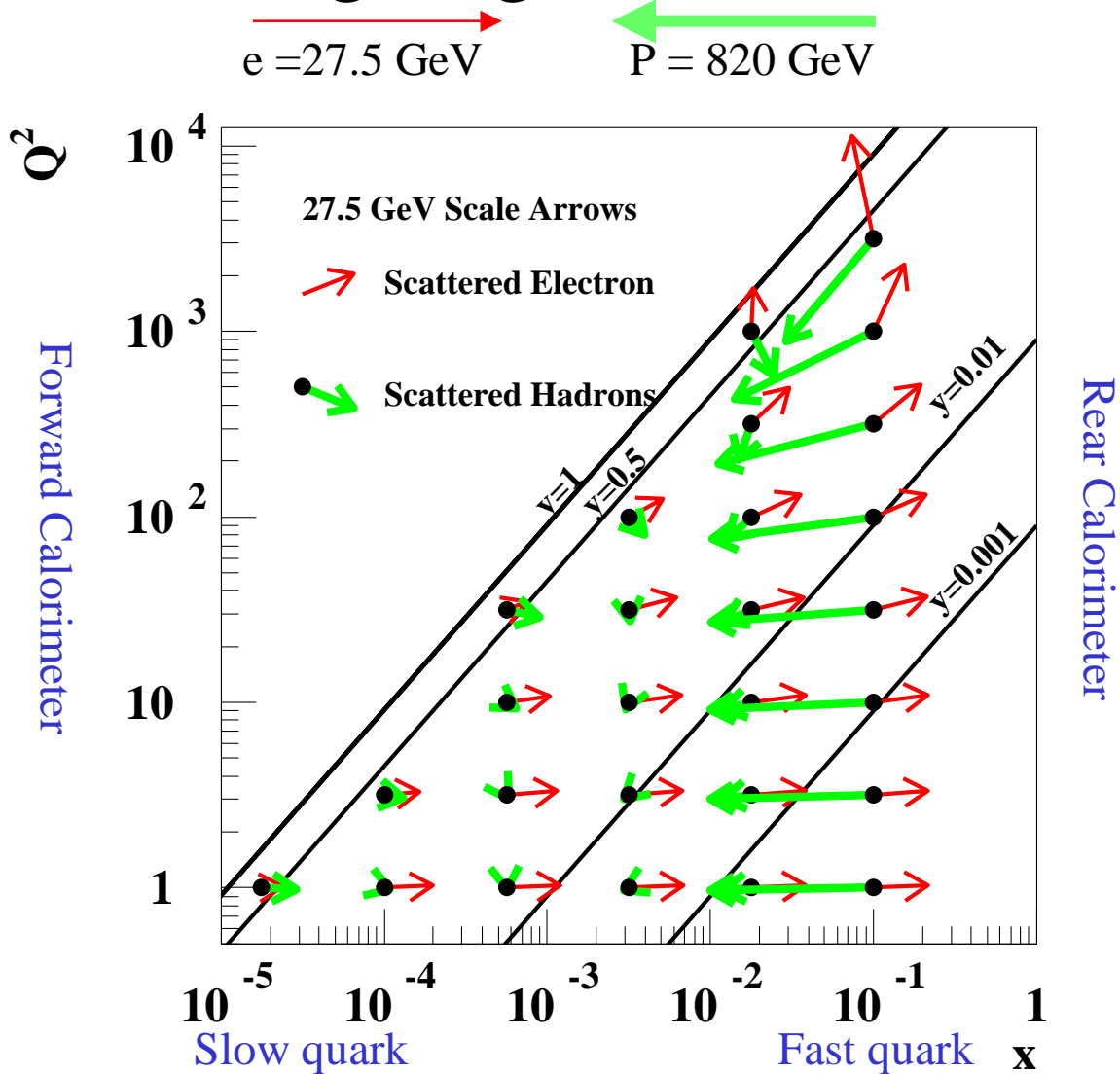
Highest Q^2 Event for 1996 and 1997



E= 540.0 Et= 198.7 pt= 6.2 pz= 491.6 E-pz= 48.4 Ef= 409.2 Eb= 130.7 Er= 0.1
 Tf= 0.5 Tr= 99.0 Le= 0.0 Lg= 0.0 FNC= 0 BCN=182 FLT=20A23F00 10000000
 e- x=-5727 y=-766 Q2=39588 DA x=-6770 Q2=44882 JB y=-629 phi [0.180]
 Zeus Run 22204 Event 11227
 12-Oct-1996 12:30:49.232 File ...s/data/mini96/r022204-z



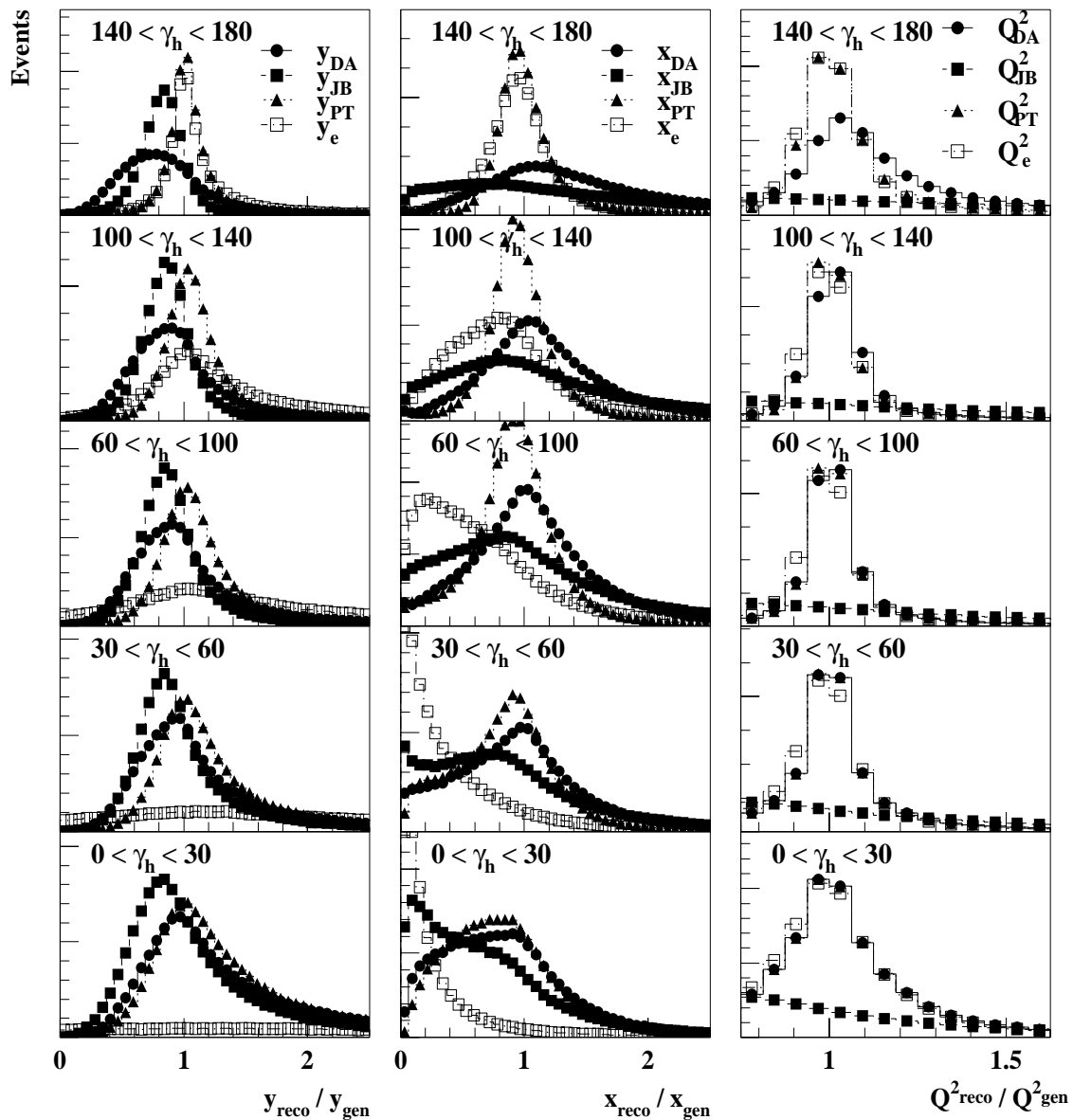
Outgoing Particles.



- Any 2 variables $(E_e, \theta_e, E_h, \gamma_h)$ determine x, Q^2
 - Electron Method (E_e, θ_e)
 - Jacquet Blondel Method (E_h, γ_h)
 - Double Angle Method (θ_e, γ_h)
 - Can also be used to predict $E_e(\theta_e, \gamma_h)$!
 - P_T method combines $(E_e, \theta_e, E_h, \gamma_h)$
- Beam pipe holes limit acceptance at low y and high y

Reconstruction Resolutions

- P_T method
 - Uses a combination of electron variables (E_e, θ_e) and hadronic variables ($P_{T,h}, (E-Pz)_h \rightarrow \gamma_h$) to optimize the reconstructed x, Q^2

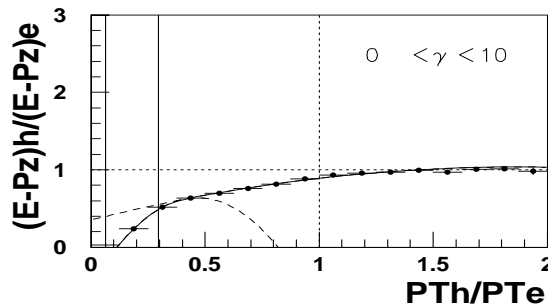


- P_T method superior over widest kinematic range.

P_T Method of x, Q^2 Recon.

$$\gamma = A \cos \left(\frac{P_{T,h}^2 - (E - P_z)_h^2}{P_{T,h}^2 + (E - P_z)_h^2} \right)$$

- Improve γ by improving $P_{T,h}$ and $(E-Pz)_h$
 - E-Pz is the dominant contribution to γ
- Replace $P_{T,h}$ with $P_{T,e}$
 - electron energy better calibrated
- Scale $(E-Pz)_h$ by a correction related to $P_{T,h}/P_{T,e}$



- P_T energy loss must also be E-Pz energy loss
- small Q^2 and γ dependence to correction
- Electron and hadrons at the same level of energy corrections.
- Scale $(E-Pz)_h$ by $2E_{\text{beam}}/(E-Pz)_{\text{total}}$
 - Corrects overall E-Pz energy scale
- Use angles (θ_e, γ_h) method to reconstruct x, Q^2

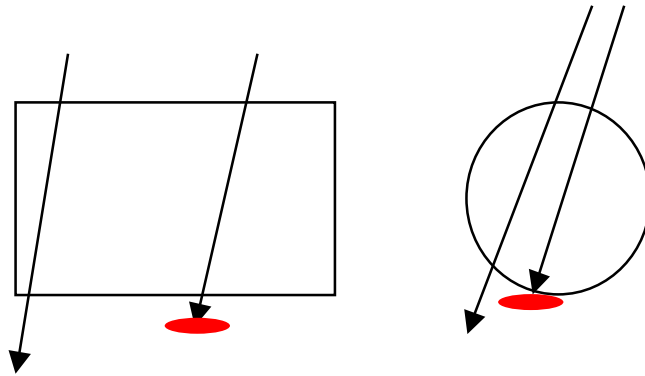
F₂ Measurement

- Select deep inelastic scattering events
 - Backgrounds
 - Selection cuts
- Reconstruct x , Q^2
 - Electron
 - Energy E_e
 - Angle θ_e
 - Event Vertex
 - Electron position in the detector
 - Hadronic
 - $P_T, E-P_z$ to calculate γ_h
- Unfold F_2
 - Count events in bins of x , Q^2
 - Divide by luminosity to produce a cross section.
 - Correct for detector acceptance (Monte Carlo)
 - Interpret cross section in terms of F_2

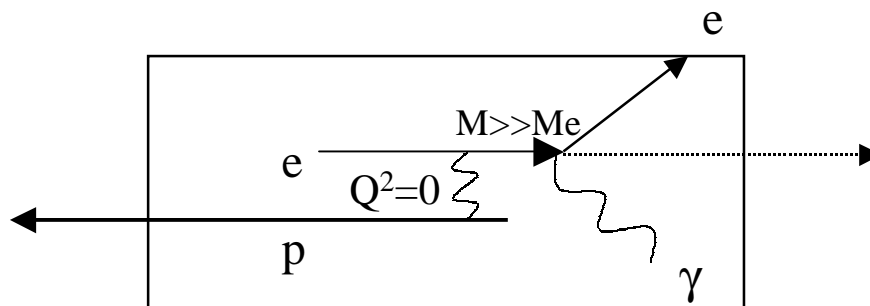
-
- Below $Q^2=800 \text{ GeV}^2$ the F_2 error is dominated by the understanding of the measured variables

Backgrounds

- Any event with the DIS electron in the detector is a signal event!
- Photoproduction
 - Electron escapes down the rear beam pipe $Q^2 \sim 0$
 - Hadronic activity fakes an electron
 - Dominant low Q^2 background
 - E-Pz Cut!
- Cosmic muons
 - Reject with P_T balance cuts and vertex cuts

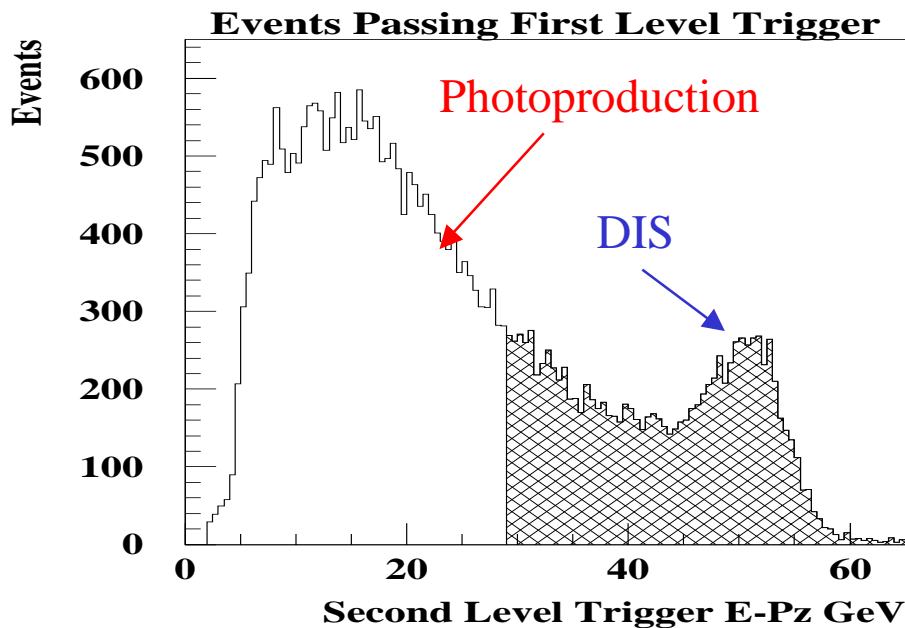
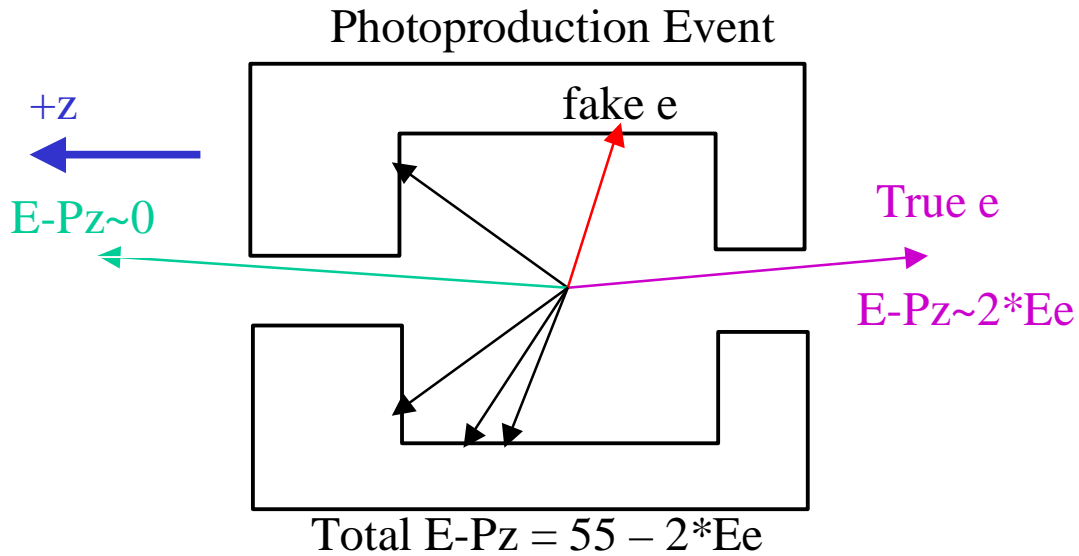


- QED Comptons ($Q^2=0$)
 - Reject events with only 2 electro-magnetic clusters



E-Pz Cut

- E-Pz is conserved
 - Before the event $E-Pz = E-Pz(e \text{ beam}) + E-Pz(P \text{ beam}) = 27.5 - (-27.5) + 820 - (820) \text{ GeV} = 55 \text{ GeV}$
 - If all particles are measured, one should measure 55 GeV.



Selection Cuts

- **Event Selection**

- First Level Trigger
 - Isolated calorimeter electromagnetic trigger tower $E > 2$ GeV in Rcal
 - Rcal EMC $E > 3.4$ GeV, Bcal EMC $E > 4.78$ GeV
 - Calorimeter $E_T > 30$ GeV, $E_T > 11.6$ GeV and CTD track
- Timing consistent with e+p interaction
- $38 < E - P_z < 65$ GeV (remove events with electron lost in rear beam pipe)
- $-50 < \text{Vertex } z < 50$ cm
- $P_T/E_T < 0.7$ and $P_T/\sqrt{E_T} < 3 \sqrt{\text{GeV}}$ (cuts cosmic events)
- $P_{T,h} / P_{T,e} > 0.3$ (well contained hadrons)

- **Electron Cuts**

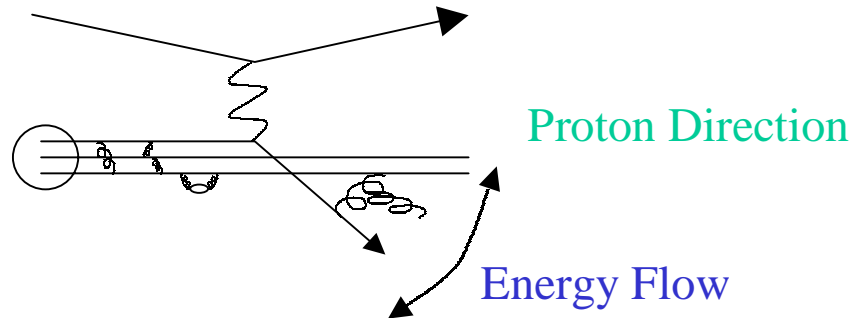
- Neural Net probability cut (function of energy)
- Energy > 8 GeV
- Position away from rear beam pipe
- Track Match (in tracking acceptance)
 - Distance of Closest Approach < 10 cm
 - Track Momentum > 5 GeV
- Non-electron energy in cone around electron < 5 GeV
- Number of Calorimeter cells < 30 in forward direction. (cut hadronic jets found as electrons)
- $y_{el} < 0.95$ (removes high Q^2 , high y fakes)

- **2 Data sets**

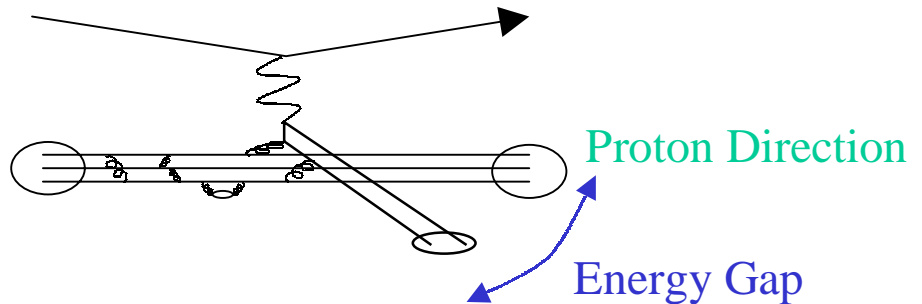
- $1.5 \text{ GeV}^2 < Q^2 < 27 \text{ GeV}^2$ 2.21 pb⁻¹ of luminosity
- $Q^2 > 27 \text{ GeV}^2$ 30.66 pb⁻¹ of luminosity

Monte Carlos

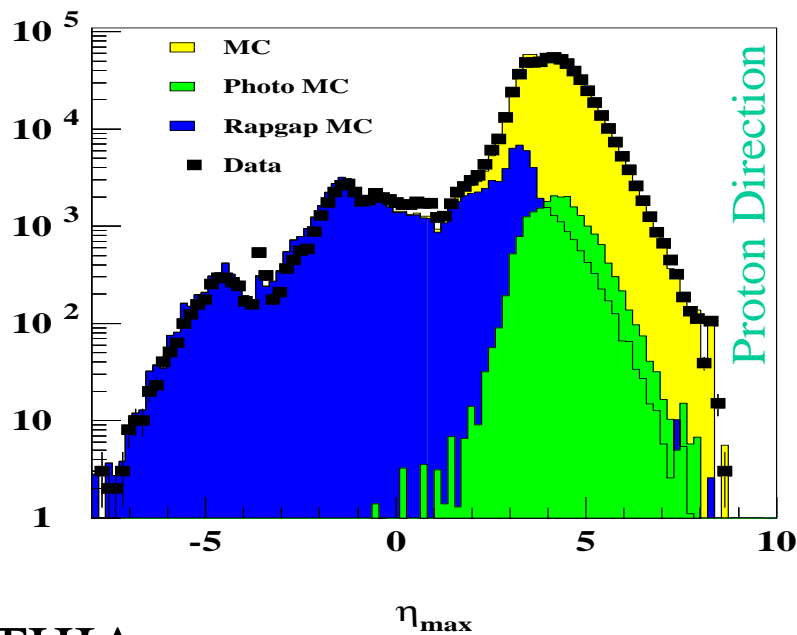
- DJANGO



- RAPGAP



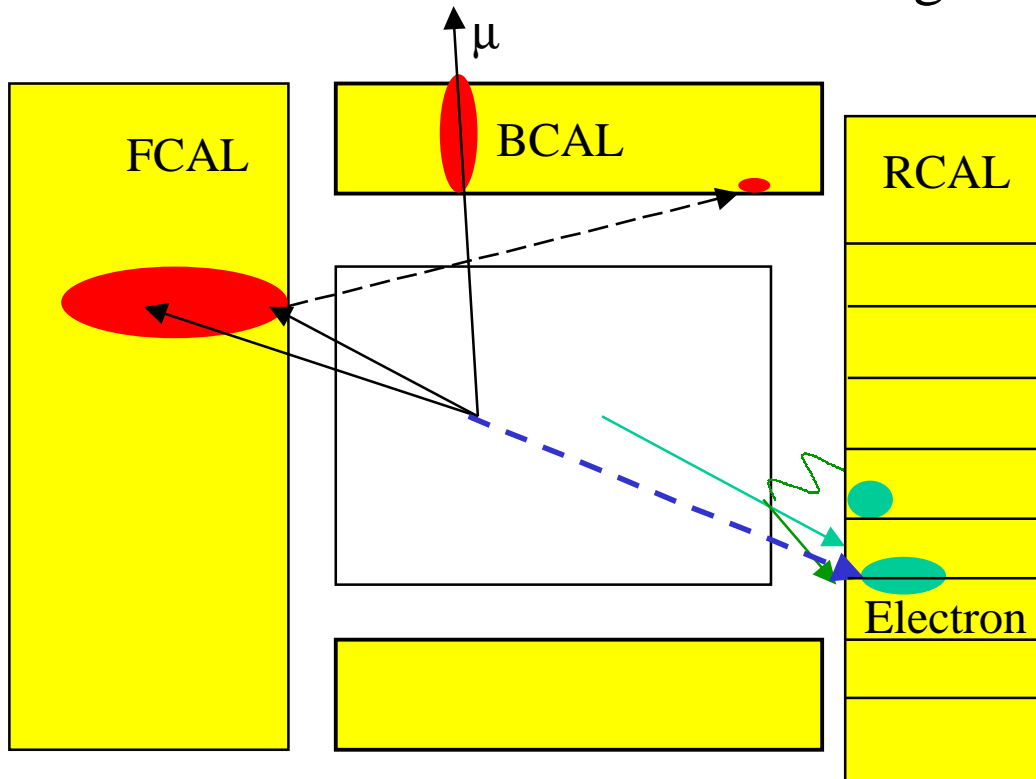
“Angle” of Energy Deposit Closest to Proton’s Direction



- PYTHIA

- Used for very low Q^2 background events where the electron escapes down the rear beam pipe

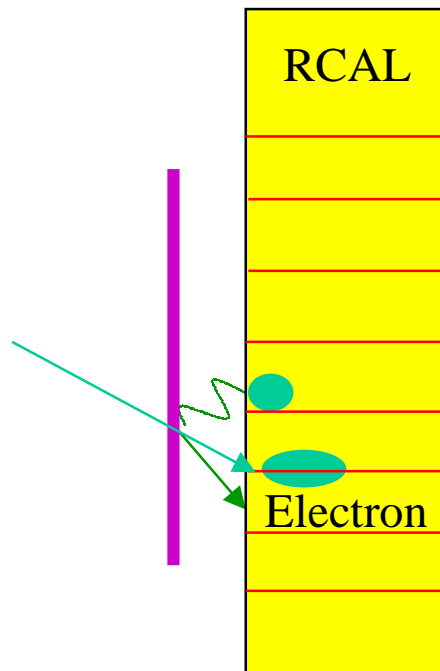
What are the reconstruction challenges?



- E_e
 - Calibration, Calorimeter edges, Loss in inactive material
- θ_e
 - Vertex position, Detector alignment
- Hadrons
 - Energy Calibration, μ energies, Backsplash
 - Vertex position, Detector alignment
- **Monte Carlo – Data Agreement**
 - Monte Carlo used to remove detector effects!

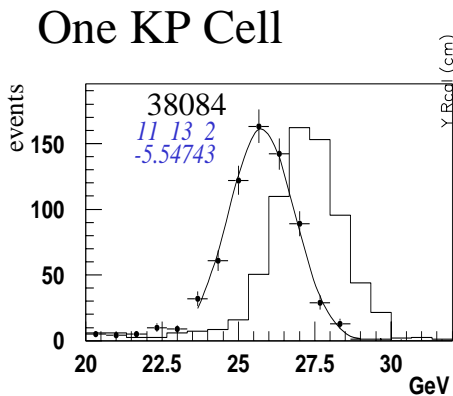
Electron Energy

- Important to calibrate the energy
 - Electron Energy Cut for DIS event selection
 - Used in P_T balance with hadrons to correct γ_h
- Calibration Steps
 - Calorimeter Cell calibration
 - Non-uniformity (Calorimeter Edges)
 - Dead Material Correction

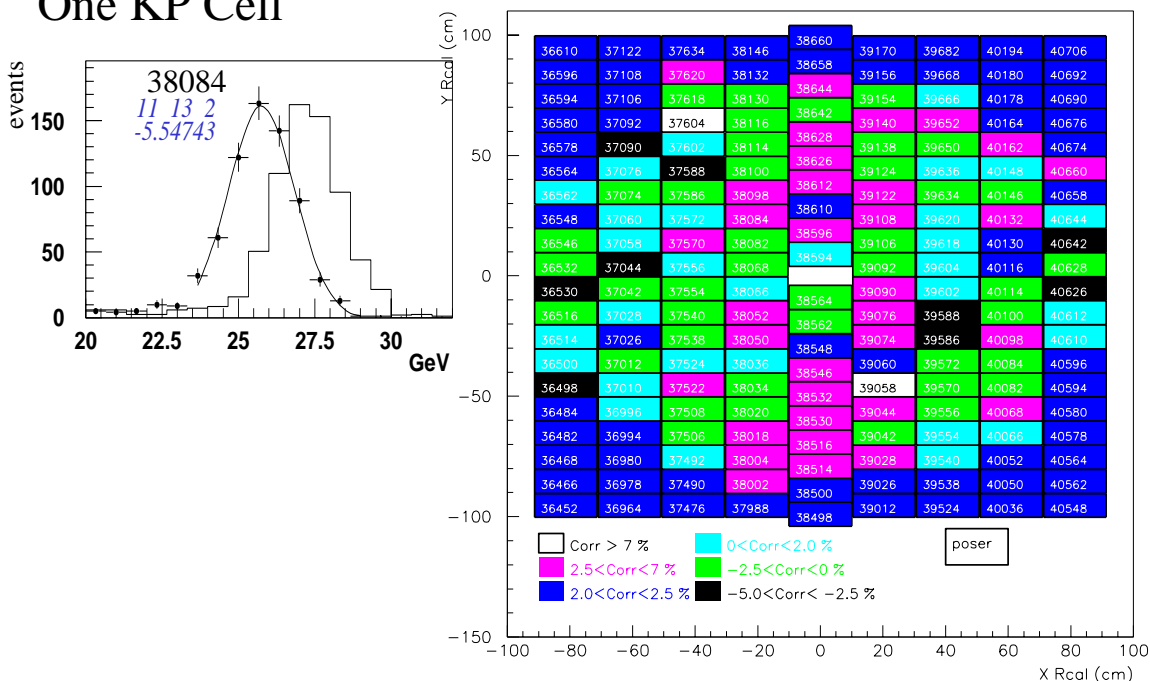


Calorimeter Cell Calibration

- Need a beam of known energy electrons
 - Double Angle Events.
 - From hadron and electron angles, the electron energy can be predicted. (2 independent variables)
 - Kinematic Peak events.
 - Special class of double angle events in which no hadronic activity is measured. ($\gamma_h = 0$)
- Force Data and Monte Carlo to be the same.

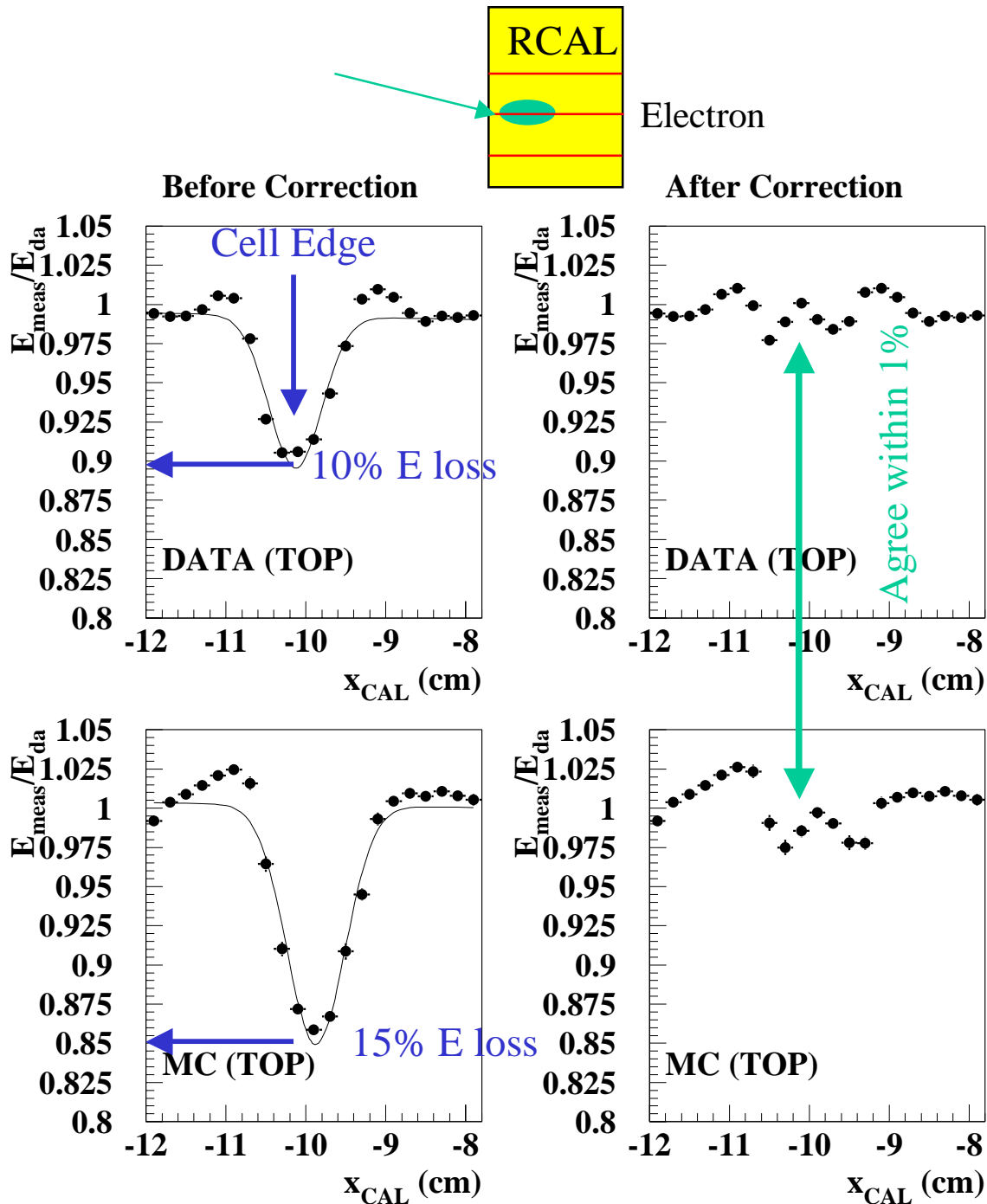


Cell by Cell Correction Factors Version 6.0



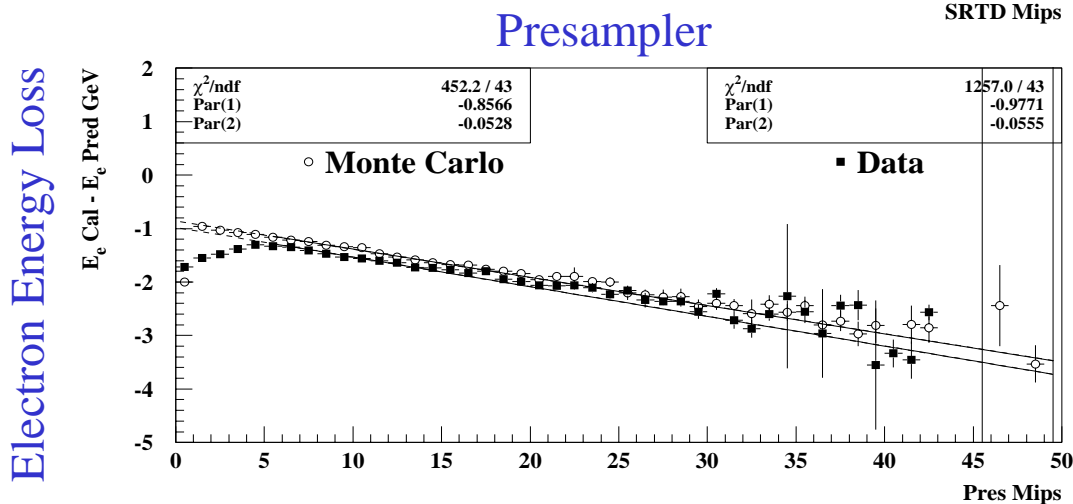
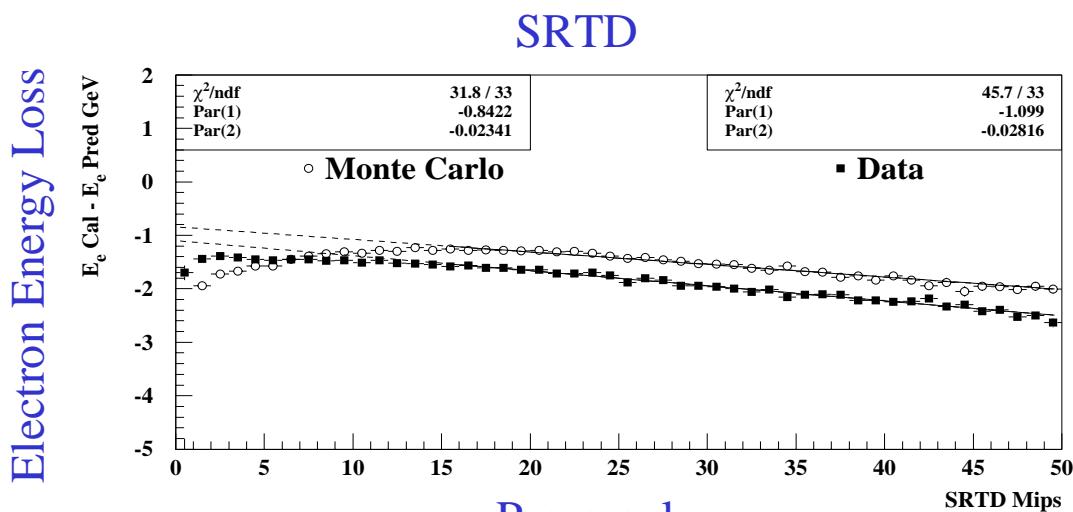
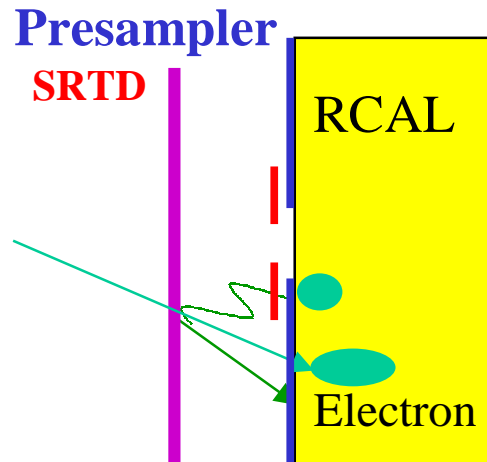
Non-Uniformity Correction

- Energy loss in calorimeter cell edges not well modeled by Monte Carlo.

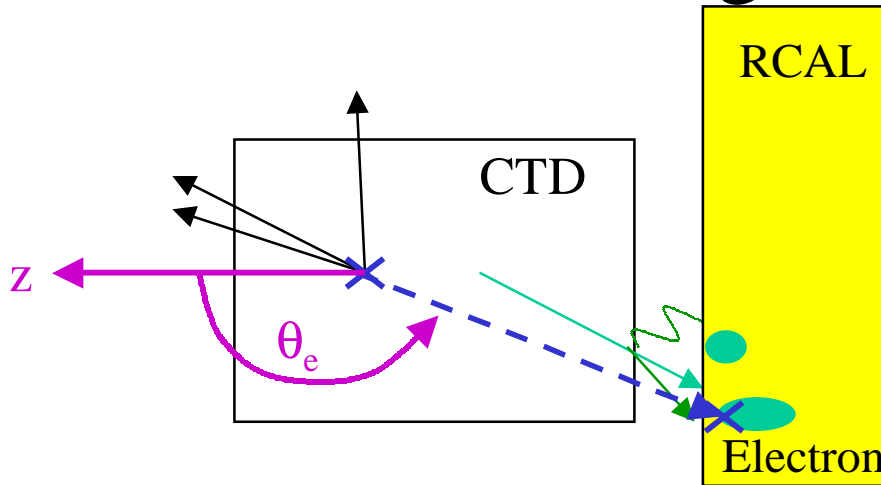


Dead Material Correction

- Electrons passing through material before reaching the calorimeter lose energy and produce a shower of particles.
- “Presampler” Method
 - Measure the number of charged particles at the face of the calorimeter and correct for the energy loss.



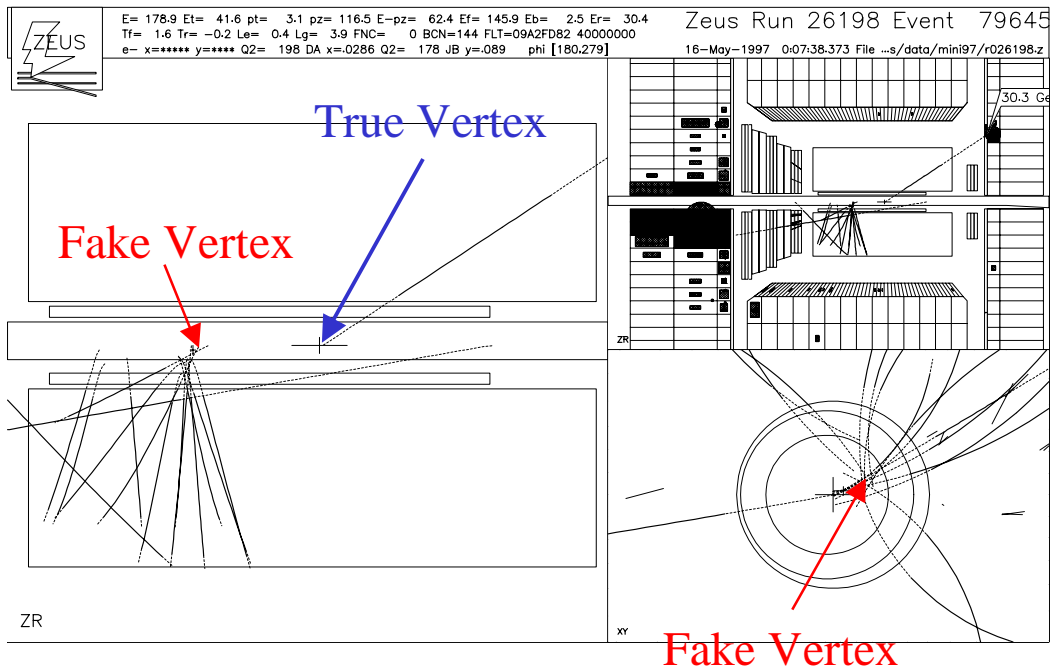
Electron Angle



- “True” Scattering Angle
 - Tracking angle at vertex
- A 2 point method is used when outside of tracking
 - Correct vertex z position
 - Electron track seed vertex
 - Fcal Timing vertex
 - Detector alignment
 - Aligned to track end point

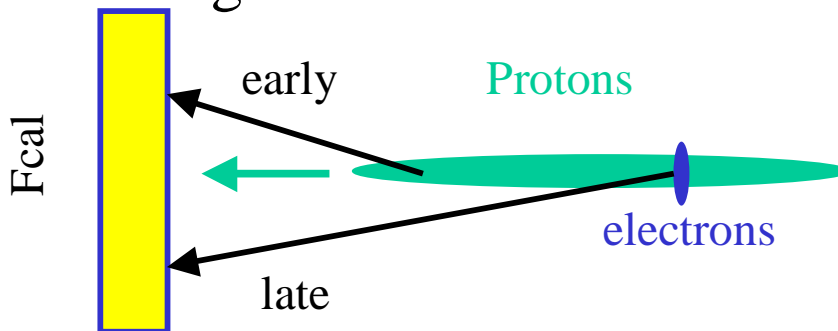
Event Vertex

- Proton bunch length 12 cm long
- Event by event vertex measurement
 - Angles of electron and hadrons measured with respect to the event vertex.
 - Electron seed vertexing method. 4 mm resolution. (Finds true vertex)



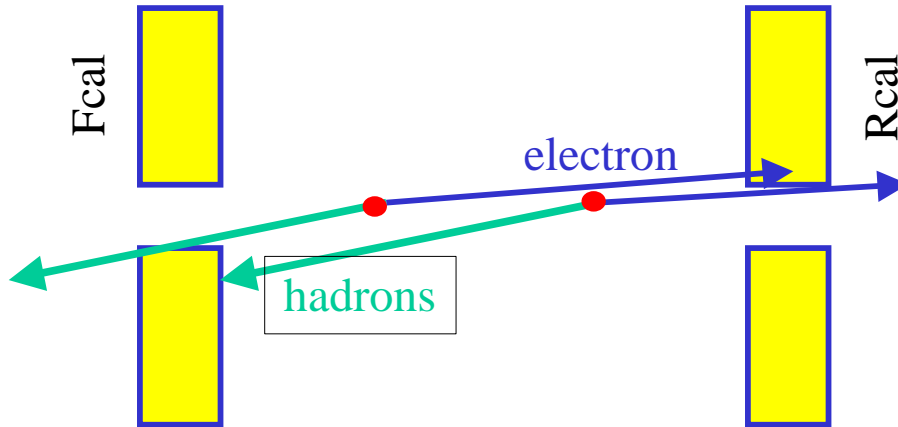
- Cut off axis vertices

- Fcal timing for no track events 8.6 cm resolution

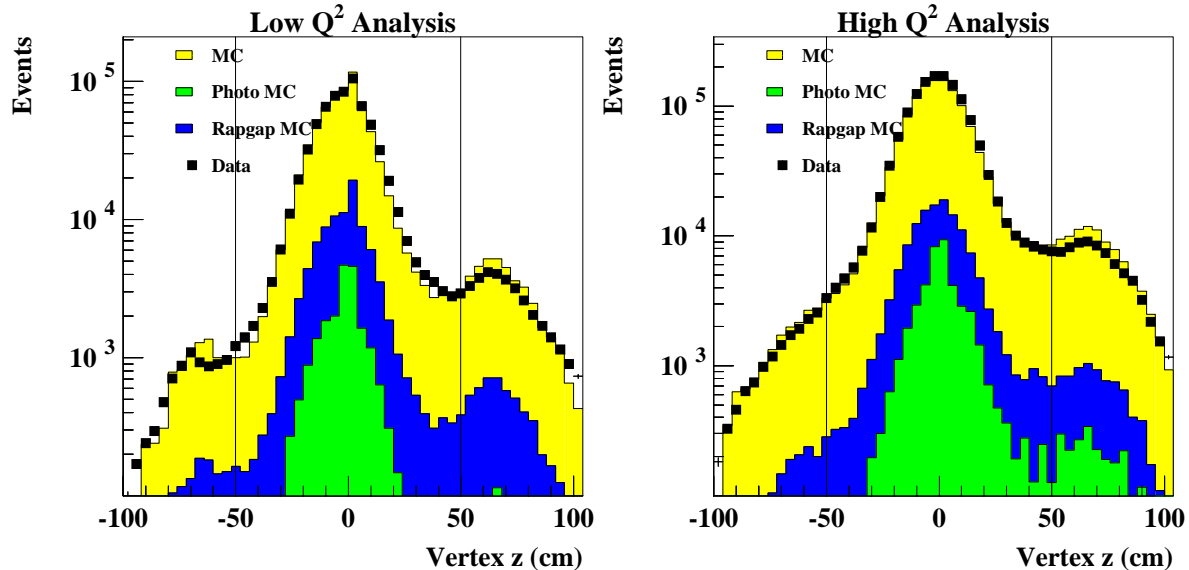


Vertex Distribution

- Acceptance depends on vertex



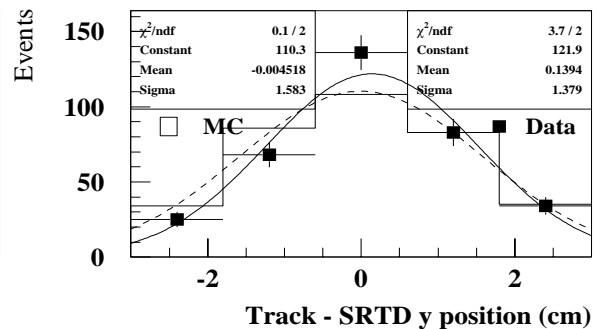
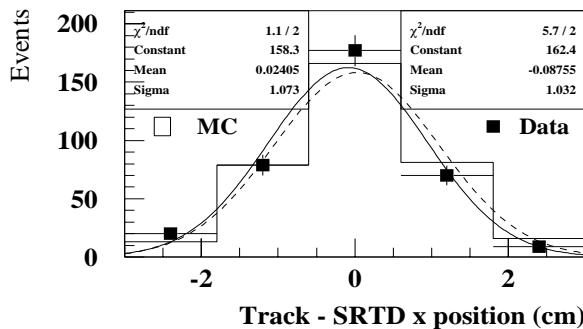
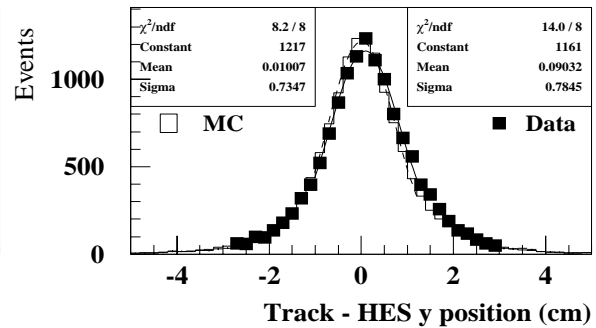
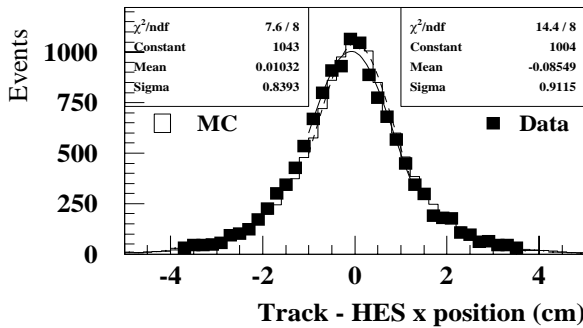
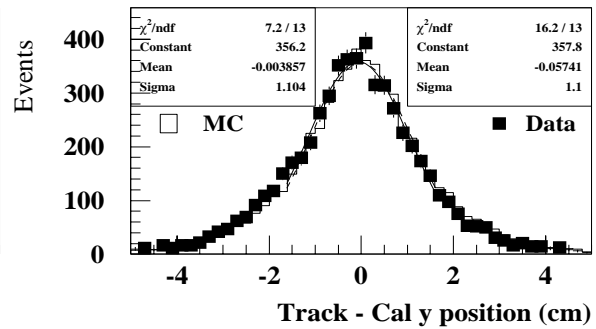
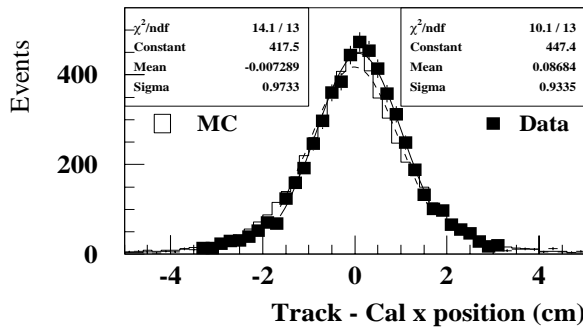
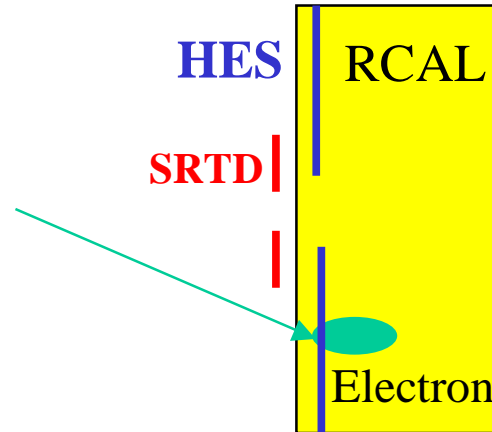
- Need same vertex distribution in data and Monte Carlo simulation!
 - Use unbiased event sample to measure data vertex distribution
 - Generate Monte Carlo events with this distribution.



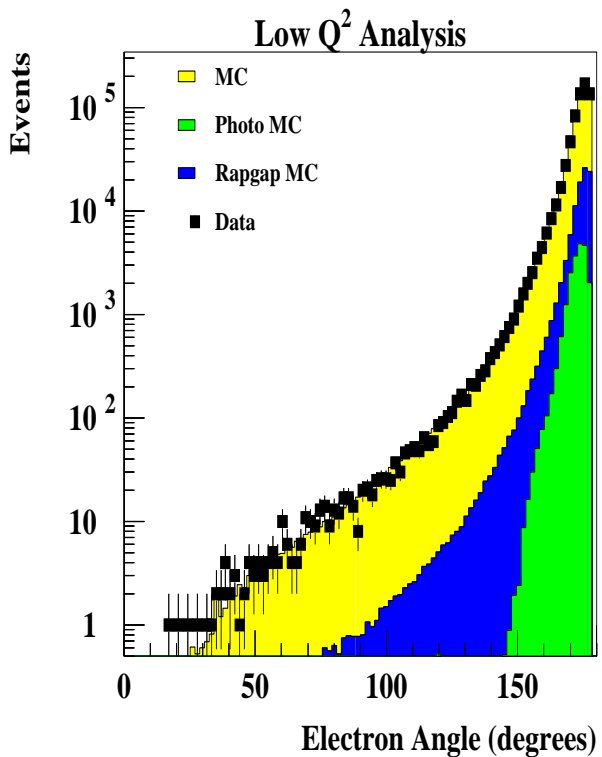
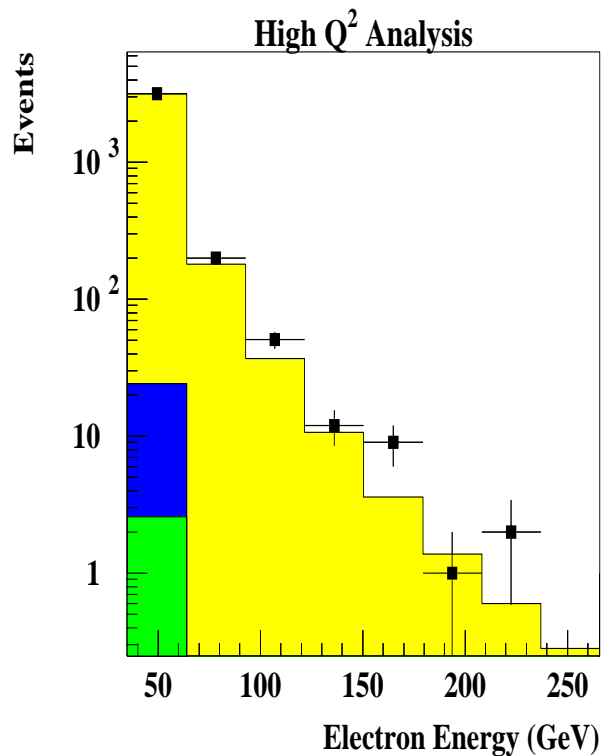
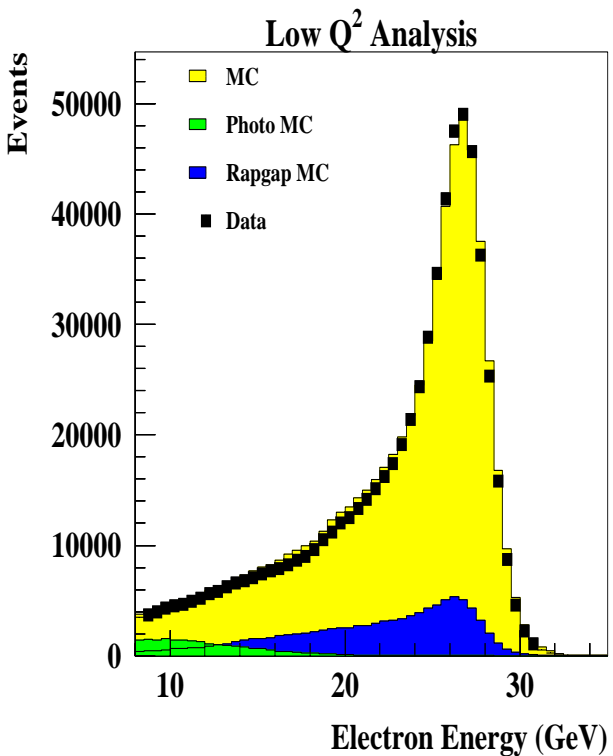
- Good Agreement in final DIS event sample

Detector Alignment

- Aligned to tracking
 - Calorimeter
 - HES (3cm x 3cm diodes 3cm inside rear calorimeter)
 - SRTD (1 cm wide scintillator strips near rear beam pipe)

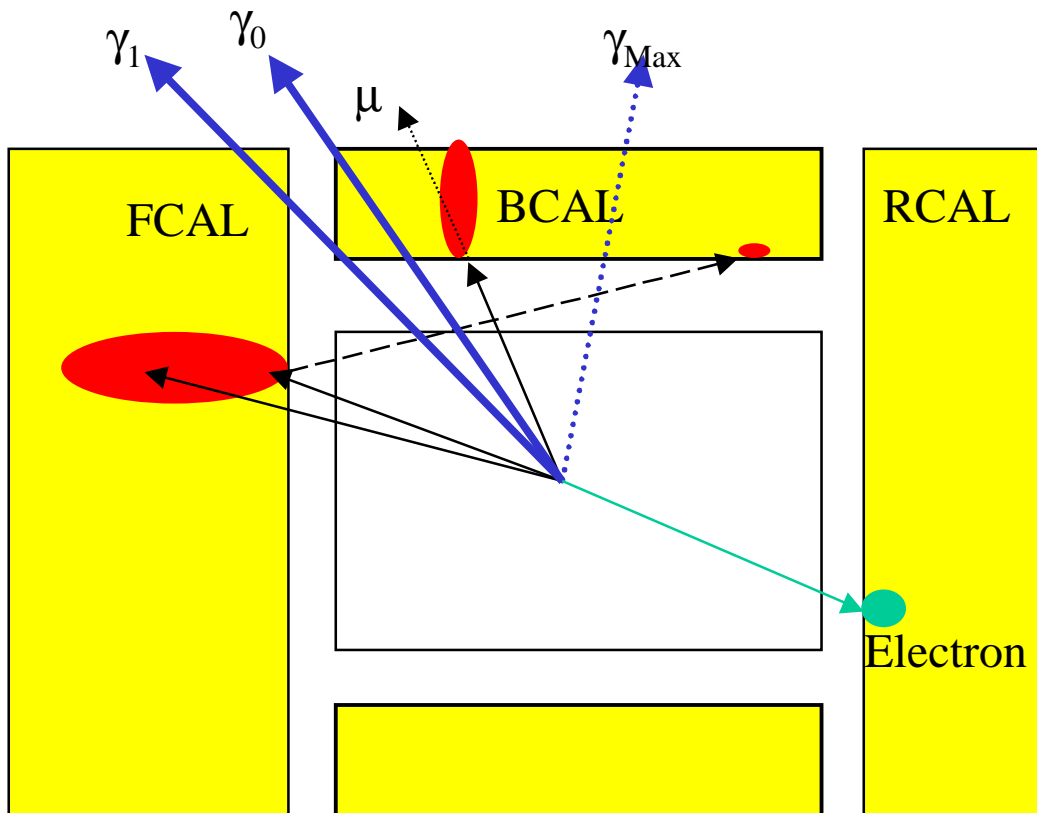


Electron Variables



- Electron energy
 - Photoproduction at low energies
 - High energy side of MC peak doesn't agree
 - accounted for in systematic error
 - Excess of data at very high energies (high Q^2)
- Electron Angle
 - agrees over 5 orders of magnitude

Hadronic Reconstruction



- Challenges... μ energies, Backsplash
- Combine tracking and calorimeter info. (P_{Th} , $E-Pz_h$)
 - Find calorimeter islands and match tracks to islands
 - Use tracking momentum if calorimeter energy \ll track momentum
 - Improves μ energies
 - Use tracking angle for 1 to 1 matches

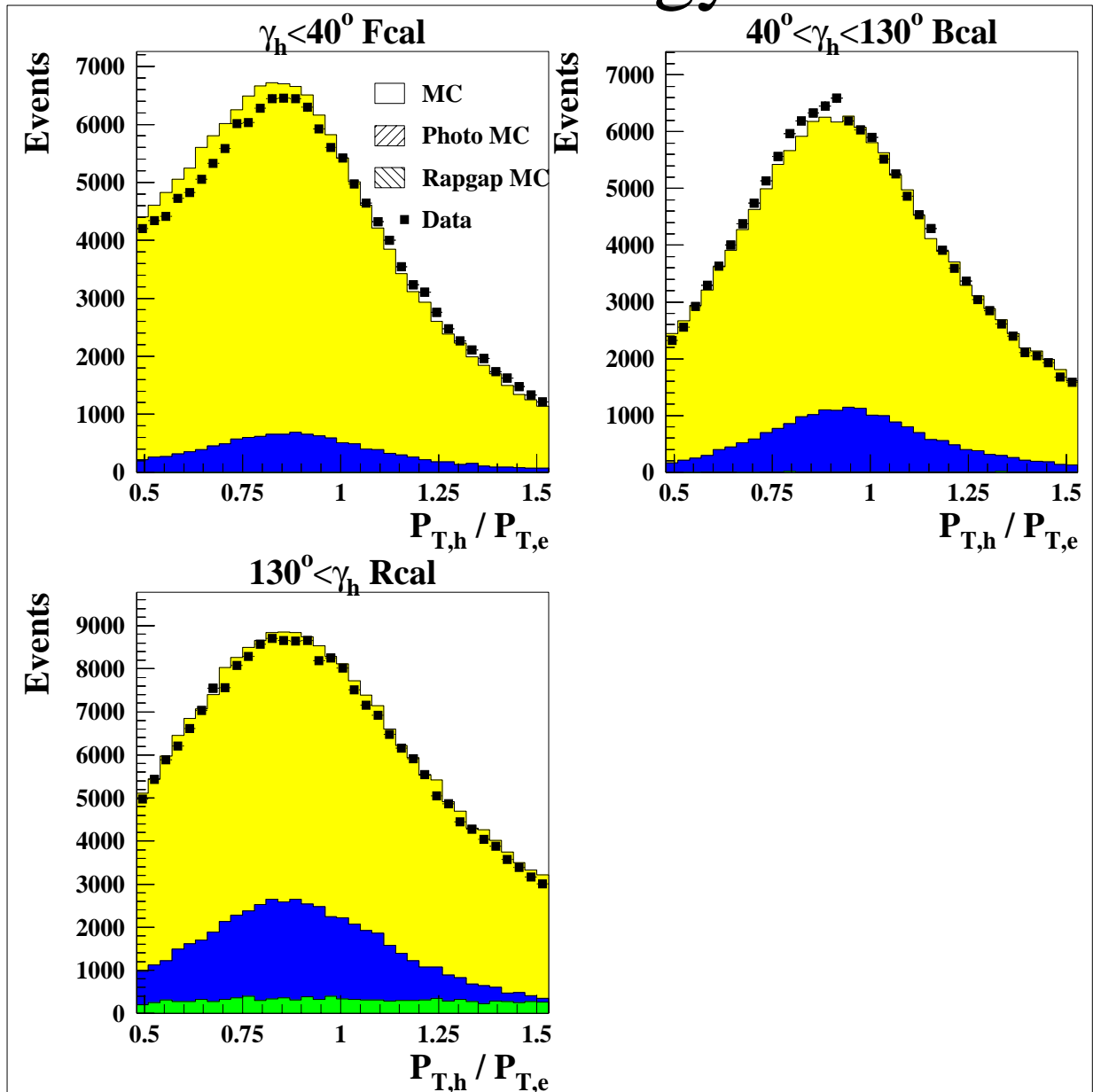
- Remove backslash

- Calculate γ from islands and $\gamma_{Max} = \gamma + 50^\circ$

$$\gamma = A \cos \left(\frac{P_{T,h}^2 - (E - P_z)_h^2}{P_{T,h}^2 + (E - P_z)_h^2} \right)$$

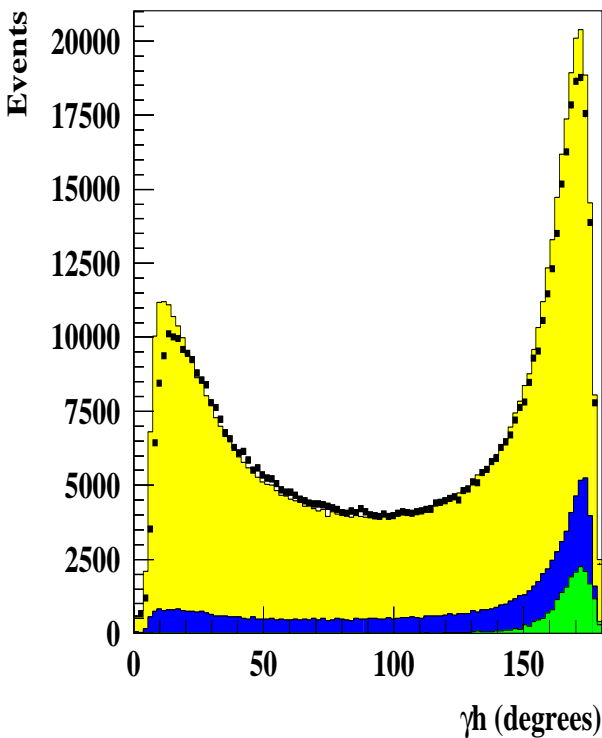
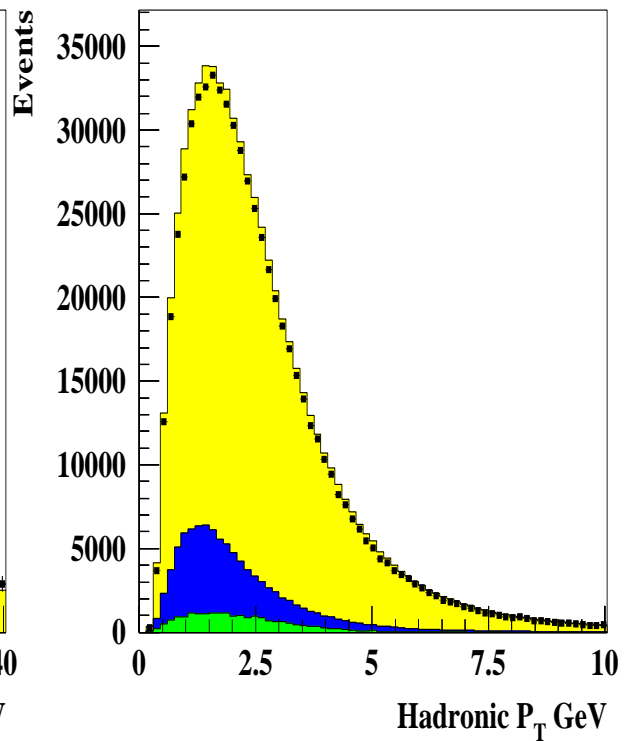
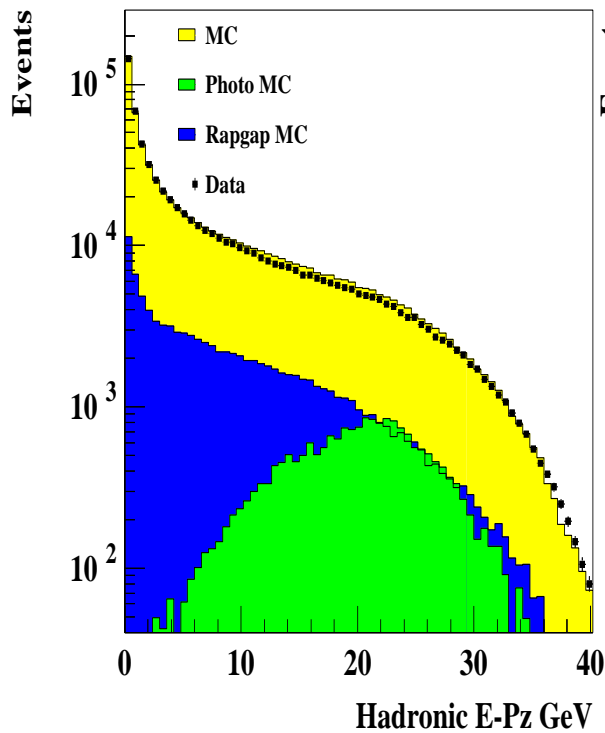
- remove low energy, trackless, isolated islands beyond γ_{Max}
- repeat up to 3 times or until γ is stable within 1%

Hadronic Energy Scale



- Need a hadronic “test beam” of known energy.
 - P_T balance with calibrated electron!
- Want Monte Carlo to match data.

Hadronic Variables



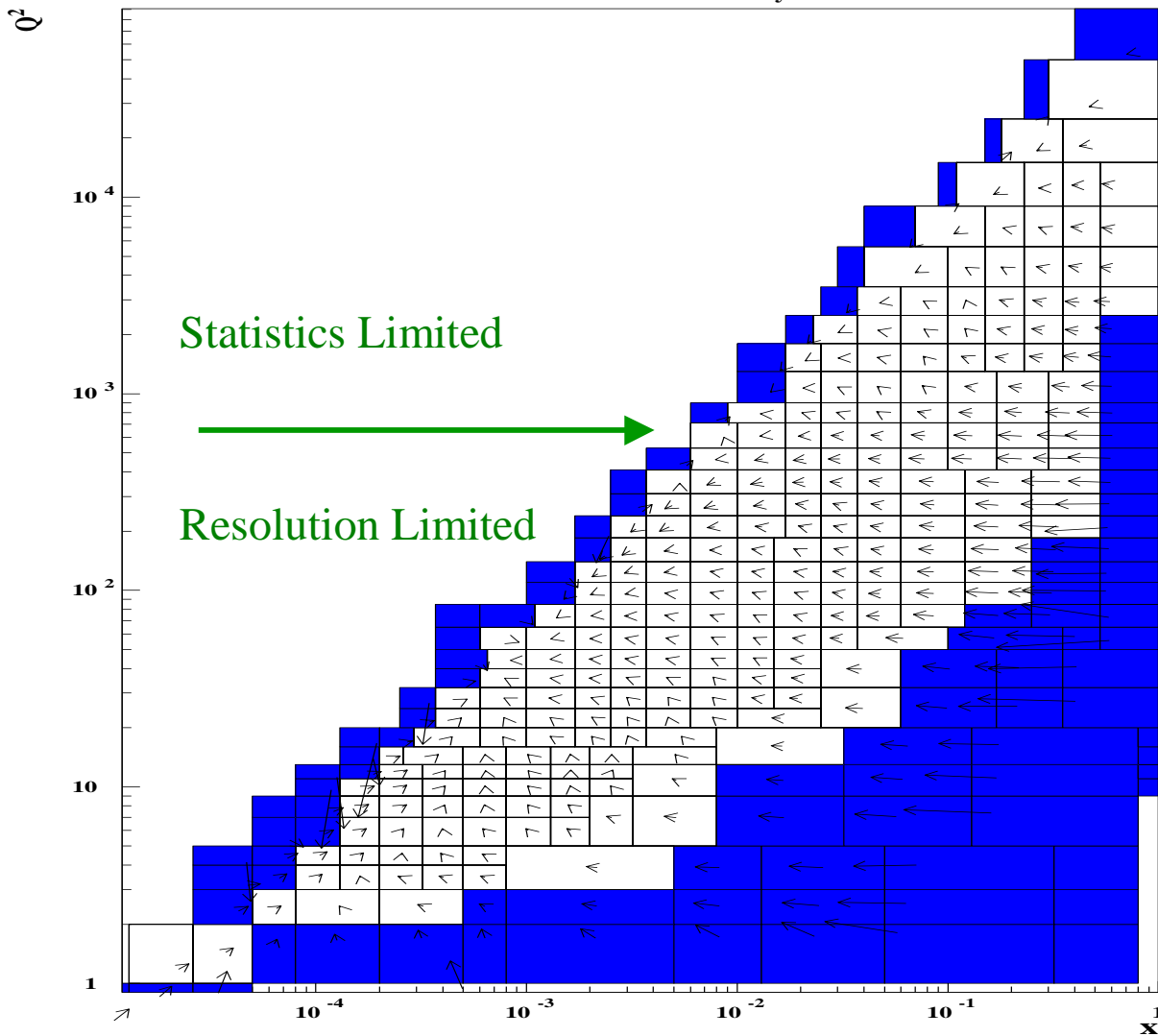
•RAPGAP/DJANGO fraction

- Medium E-Pz
- low P_T
- γ_h
 - very low γ not used
 - High γ RAPGAP fraction

Binning and Unfolding

MC Migration

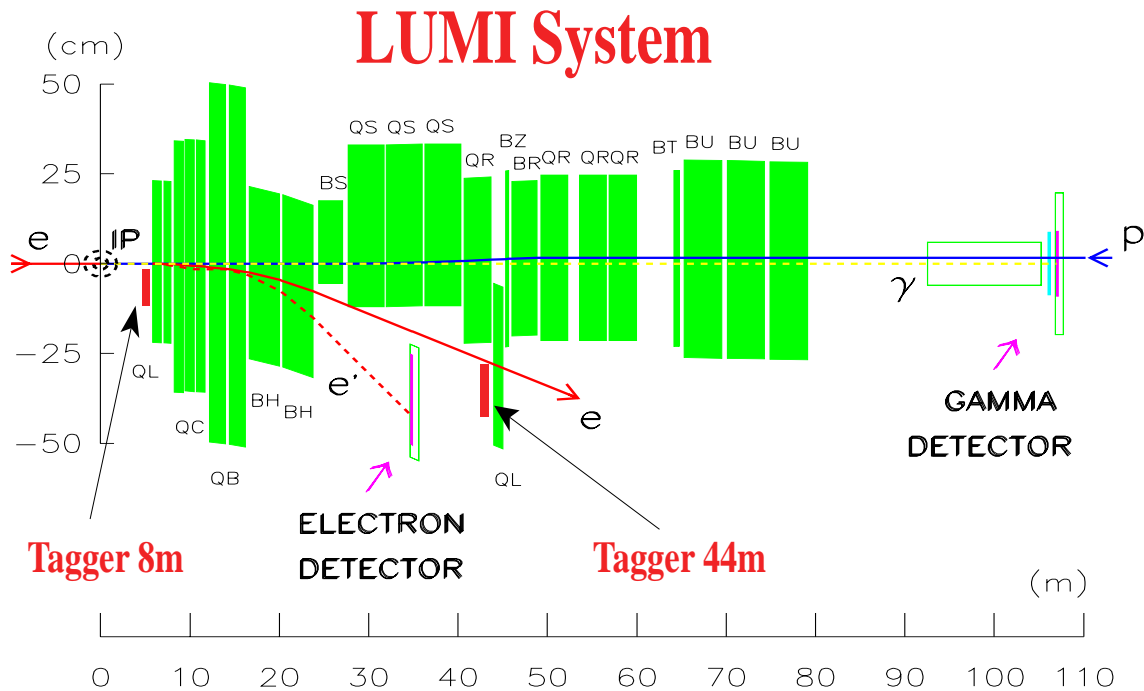
For Generated events in Bin where they are measured



- Count events in bin in data N_{data} and Monte Carlo N_{MC} for the same luminosity
- Extract F_2
$$F_2^{\text{data}}(x, Q^2) = F_2^{\text{MC}}(x, Q^2) \times \frac{N_{\text{data}}}{N_{\text{MC}}}$$
- Fit F_2^{data} , Reweight MC to fit, and repeat

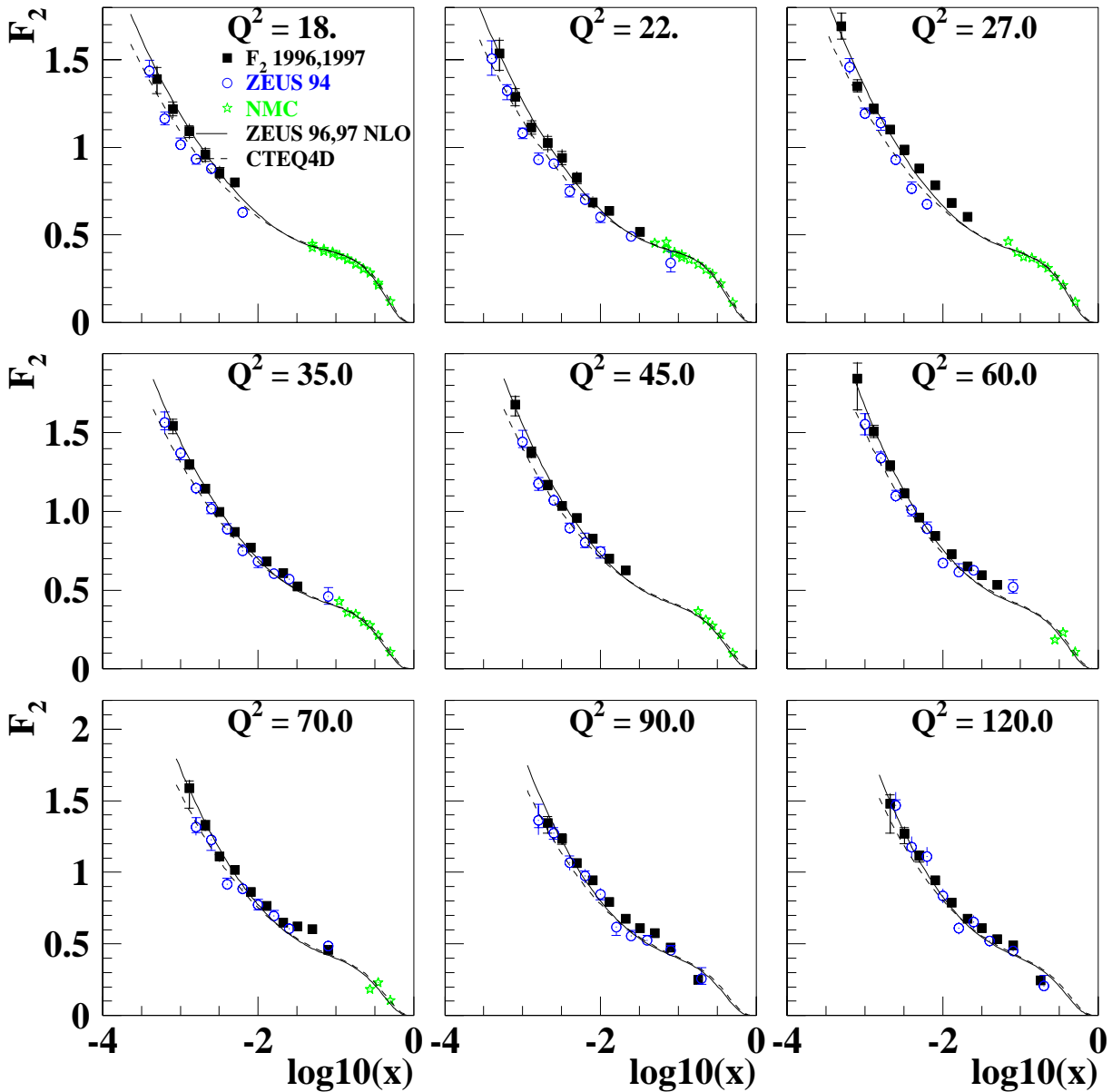
Luminosity

- Luminosity = how many particles pass per unit area (per unit time).



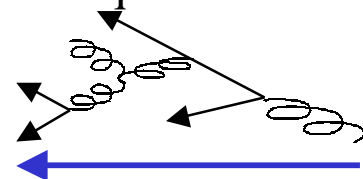
- Use precisely known cross section $\sigma(ep \rightarrow ep\gamma)$
 - Photon measured in gamma detector
 - Electron measured in 35 m tagger is used to calibrate the gamma detector energy.
- Luminosity = $\# \gamma \text{ events} / \sigma(ep \rightarrow ep\gamma)$
- Sets Normalization of F_2

Med Q^2 F_2

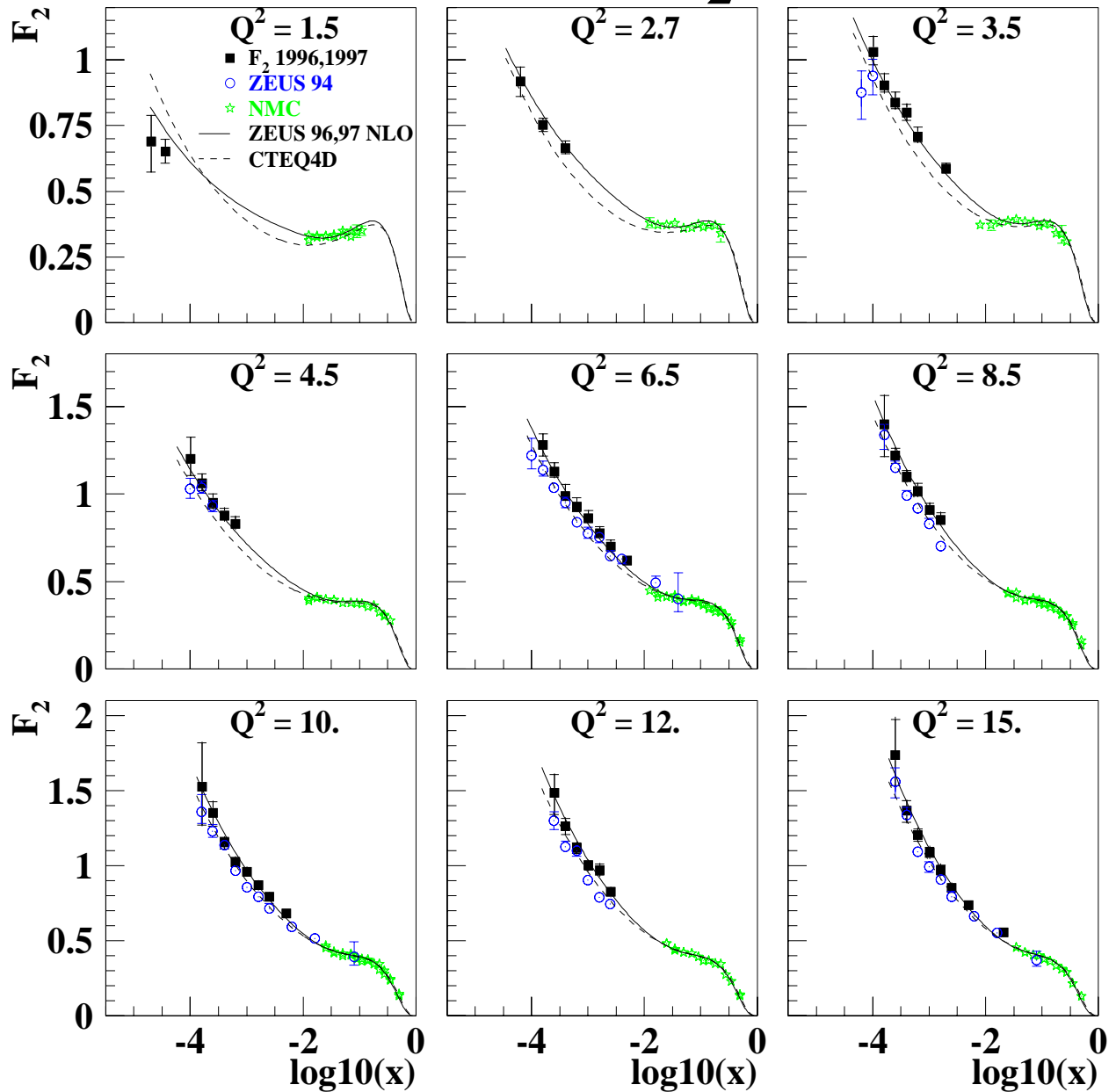


- CTEQ4D is a QCD DGLAP fit to the 1994 DATA
- ZEUS96&97 is a QCD DGLAP fit to this data
- Strong rise in F_2 at low x is from sea quarks

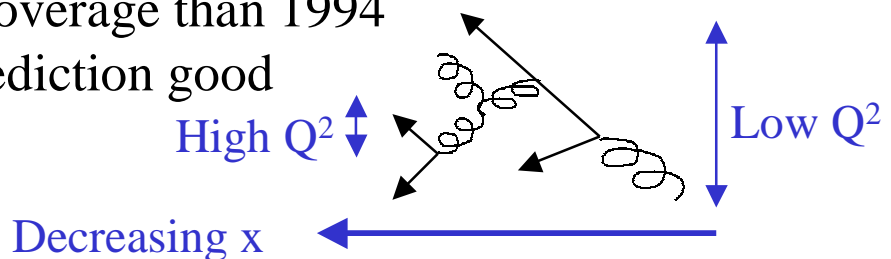
Decreasing x



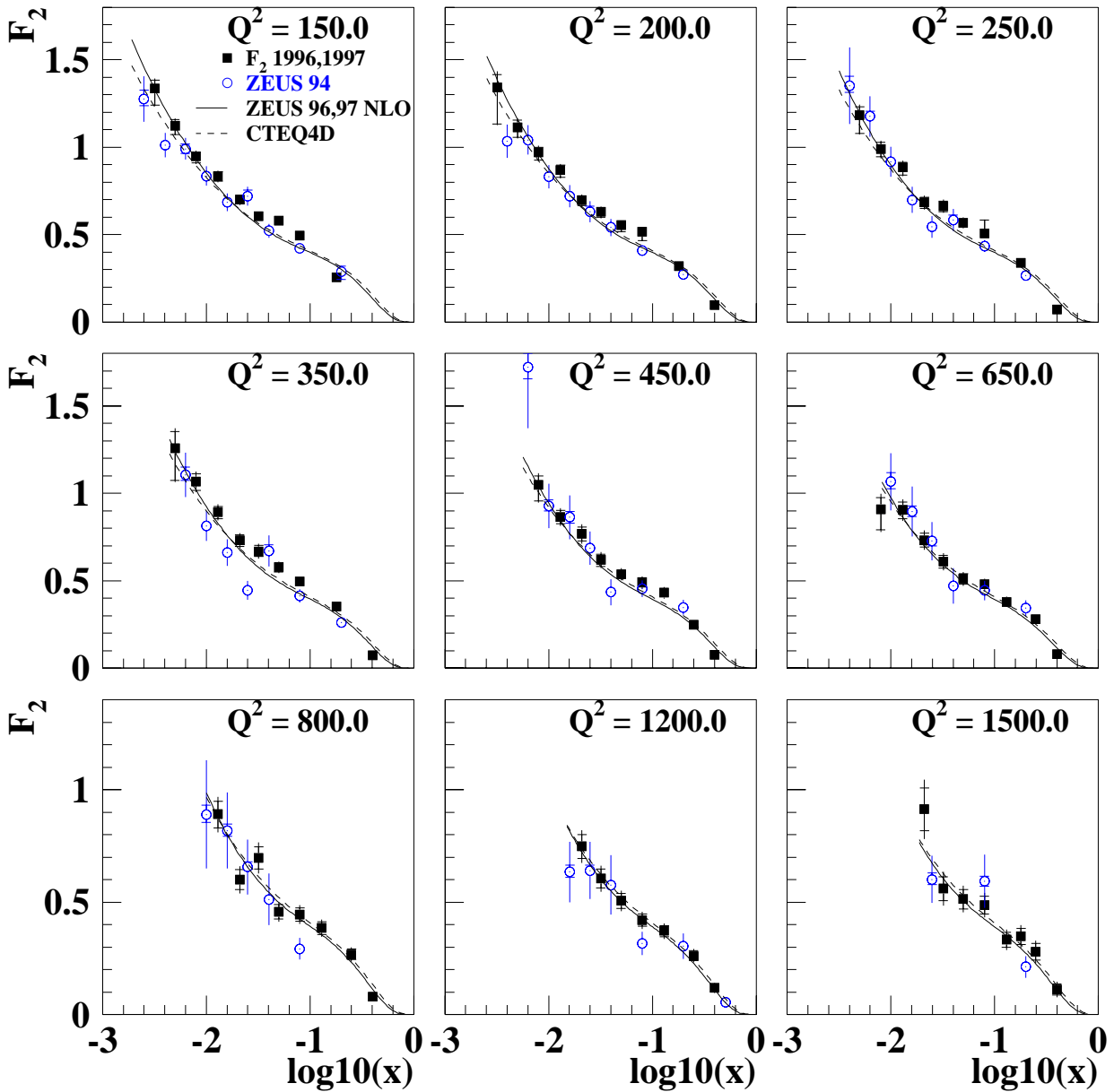
Low Q^2 F_2



- Strong rise at low x seen at very low Q^2 !
- Lower Q^2 coverage than 1994
- DGLAP prediction good

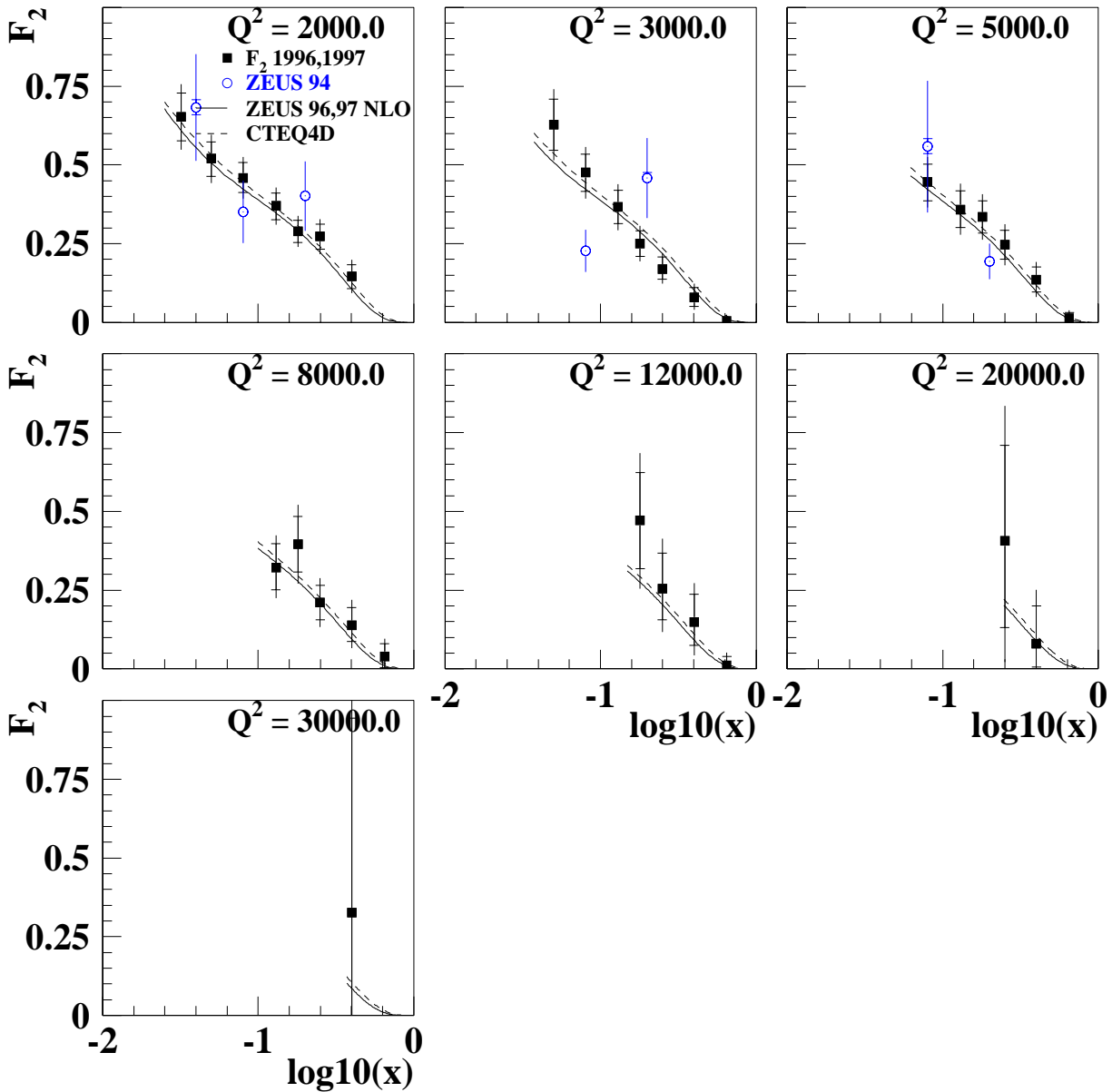


High Q^2 F_2



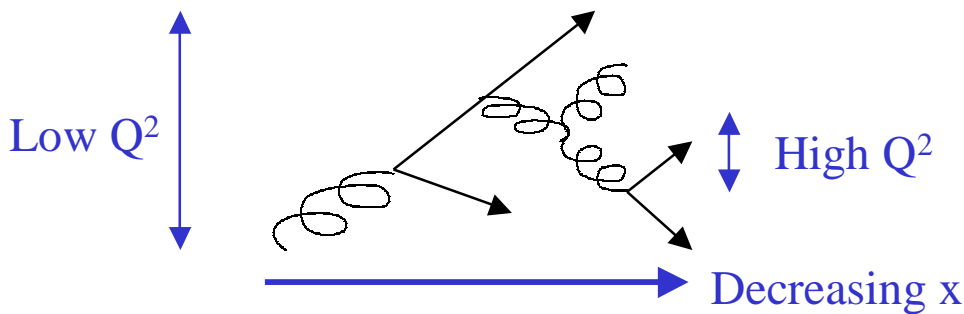
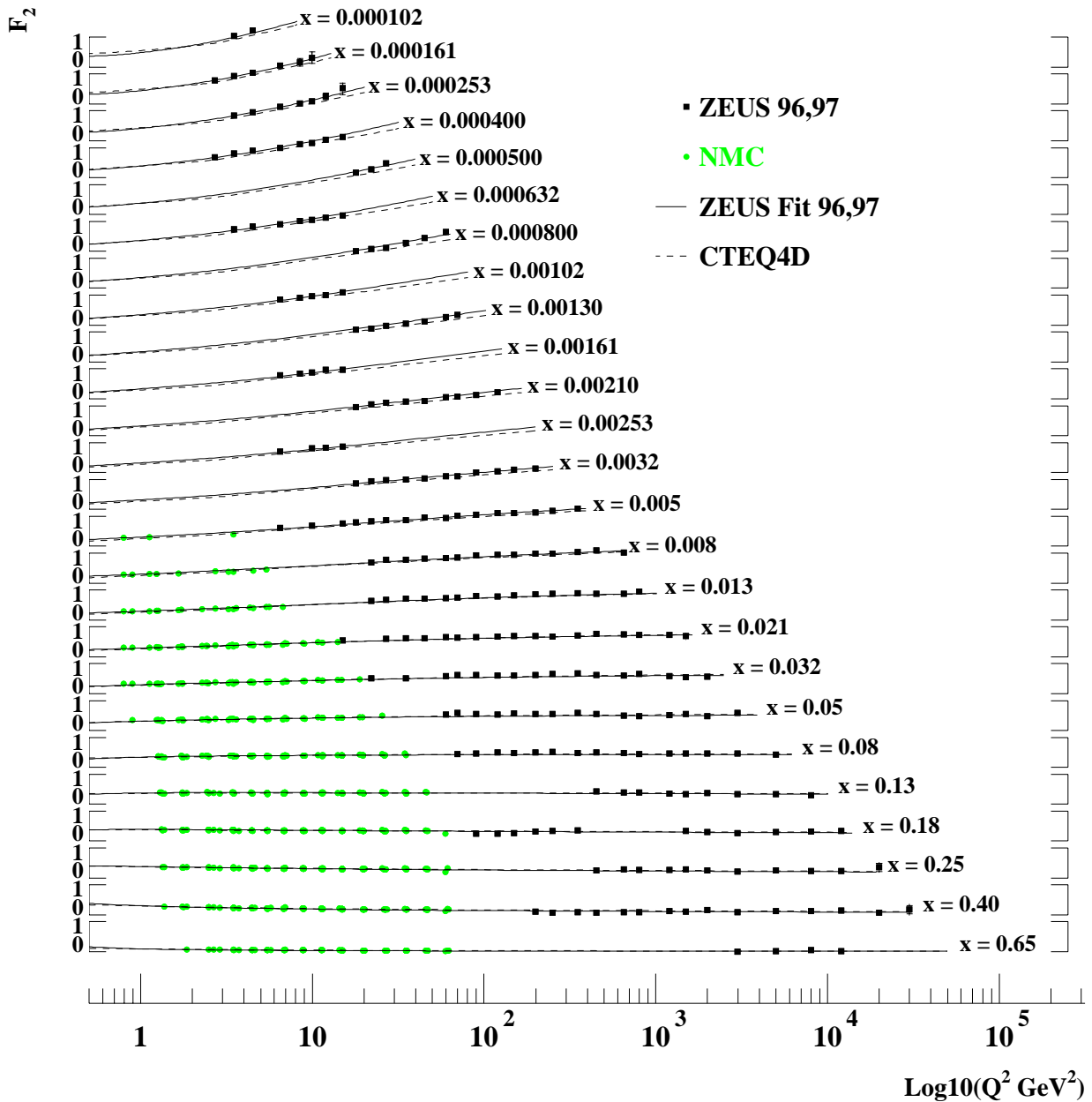
- Finer binning than 1994 at high Q^2
- Detector understanding (systematic error) is the dominant error up to $Q^2=800 \text{ GeV}^2$

Very High Q^2 F_2

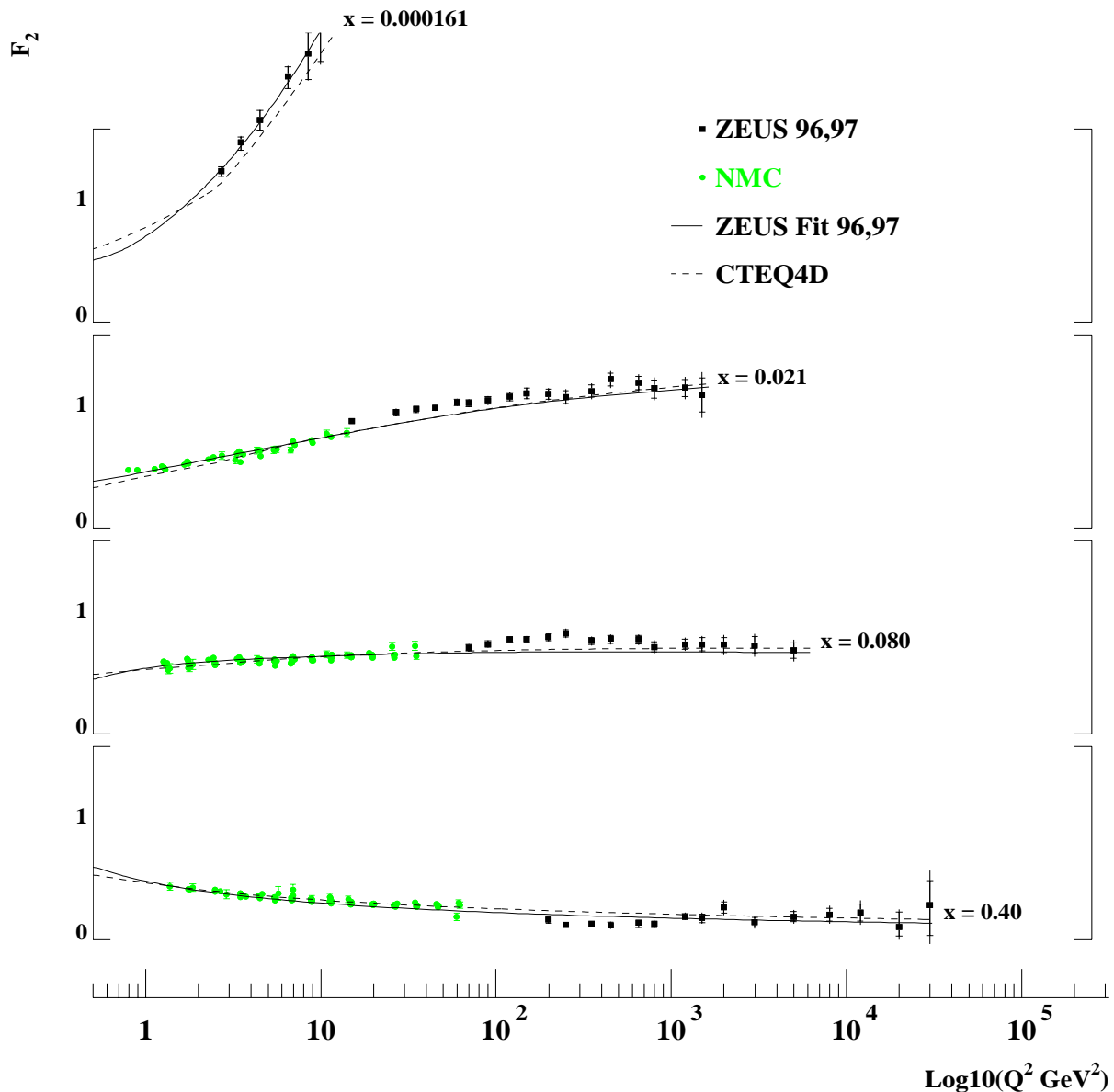


- Strong rise in F_2 seen at our highest Q^2
 - Over 4 orders of magnitude in Q^2
- DGLAP prediction good over whole range
- Increased coverage by an order of magnitude from 1994

F_2 vs. Q^2



F_2 vs. Q_2^2 (Few x bins)



- High x , F_2 falls with Q^2
 - valance quarks contribution falls
 - gluon small at high x
- Low x , F_2 rises strongly with Q^2
 - Strong gluon \rightarrow sea quark contribution

Conclusions

- F_2 is a measure of the quark structure of the proton.
- It has been measured from 1996 and 1997 ZEUS data from e+p collisions.
 - $1.5 < Q^2 < 30000 \text{ GeV}^2$
 - Wider Coverage at high and low Q^2 compared to 1994
 - $2(10^{-5}) < x < 0.65$
- F_2 shows a strong rise with decreasing x
 - Gluon quark cascade produces more quarks at lower momenta
 - Rise is seen at very low Q^2 (not expected!)
- The rise in F_2 becomes stronger with increasing Q^2
 - Parton splitting increases at smaller distances
- F_2 is described well by QCD DGLAP evolution over the full kinematic range covered.
- Dominant errors on F_2
 - Detector understanding for $Q^2 < 800 \text{ GeV}^2$
 - Event statistics for $Q^2 < 800 \text{ GeV}^2$

Future of ZEUS Structure Functions

$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2}(x, Q^2) = \frac{2\pi\alpha^2}{xQ^4} \left[(1 + (1-y)^2) F_2(x, Q^2) - y^2 F_L(x, Q^2) \mp (1 - (1-y)^2) xF_3(x, Q^2) \right]$$

- e+ and e- running to measure $x F_3$
 - ZEUS has enough e+ data (used in this analysis)
 - HERA 70.92 pb⁻¹ , ZEUS 48.08 pb⁻¹
 - Plans are forming to get additional e- data
 - HERA 27.37 pb⁻¹ ZEUS 17.61 pb⁻¹
- Lower P beam energy to measure F_L
 - $sxy = Q^2$ Center of mass energy = \sqrt{s}
 - Change \sqrt{s} and measure at two different y for the same x and Q^2
 - After 2000 Luminosity upgrade
- Continue high $Q^2 > 800 \text{ GeV}^2$ Analysis
- Longitudinal polarized electron beam
- Far future?...Polarized proton beams to measure polarized proton structure functions.