## Find a signal

$$
\begin{aligned}
\mathbf{N}_{\text {obs }}\left(\mathbf{p}_{\mathbf{T}}, \phi, \eta, \mathbf{z}\right)= & B+\sigma A \epsilon \int \mathcal{L} d t \\
= & B\left(p_{T}, \phi, \eta, z, \mathcal{L}\right)+ \\
& \int \sigma\left(p_{T}^{0}, \phi^{0}, \eta^{0}\right) \times \mathcal{L}(t, z) \times A\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0}\right) \times \\
& R\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0} ; p_{T}, \phi, \eta, z, \mathcal{L}\right) \times \\
& \epsilon\left(p_{T}, \phi, \eta, z, \mathcal{L}\right) d p_{T}^{0} d \phi^{0} d \eta^{0} d t
\end{aligned}
$$

$\sigma\left(X_{\alpha} \ldots\right)$ is the true cross section as a function of true variables $X_{a} \ldots$
$\mathcal{L}(t, z)$ is the luminosity as a function of time.
$A\left(X_{\alpha} ..\right)$ is the geometrical Acceptance as a function of true variables.
$R\left(X_{\alpha \ldots} \ldots, \mathcal{L} ; X_{a} \ldots\right)$ is the Resolution function which smears true $X_{\alpha} \ldots$ to detected $X_{a} \ldots$
$\epsilon\left(X_{a} \ldots, \mathcal{L}\right)$ is the probability that a particle is actually detected by a physical detector.
$B\left(X_{a} \ldots, \mathcal{L}\right)$ is the background

## Look for Z's In our detector


$\square$ Expect 2
"electrons"
$\square \mathrm{PT} \sim M Z / 2$
$\square$ Calorimeter energy
$\square$ Shower looks like an electron
$\square$ Track momentum

- Tracks come from the same point

7/12/11

## $Z \rightarrow$ ee

ET scale: 168 GeV


This ZO has high transverse momentum!

10/30/2009 - Louisville

Electrons have a track and EM energy
A $\mathbf{Z}^{0} \boldsymbol{\rightarrow} \mathrm{e}^{+} \mathrm{e}^{-}+\mathrm{jjevent}$
Jets have many tracks and EM+Hadron energy
$M_{e e} \sim 91 \mathrm{GeV} / \mathrm{c}^{2}$


End view

## What do we expect

$$
\begin{gathered}
\frac{d \sigma\left(f \bar{f} \rightarrow e^{-} e^{+}\right)}{d \cos \theta^{*}}=(1 / 128 \pi \hat{s})\left[\left(\left|A_{L L}\right|^{2}+\left|A_{R R}\right|^{2}\right)\left(1+\cos \theta^{*}\right)^{2}\right. \\
\left.+\left(\left|A_{L R}\right|^{2}+\left|A_{R L}\right|^{2}\right)\left(1-\cos \theta^{*}\right)^{2}\right] \\
\frac{d \sigma(g g \rightarrow g g)}{d \cos \theta^{*}} \propto \frac{1}{\sin ^{4} \frac{\theta^{2}}{2}}
\end{gathered}
$$

The leptons are more isotropic than the backgrounds

## Jacobian peak



## Now let's do some data analysis

$\square$ I happen to have a di-electron root file lying around.

## Invariant mass of raw electron pairs



## After quality cuts



[^0]
## Phi



## Detector $\eta$



## Detector and physics coordinates

$\square$ Detector is designed around $z=0$
$\square$ The beam is actually $25-30 \mathrm{~cm}$ long

## PT of leptons


$\mathrm{p}_{\mathrm{T}}, \mathrm{GeV}$

## Z position of vertex in cm



## The signal

## Invariant mass - Z candidates(All)



## Luminosity

$$
\begin{aligned}
N_{o b s}\left(p_{T}, \phi, \eta, z\right)= & B+\sigma A \epsilon \int \mathcal{L} d t \\
= & B\left(p_{T}, \phi, \eta, z, \mathcal{L}\right)+ \\
& \int \sigma\left(p_{T}^{0}, \phi^{0}, \eta^{0}\right) \times \mathcal{L}(t, z) \times A\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0}\right) \times \\
& R\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0} ; p_{T}, \phi, \eta, z, \mathcal{L}\right) \times \\
& \epsilon\left(p_{T}, \phi, \eta, z, \mathcal{L}\right) d p_{T}^{0} d \phi^{0} d \eta^{0} d t
\end{aligned}
$$

$\sigma\left(X_{\alpha} \ldots\right)$ is the true cross section as a function of true variables $X_{a} \ldots$.
$\mathcal{L}(t, z)$ is the luminosity as a function of time.
$A\left(X_{\alpha} ..\right)$ is the geometrical Acceptance as a function of true variables.
$R\left(X_{\alpha} \ldots, \mathcal{L} ; X_{a} \ldots\right)$ is the Resolution function which smears true $X_{\alpha} \ldots$ to detected $X_{a} \ldots$
$\epsilon\left(X_{a} \ldots, \mathcal{L}\right)$ is the probability that a particle is actually detected by a physical detector.
$B\left(X_{a} \ldots, \mathcal{L}\right)$ is the background

## Luminosity

Luminosity is a measure of how often protons/antiprotons get close enough to interact

$$
L=f \frac{n_{1} n_{2}}{4 \pi s_{x} s_{y}}
$$


$\mathrm{f}=$ beam crossing frequency 396 nsec
$\mathrm{n}=$ protons/bunch $10^{11}$
$\mathrm{s}=$ transverse beam size $=0.0001 \mathrm{~m}$
$\mathrm{L} \sim 10^{32}$ crossings $/ \mathrm{cm}^{2} / \mathrm{sec}$

## Typical Cross Sections

Total proton/antiproton cross section is $7 \times 10^{-30} \mathrm{~m}^{2}$

Unit of Barns $(b)=10^{-28} \mathrm{~m}^{2}$

$$
\sigma(p \bar{p} \Rightarrow X) \approx 70 m b
$$

Run II L $\sim 10^{32}$ crossings $/ \mathrm{cm}^{2} / \mathrm{sec}$
$\mathrm{N} / \mathrm{sec} \sim \sigma \mathrm{L}=7 \times 10^{6} / \mathrm{sec}$
$>3$ interactions per beam crossing!

Cross Section for top production:
$\sigma(p \bar{p} \Longrightarrow t \bar{t}+X) \approx 5 p b$
This is around $1 / 10^{10}$ of total
$\mathrm{N} / \mathrm{sec} \sim \sigma \mathrm{L}=5 \times 10^{-4} / \mathrm{sec}$

A couple were created/hour but we only saw a small \%

## Luminosity

Rate $=\int \sigma_{\text {int }}$

pbars

$$
\begin{array}{r}
\mathrm{P}_{\text {int }}=\mathrm{N}_{\text {prot }} \sigma_{\text {int }} / \mathrm{A} \\
L=\mathrm{f}_{\text {rev }} \mathrm{N}_{\text {pbar }} \mathrm{N}_{\text {prot }} / \mathrm{A}
\end{array}
$$

## Not easy to do this

$\square$ Hard to measure currents accurately
$\square$ the beam density depends on the intrinsic beam lengths and widths and on the local beam optics.

$$
\begin{aligned}
& L=2 \mathrm{f}_{\mathrm{rev}} \iiint \int \rho_{1} \rho_{2} \mathrm{dx} \mathrm{dy} \mathrm{dz} \mathrm{~d}(\mathrm{ct}) \\
& \rho(x, y, z, c t)=\frac{N_{1}}{\sqrt{2 \pi \sigma_{x}}} \quad \exp \left[-(x+\Delta x / 2)^{2} / 2 \sigma_{x}{ }^{2}\right] \\
& \sqrt{\frac{1}{2 \pi \sigma_{y}}} \exp \left[-(y+\Delta y / 2)^{2} / 2 \sigma_{y}{ }^{2}\right] \\
& \frac{1}{\sqrt{2 \pi \sigma_{z}}} \exp \left[-(\underset{\text { СтеQ2011 }}{(z+c t-C t})^{2} / 2 \sigma_{z}{ }^{2}\right]
\end{aligned}
$$

You can measure the beam spot size as a function of distance along the beam to get $\sigma_{x}(z), \sigma_{y}(z)$
$\square$ Also need estimates of beam current
$\square \sigma_{z}$ can be estimated from beam timing or $N(z)$
$\square$ Estimates are good to ~10\% at the Tevatron.
$\square$ The beam position and size vary with time as the beam "heats up" and "cools down"


7/12/11

## $\mathscr{L}(z)$ distribution for beam collisions



## Alternate method

$\square$ Measure the total inelastic proton antiproton cross section in a special run.
$\square$ Count inelastic collisions as you run your regular experiment.
$\square$ Use those to get an estimate of your integrated luminosity.

## DO Luminosity System

24
proton direction

LM


Count fraction, $P(0)$, of beam crossings without inelastic interactions.

## Method

$\square$ Measure $\mu$, the average number of interactions/crossing
$\square$ Apply acceptance and efficiency corrections and then compare to the inelastic pbar p total cross section to get the luminosity integrated over the 396 nsec crossing time.
$\square$ Typical $\mu$ are 1-10 at the Tevatron

## Interactions/crossing

Typical running at the Tevatron


## Luminosity errors

$\square$ The luminosity measurement has 3 sources of error

- The inelastic cross section is only known to 3-4\%
$\square$ The luminosity detector is not perfect... 3-4\%
$\square$ Bookkeeping - you have to match the luminosity data stream up with your real data.
- What happens if you lose a data tape? Flag data as bad?
- We track this by matching "luminosity blocks", 1 minute periods of data.


## So far we

$\square$ Have a signal
$\square$ Have a normalization
$\square$ How do we find backgrounds?
$\square$ How do we correct for detector efficiency and resolution?
$\square$ How do we correct for Acceptance?

## Estimating backgrounds

$\square$ Simulate and subtract
$\square$ Requires good modeling
$\square$ Hard to normalize absolutely
$\square$ Simulate and fit
$\square$ Template method
$\square$ Vary cuts - matrix method

## Simulated backgrounds

Muon channel



## Background from template fit

## MC Template Fit to $\mathbf{Z} \gamma$ Data

Do this for $Z+\gamma$ signal
Neural network discriminant for EM vs jet


## Vary cuts - matrix method

## Vary cuts

$$
\begin{aligned}
& N_{\text {loose }}=\epsilon_{\text {loose }} N_{e}+f_{\text {loose }} N_{Q C D} \\
& N_{\text {tight }}=\epsilon_{\text {tight }} N_{e}+f_{\text {tight }} N_{Q C D}
\end{aligned}
$$

$\epsilon$ is efficiency for signal $\approx 70-100 \%$.
$f$ is probability to create a fake $\approx 0.1-10 \%$

$$
N_{e} \approx \frac{f_{\text {tight }} N_{\text {loose }}-f_{\text {loose }} N_{\text {tight }}}{\epsilon_{\text {loose }} f_{\text {tight }}-\epsilon_{\text {tight }} f_{\text {loose }}}
$$

## Efficiency

$$
\begin{aligned}
\mathbf{N}_{\text {obs }}\left(\mathbf{p}_{\mathbf{T}}, \phi, \eta, \mathbf{z}\right)= & B+\sigma A \in \int \mathcal{L} \mathbf{d} \mathbf{t} \\
= & B\left(p_{T}, \phi, \eta, z, \mathcal{L}\right)+ \\
& \int \sigma\left(p_{T}^{0}, \phi^{0}, \eta^{0}\right) \times \mathcal{L}(\mathbf{t}, \mathbf{z}) \times A\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0}\right) \times \\
& R\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0} ; p_{T}, \phi, \eta, z, \mathcal{L}\right) \times \\
& \epsilon\left(\mathbf{p}_{\mathrm{T}}, \phi, \eta, \mathbf{z}, \mathcal{L}\right) d p_{T}^{0} d \phi^{0} d \eta^{0} d t
\end{aligned}
$$

$\sigma\left(X_{\alpha} \ldots\right)$ is the true cross section as a function of true variables $X_{a} \ldots$
$\mathcal{L}(t, z)$ is the luminosity as a function of time.
$A\left(X_{\alpha} ..\right)$ is the geometrical Acceptance as a function of true variables.
$R\left(X_{\alpha} \ldots, \mathcal{L} ; X_{a \ldots} \ldots\right.$ is the Resolution function which smears true $X_{\alpha} \ldots$ to detected $X_{a} \ldots$
$\epsilon\left(\mathbf{X}_{\mathbf{a}} \ldots, \mathcal{L}\right)$ is the probability that a particle is actually detected by a physical detector.
$B\left(X_{a} \ldots, \mathcal{L}\right)$ is the background

## Efficiency

$\square$ First pass
$\square$ Define an "acceptance region" where you would expect your detector to work

- (say $|\eta|<3, \mathrm{p}_{\mathrm{T}}>10 \mathrm{GeV}$ )
- Take a particle level MC like pythia
$\square$ Reweight kinematics to reflect best knowledge
$\square$ Trace particles through detector
$\square$ Overlay noise and interactions from other events
$\square$ Count how often you reconstruct the particle


## The raw simulation isn't good enough

$\square$ Use scale factors

Find two ways of measuring something
Compare the two in data and in simulation
Correct by the ratio

$$
\epsilon\left(p_{T}, z\right)=\frac{(P \cdot T)_{d a t a} / T_{d a t a}}{(P \cdot T)_{s i m} / T_{s i m}} \epsilon_{s i m}\left(p_{T}, z\right)
$$

## Tag and probe


$\square$ Find a really good electron (TAG)
$\square$ Find a track which might be from a $Z$ decay (PROBE)
$\square$ Did you find a shower
$\square$ Do the same for a shower (PROBE) that might have a track

Total efficiency : caleta_vtxz


## Typical electron efficiencies

## Efficiency versus Instantaneous Luminosity



## Resolutions

Step 1 - Simulate$\square$ Step 2 - Cry in frustration
$\square$ Step 3 - use the data
$\square$ Tune to the $Z$ peak width or
$\square$ Compare energy of back to back jets

$$
A=\frac{E_{T}^{1}-E_{T}^{2}}{E_{T}^{1}+E_{T}^{2}}
$$



## Check the tuned full detector simulation



## Z v†x check



## Curvature check


$\square$ Ok, looks like we understand our efficiencies and resolutions pretty well in our simulation

Correct data for backgrounds and efficiency
Now on to unfolding the resolution effects

## Unfolding the resolution effects

$$
\begin{aligned}
& \mathbf{N}_{\text {obs }}\left(\mathbf{p}_{\mathbf{T}}, \phi, \eta, \mathbf{z}\right)= B+\sigma A \in \int \mathcal{L} \mathbf{d t} \\
&= B\left(p_{T}, \phi, \eta, z, \mathcal{L}\right)+ \\
& \int \sigma\left(p_{T}^{0}, \phi^{0}, \eta^{0}\right) \times \mathcal{L}(\mathbf{t}, \mathbf{z}) \times A\left(p_{T}^{0}, \phi^{0}, \eta^{0}, z^{0}\right) \times \\
& \quad \mathbf{R}\left(\mathbf{p}_{\mathbf{T}}^{\mathbf{0}}, \phi^{\mathbf{0}}, \eta^{\mathbf{0}}, \mathbf{z}^{\mathbf{0}} ; \mathbf{p}_{\mathbf{T}}, \phi, \eta, \mathbf{z}, \mathcal{L}\right) \times \\
& \epsilon\left(\mathbf{p}_{\mathrm{T}}, \phi, \eta, \mathbf{z}, \mathcal{L}\right) d p_{T}^{0} d \phi^{0} d \eta^{0} d t
\end{aligned}
$$

$\sigma\left(X_{\alpha} \ldots\right)$ is the true cross section as a function of true variables $X_{a} \ldots$
$\mathcal{L}(t, z)$ is the luminosity as a function of time.
$A\left(X_{\alpha} ..\right)$ is the geometrical Acceptance as a function of true variables.
$\mathbf{R}\left(\mathbf{X}_{\alpha} \ldots, \mathcal{L} ; \mathbf{X}_{\mathbf{a}} \ldots\right)$ is the Resolution function which smears true $X_{\alpha} \ldots$ to detected $X_{a} \ldots$
$\epsilon\left(\mathbf{X}_{\mathbf{a}} \ldots, \mathcal{L}\right)$ is the probability that a particle is actually detected by a physical detector.
$B\left(X_{a} \ldots, \mathcal{L}\right)$ is the background

## Resolution Smearing


(a)

(b)

## Unfolding method 1 - matrix

$\square$ Invert the resolution matrix

$$
N_{a}=B_{a}+\epsilon_{a} R_{a \alpha} A_{\alpha} \sigma_{\alpha} \int \mathcal{L} d t
$$

Invert

$$
\sigma_{\alpha}=\frac{1}{A_{\alpha}} R^{-1} \alpha a \frac{\left(N_{a}-B_{a}\right)}{\epsilon_{a}} \int \mathcal{L} d t
$$

$\square$ Common to regularize the matrix
$\square$ Improved stability
$\square$ Bias towards smoothing the function

## Smearing Matrix $R\left(X_{\alpha}->X_{a}\right)$



## Unfolding method 2 - Ansatz

$\square$ Convolute a trial unsmeared function with the resolution as a function of (pt, $\eta$ )

$$
f\left(p_{T}, \eta\right)=N_{0}\left(\frac{p_{T}}{100 \mathrm{GeV} / \mathrm{c}}\right)^{-\alpha}\left(1-\frac{2 p_{T} \cosh \left(y_{\min }\right)}{\sqrt{s}}\right)^{\beta} \exp \left(-\gamma p_{T}\right) .
$$

$\square$ Fit this convoluted "smeared" function to your data
$\square$ Correct data by the ratio of the unsmeared to smeared "ansatz" function.

$$
R\left(X_{\alpha} \rightarrow X_{a}\right)=\delta_{\alpha a} \frac{\operatorname{Smeared}\left(X_{a}\right)}{\operatorname{Unsmeared}\left(X_{\alpha}\right)}
$$

$\square$ This works if a simple functional form can fit your data well and the resolution function is well understood.

## Ansatz fit $\quad R\left(X_{\alpha} \rightarrow X_{a}\right)=\delta_{\alpha a} \frac{\operatorname{Smeared}\left(X_{a}\right)}{\operatorname{Unsmeared}\left(X_{\alpha}\right)}$




## Now for the final step: Acceptance

$\square$ Correct from the region we cover to the rest of phase space
$\square$ Why do this instead of including in the "efficiency"?
$\square$ This correction is big - factor of 3 for $Z$ production $\varepsilon \sim 40 \%$, $\mathrm{A} \sim 30 \%$
$\square$ Only depends on generator level quantities

- Can use the best NNLO simulations


## DO $1 \mathrm{fb}^{-1}$ sample



## Z Rapidity raw and generated



We actually define acceptance as $\left|\eta_{\mathrm{e}}\right|<1$ or $1.5<\left|\eta_{\mathrm{e}}\right|<2.4$ and $\mathrm{p}_{\mathrm{T}}>25 \mathrm{GeV}$

Cut cross section in electron rapidity and pt


## Fully corrected normalized distributions



## PHYSICAL REVIEW D 76, 012003 (2007)



## Theory



The Drell-Yan Process


Higher order QCD is important
$\mathrm{gq} \rightarrow \mathrm{Zq}, \mathrm{qq} \rightarrow \mathrm{Zg}$ etc.
Cross section*B(Z $\rightarrow$ ee)
LO ~ 180 pb
NLO ~ 250 pb
NNLO ~ 260 pb
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## $Z \rightarrow$ ee Errors

| Statistical | $0.62 \%$ |
| :--- | :--- |
| Preselection efficiency | $0.85 \%$ |
| Radiative corrections | $<0.5 \%$ |
| ID eff stat error | $0.4 \%$ |
| Tag-probe bias | $0.3 \%$ |
| Noise corrections | $0.22 \%$ |
| Vertex z | $0.6 \%$ |
| Cut variations | $1.5 \%$ |
| Total systematic error | $2.0 \%$ |
| PDF error on sigma(tot) <br> Luminosity | $6.1 \% \%-1.7 \%$ |

## More fun

$\square$ Radiative corrections
$\square$ Calibration
$\square$ Jet definitions
$\square$ B-tagging
$\square$ My list of things to check is around 80 items long


[^0]:    CTEQ2011 Schellman 7/12/11

