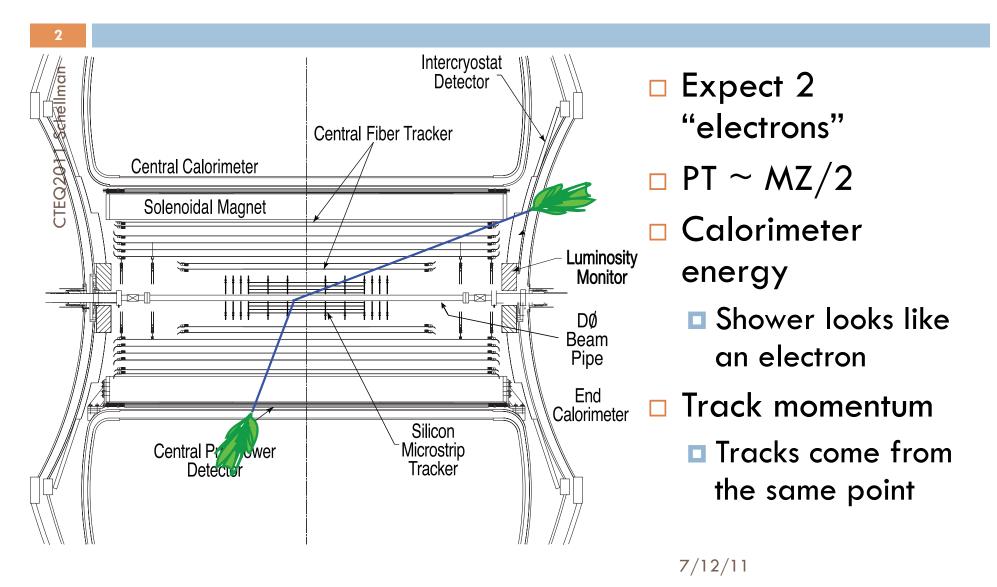
Find a signal

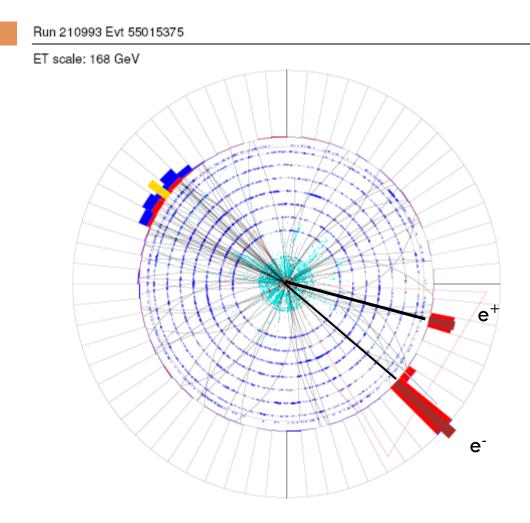
$$\begin{split} \mathbf{N_{obs}}(\mathbf{p_T}, \boldsymbol{\phi}, \boldsymbol{\eta}, \mathbf{z}) &= B + \sigma A \epsilon \int \mathcal{L} dt \\ &= B(p_T, \boldsymbol{\phi}, \boldsymbol{\eta}, z, \mathcal{L}) + \\ &\int \sigma(p_T^0, \boldsymbol{\phi}^0, \boldsymbol{\eta}^0) \times \mathcal{L}(t, z) \times A(p_T^0, \boldsymbol{\phi}^0, \boldsymbol{\eta}^0, z^0) \times \\ &R(p_T^0, \boldsymbol{\phi}^0, \boldsymbol{\eta}^0, z^0; p_T, \boldsymbol{\phi}, \boldsymbol{\eta}, z, \mathcal{L}) \times \\ &\epsilon(p_T, \boldsymbol{\phi}, \boldsymbol{\eta}, z, \mathcal{L}) dp_T^0 d\boldsymbol{\phi}^0 d\boldsymbol{\eta}^0 dt \end{split}$$

- $\sigma(X_{\alpha}...)$ is the true cross section as a function of true variables $X_{a}...$
- $\mathcal{L}(t,z)$ is the luminosity as a function of time.
- $A(X_{\alpha}..)$ is the geometrical Acceptance as a function of true variables.
- $R(X_{\alpha}...,\mathcal{L};X_{a}...)$ is the Resolution function which smears true $X_{\alpha}...$ to detected $X_{a}...$
- $\epsilon(X_a..., \mathcal{L})$ is the probability that a particle is actually detected by a physical detector.
- $B(X_a..., \mathcal{L})$ is the background

Look for Z's In our detector



$Z \rightarrow ee$



This ZO has high transverse momentum!

10/30/2009 – Louisville

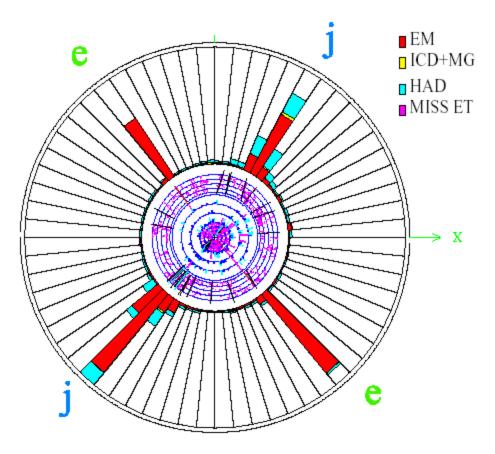
Electrons have a track and EM energy

A $Z^0 \rightarrow e^+e^- + jjevent$



 $M_{ee} \sim 91 \text{ GeV/c}^2$

Jets have many tracks and EM+Hadron energy



End view

ouisville

What do we expect

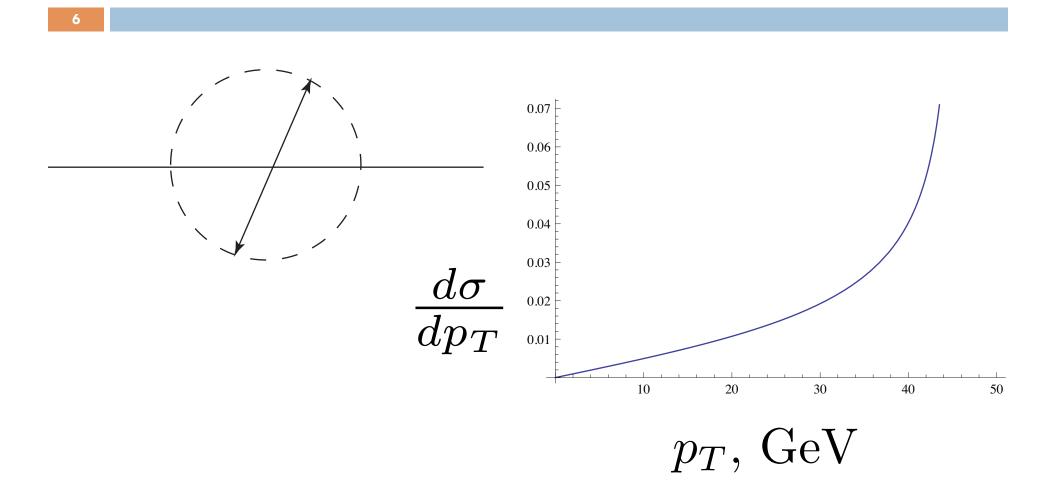
$$\frac{d\sigma(f\bar{f} \to e^- e^+)}{d\cos\theta^*} = (1/128\pi\hat{s})[(|A_{LL}|^2 + |A_{RR}|^2)(1 + \cos\theta^*)^2 + (|A_{LR}|^2 + |A_{RL}|^2)(1 - \cos\theta^*)^2] \quad .$$

$$\frac{d\sigma(gg \to gg)}{d\cos\theta^*} \propto \frac{1}{\sin^4\frac{\theta}{2}^*}$$

The leptons are more isotropic than the backgrounds

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Jacobian peak



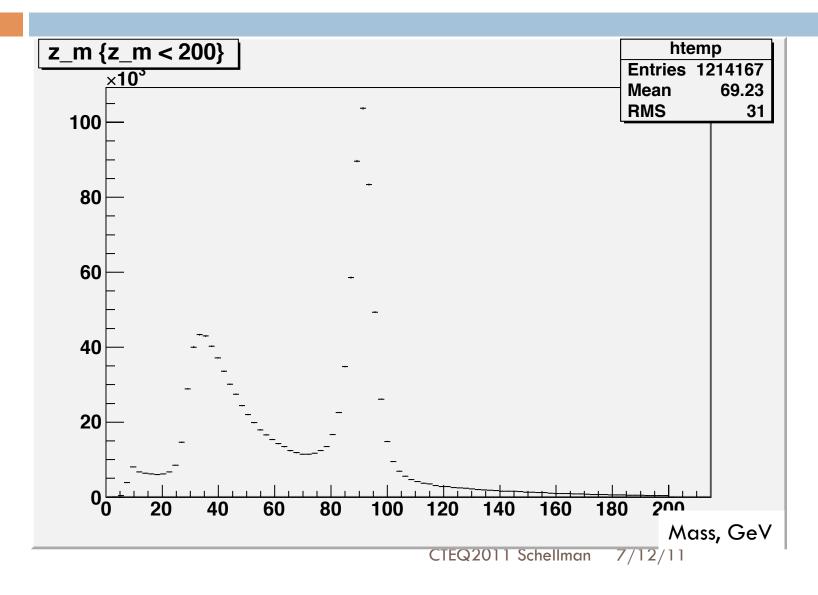
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Now let's do some data analysis

I happen to have a di-electron root file lying around.

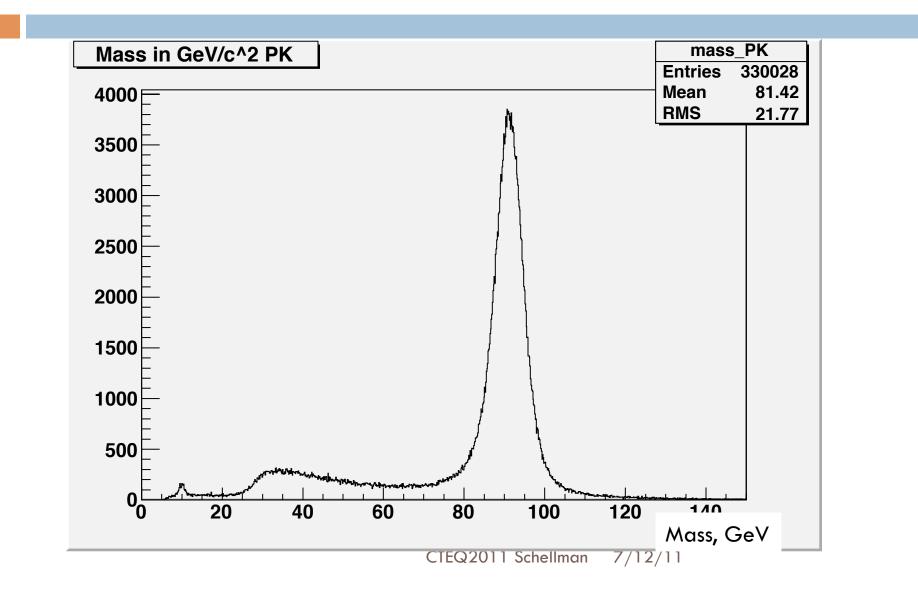
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Invariant mass of raw electron pairs



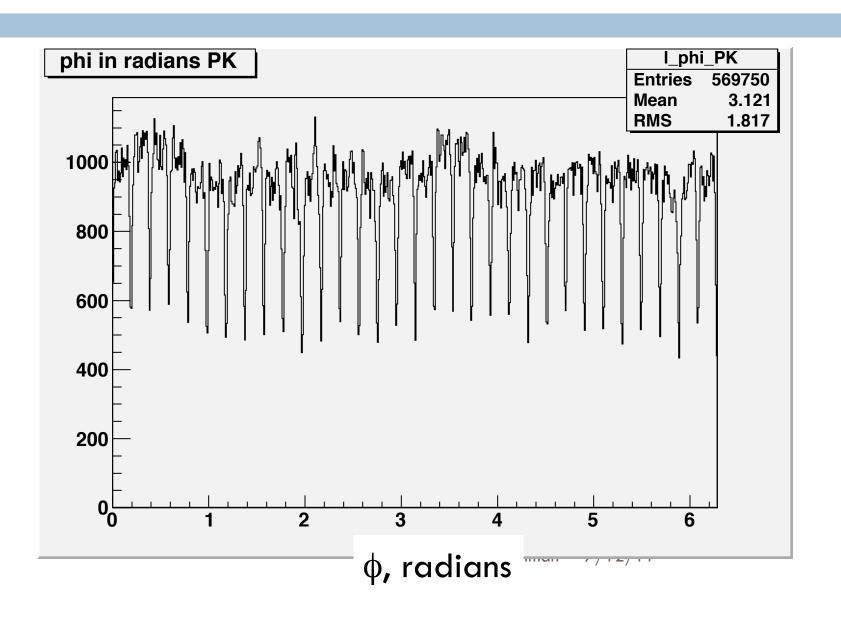
After quality cuts

9

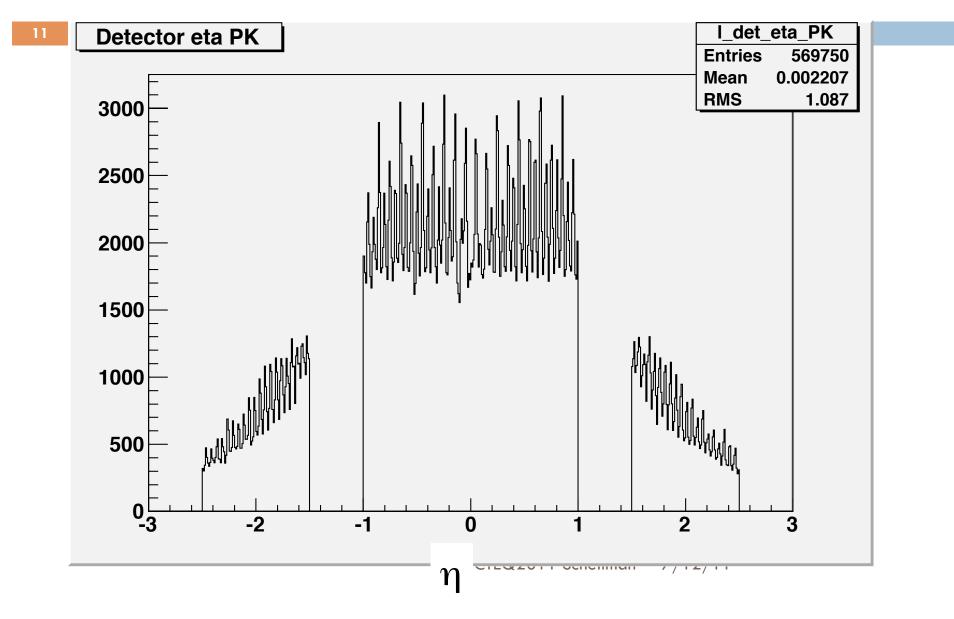


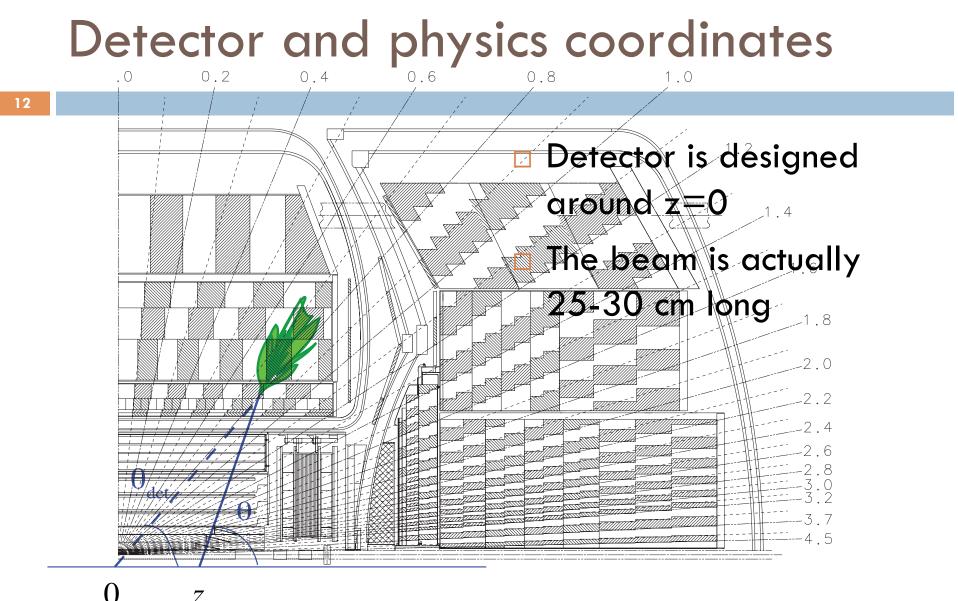
Phi

10



Detector η

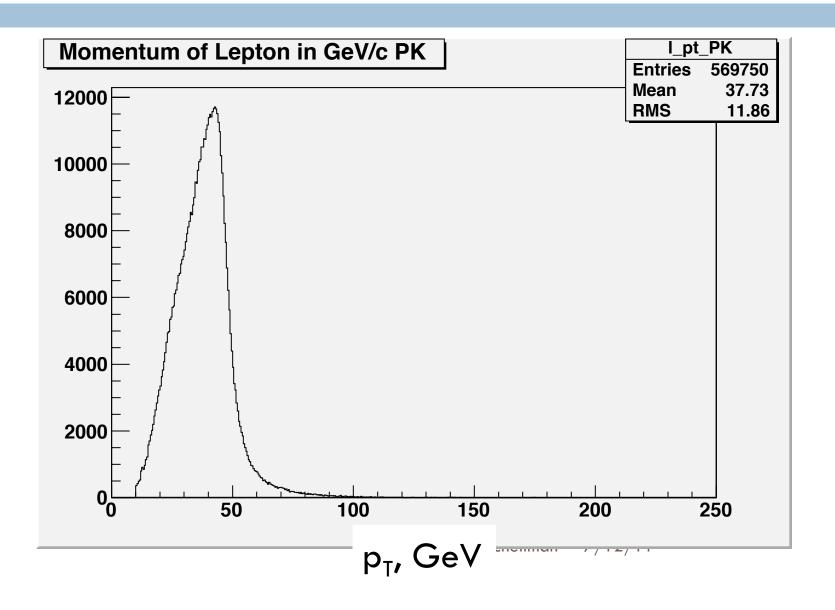




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Ζ.

PT of leptons

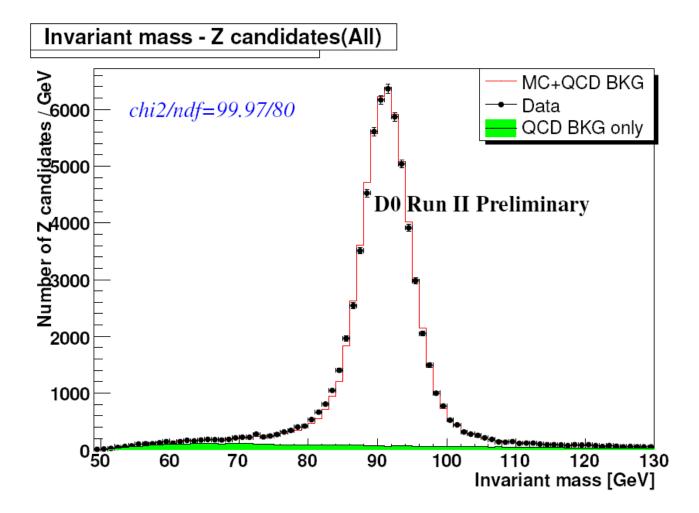


Z position of vertex in cm

I_z0 PK I z0 PK 569750 **Entries** 0.01386 Mean RMS 17.57 5000 4000 3000 2000 1000 -100 -60 -80 -40 -20 20 40 60 80 100 0 S A V I **UCHONING** z₀, cm

14

The signal



Luminosity

16

$$\begin{split} N_{obs}(p_T,\phi,\eta,z) &= B + \sigma A\epsilon \int \mathcal{L} dt \\ &= B(p_T,\phi,\eta,z,\mathcal{L}) + \\ &\int \sigma(p_T^0,\phi^0,\eta^0) \times \mathcal{L}(t,z) \times A(p_T^0,\phi^0,\eta^0,z^0) \times \\ &R(p_T^0,\phi^0,\eta^0,z^0;p_T,\phi,\eta,z,\mathcal{L}) \times \\ &\epsilon(p_T,\phi,\eta,z,\mathcal{L}) dp_T^0 d\phi^0 d\eta^0 dt \end{split}$$

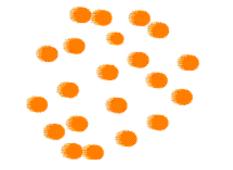
 $\sigma(X_{\alpha}...)$ is the true cross section as a function of true variables $X_{a}...$

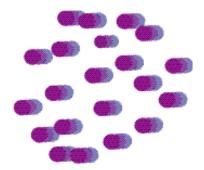
- $\mathcal{L}(t,z)$ is the luminosity as a function of time.
- $A(X_{\alpha}..)$ is the geometrical Acceptance as a function of true variables.
- $R(X_{\alpha}..., \mathcal{L}; X_{a}...)$ is the Resolution function which smears true $X_{\alpha}...$ to detected $X_{a}...$
- $\epsilon(X_a..., \mathcal{L})$ is the probability that a particle is actually detected by a physical detector.
- $B(X_a..., \mathcal{L})$ is the background CTEQ2011 Schellman 7/12/11

Luminosity

Luminosity is a measure of how often protons/antiprotons get close enough to interact

$$L = f \frac{n_1 n_2}{4\pi s_x s_y}$$





- f= beam crossing frequency 396 nsec
- n= protons/bunch 10¹¹
- s = transverse beam size = 0.0001 m
- $L \sim 10^{32} \text{ crossings/cm}^2/\text{sec}$

Typical Cross Sections

Total proton/antiproton cross section is $7x10^{-30} m^2$

Unit of Barns (b) = 10^{-28} m²

 $\sigma(p\overline{p} \Rightarrow X) \approx 70 \, mb$

Run II L $\sim 10^{32}\,crossings/cm^2/sec$

N/sec ~ $\sigma L = 7 \times 10^6$ /sec

> 3 interactions per beam crossing!

Cross Section for top production:

$$\sigma(p\overline{p} \Longrightarrow t\overline{t} + X) \approx 5\,p\,b$$

This is around $1/10^{10}$ of total

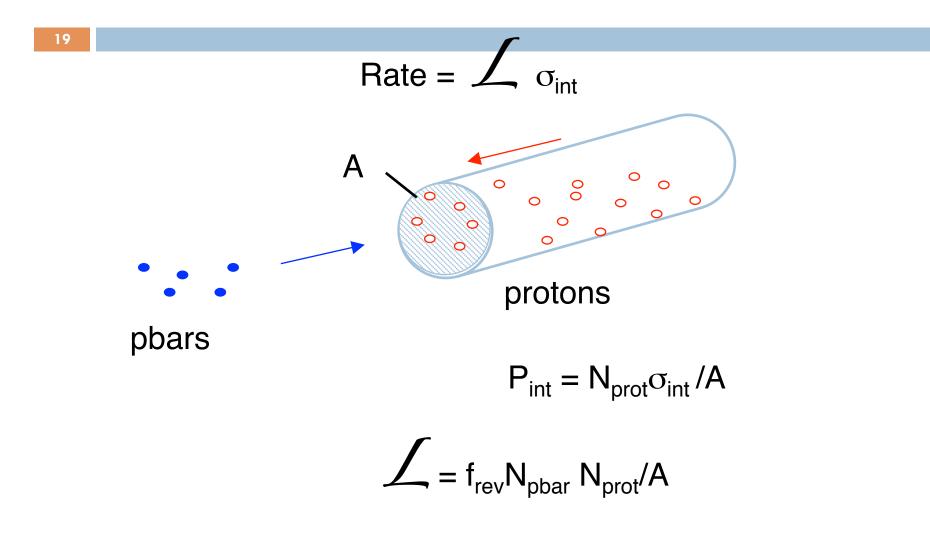
N/sec ~ $\sigma L = 5 \times 10^{-4}$ /sec

A couple were created/hour but we only saw a small %

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Luminosity



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Not easy to do this

- Hard to measure currents accurately
- the beam density depends on the intrinsic beam lengths and widths and on the local beam optics.

$$\int = 2 f_{rev} \iiint \rho_1 \rho_2 \, dx \, dy \, dz \, d(ct)$$

$$\rho(x,y,z,ct) = \frac{N_1}{\sqrt{2\pi} \sigma_x} \quad \exp[-(x+\Delta x/2)^2/2 \sigma_x^2]$$

$$\frac{1}{\sqrt{2\pi} \sigma_y} \quad \exp[-(y+\Delta y/2)^2/2 \sigma_y^2]$$

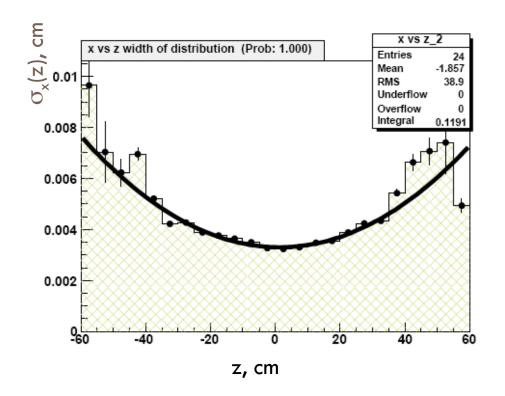
$$\frac{1}{\sqrt{2\pi} \sigma_z} \quad \exp[-(z+ct-ct_0)^2/2 \sigma_z^2]$$

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You can measure the beam spot size as a function of distance along the beam to get $\sigma_x(z)$, $\sigma_y(z)$

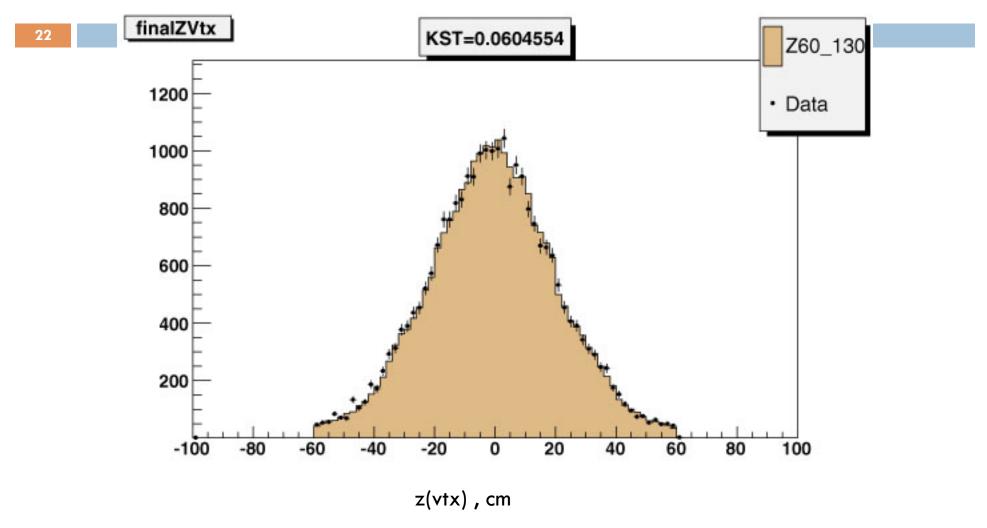
21

- Also need estimates of beam current
 - σ_z can be estimated from beam timing or N(z)
 - Estimates are good to
 ~10% at the Tevatron.
 - The beam position and size vary with time as the beam "heats up" and "cools down"



7/12/11

$\mathcal{L}(z)$ distribution for beam collisions

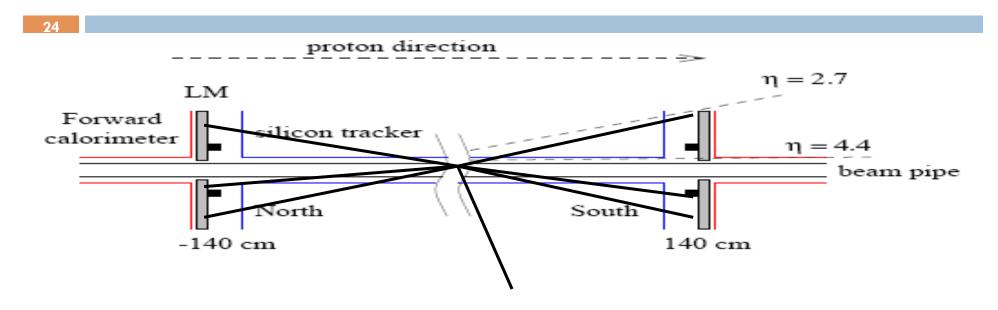


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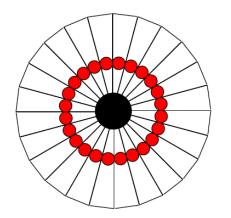
Alternate method

- 23
- Measure the total inelastic proton antiproton cross section in a special run.
- Count inelastic collisions as you run your regular experiment.
- Use those to get an estimate of your integrated luminosity.

DO Luminosity System



٠



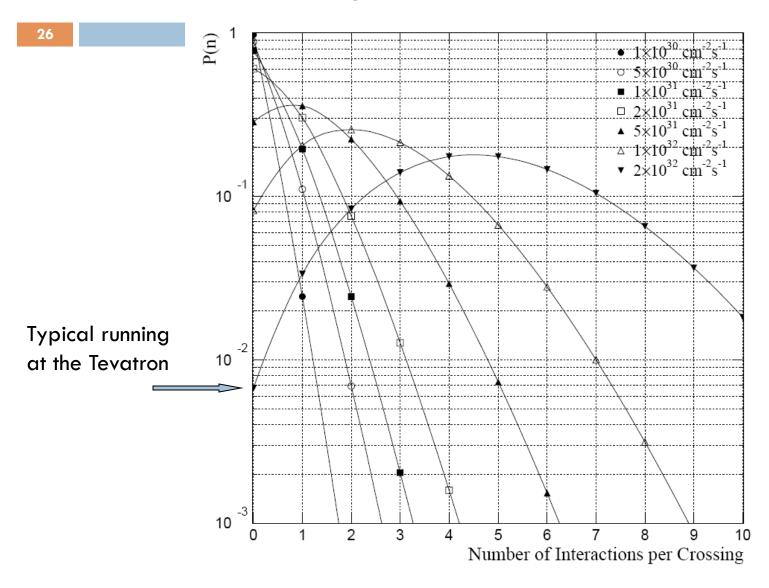
Count fraction, P(0), of beam crossings without inelastic interactions.

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Method

- \Box Measure μ , the average number of interactions/crossing
- Apply acceptance and efficiency corrections and then compare to the inelastic pbar p total cross section to get the luminosity integrated over the 396 nsec crossing time.
- $\hfill\square$ Typical μ are 1-10 at the Tevatron

Interactions/crossing



Luminosity errors

- 27
- □ The luminosity measurement has 3 sources of error
 - The inelastic cross section is only known to 3-4%
 - The luminosity detector is not perfect... 3-4%
 - Bookkeeping you have to match the luminosity data stream up with your real data.
 - What happens if you lose a data tape? Flag data as bad?
 - We track this by matching "luminosity blocks", 1 minute periods of data.

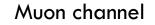
So far we

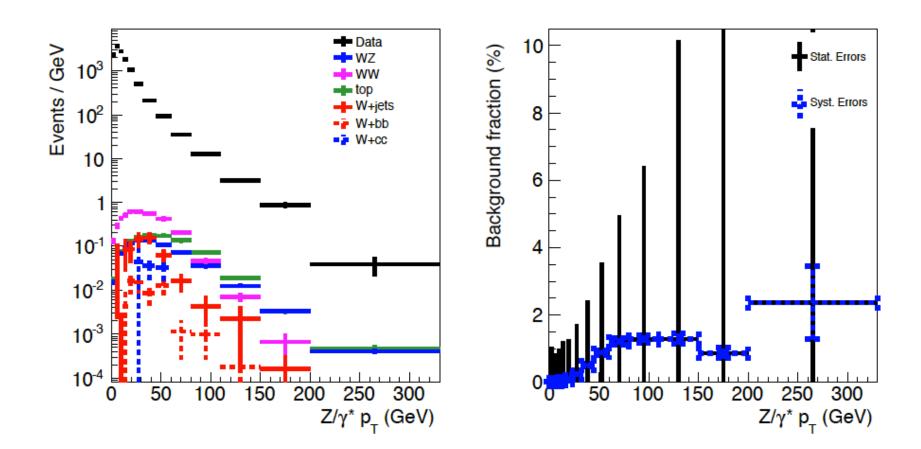
- Have a signal
- Have a normalization
- □ How do we find backgrounds?
- How do we correct for detector efficiency and resolution?
- □ How do we correct for Acceptance?

Estimating backgrounds

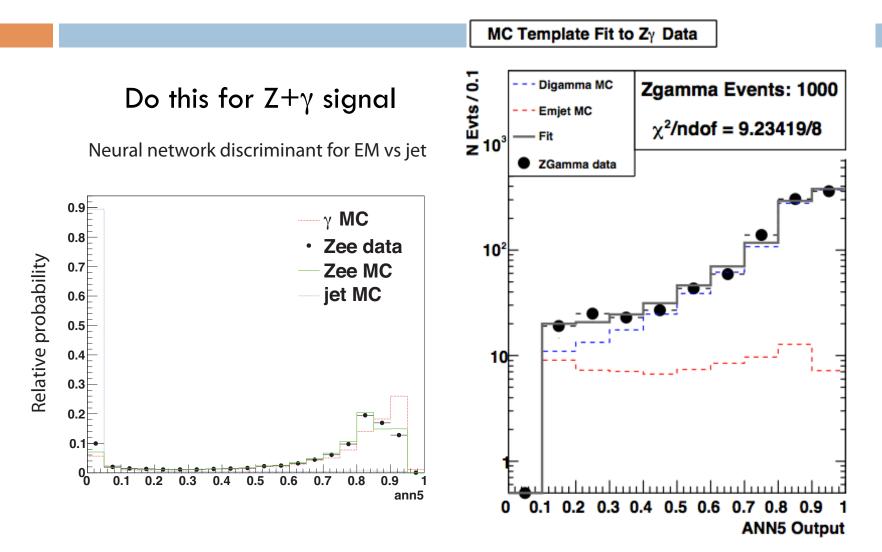
- Simulate and subtract
 - Requires good modeling
 - Hard to normalize absolutely
- Simulate and fit
 - Template method
- Vary cuts matrix method

Simulated backgrounds





Background from template fit



Vary cuts – matrix method

Vary cuts

$$N_{loose} = \epsilon_{loose} \ N_e + f_{loose} \ N_{QCD}$$

$$N_{tight} = \epsilon_{tight} N_e + f_{tight} N_{QCD}$$

 ϵ is efficiency for signal \approx 70-100%. f is probability to create a fake \approx 0.1-10%

$$N_e \approx \frac{f_{tight} N_{loose} - f_{loose} N_{tight}}{\epsilon_{loose} f_{tight} - \epsilon_{tight} f_{loose}}$$

Efficiency

33

$$\begin{split} \mathbf{N_{obs}}(\mathbf{p_T}, \phi, \eta, \mathbf{z}) &= B + \sigma A \epsilon \int \mathcal{L} \mathbf{dt} \\ &= B(p_T, \phi, \eta, z, \mathcal{L}) + \\ &\int \sigma(p_T^0, \phi^0, \eta^0) \times \mathcal{L}(\mathbf{t}, \mathbf{z}) \times A(p_T^0, \phi^0, \eta^0, z^0) \times \\ &R(p_T^0, \phi^0, \eta^0, z^0; p_T, \phi, \eta, z, \mathcal{L}) \times \\ &\epsilon(\mathbf{p_T}, \phi, \eta, \mathbf{z}, \mathcal{L}) dp_T^0 d\phi^0 d\eta^0 dt \end{split}$$

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- $R(X_{\alpha}...,\mathcal{L};X_{a}...)$ is the Resolution function which smears true $X_{\alpha}...$ to detected $X_{a}...$
- $\epsilon(\mathbf{X_a}...,\mathcal{L}) \text{is the probability that a particle is actually detected by a physical detector.}$

$$B(X_a..., \mathcal{L})$$
 is the background CTEQ2011 Schellman 7/12/11

Efficiency

First pass

Define an "acceptance region" where you would expect your detector to work

(say |η|<3, p_T > 10 GeV)

- Take a particle level MC like pythia
- Reweight kinematics to reflect best knowledge
- Trace particles through detector
- Overlay noise and interactions from other events
- Count how often you reconstruct the particle

The raw simulation isn't good enough

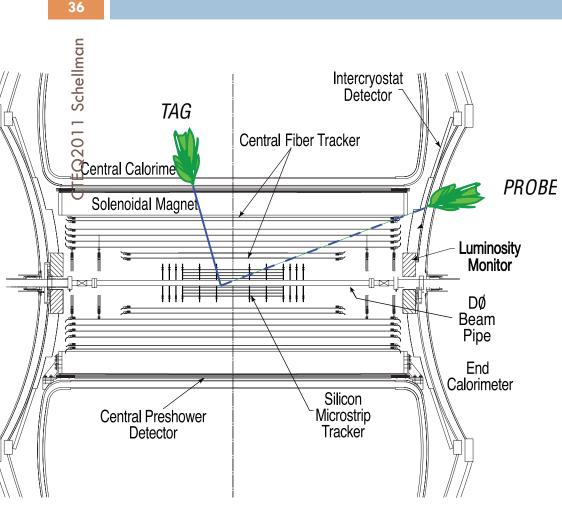
Use scale factors

- Find two ways of measuring something
- Compare the two in data and in simulation
- Correct by the ratio

$$\epsilon(p_T, z) = \frac{(P \cdot T)_{data} / T_{data}}{(P \cdot T)_{sim} / T_{sim}} \epsilon_{sim}(p_T, z)$$

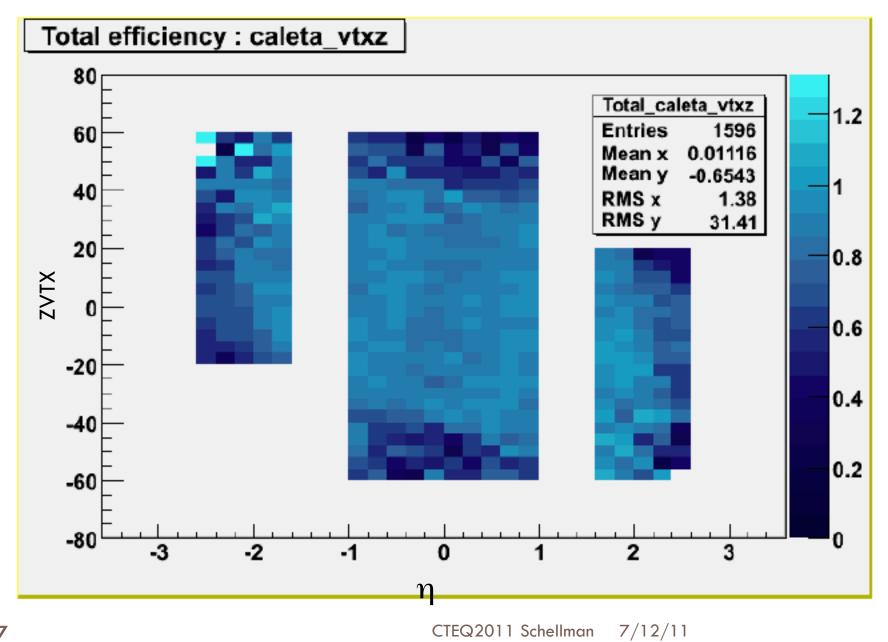
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Tag and probe



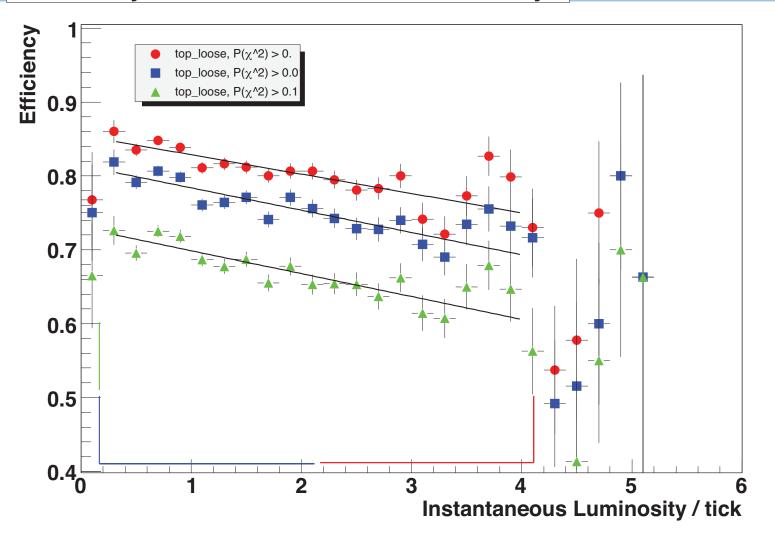
- Find a really good electron (TAG)
- □ Find a track which
- might be from a Z decay (PROBE)
- Did you find a shower

Do the same for a shower (PROBE) that might have a track



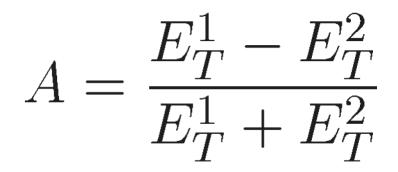
Typical electron efficiencies

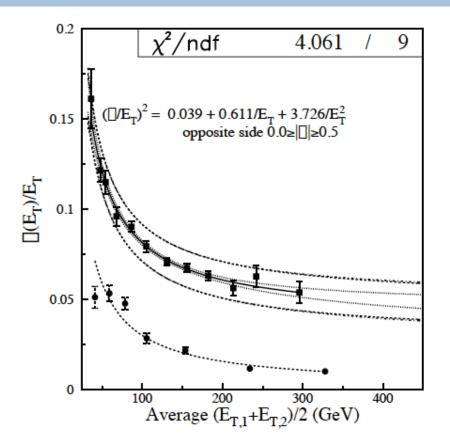
Efficiency versus Instantaneous Luminosity



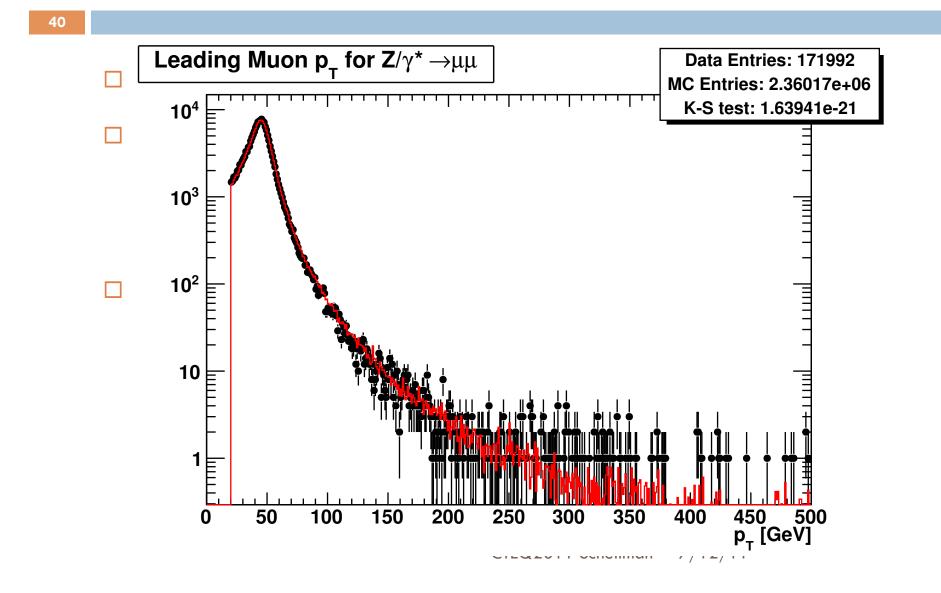
Resolutions

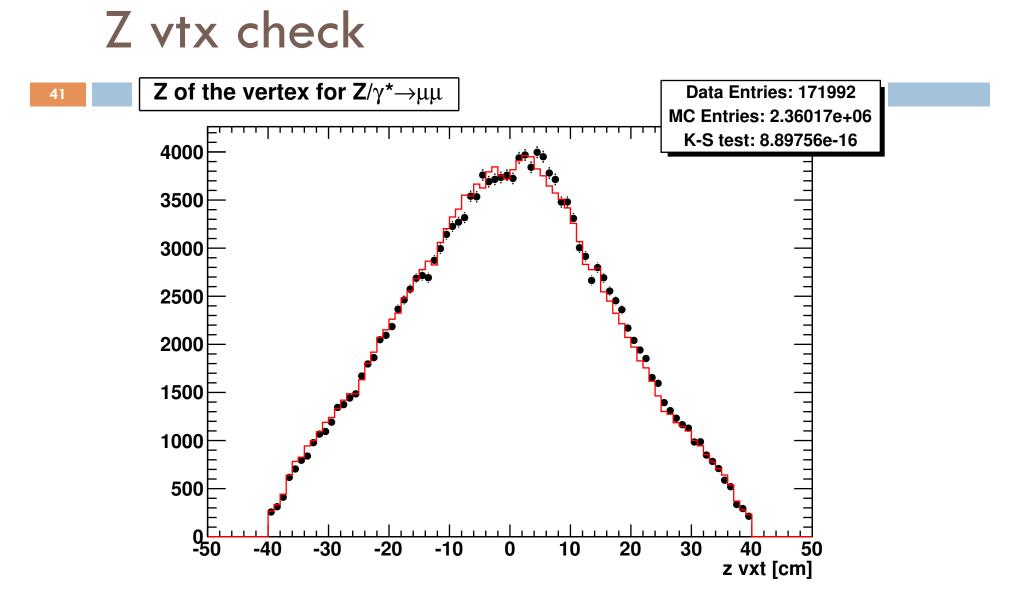
Step 1 – Simulate
Step 2 – Cry in frustration
Step 3 – use the data
Tune to the Z peak width or
Compare energy of back to back jets





Check the tuned full detector simulation

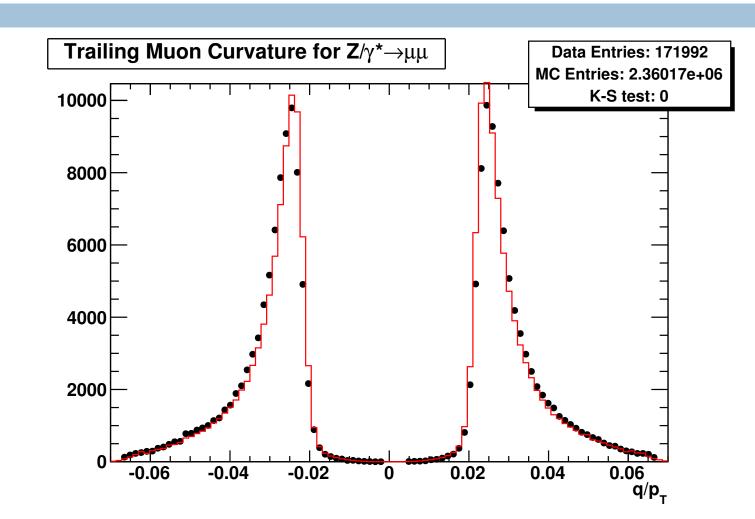




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Curvature check

42



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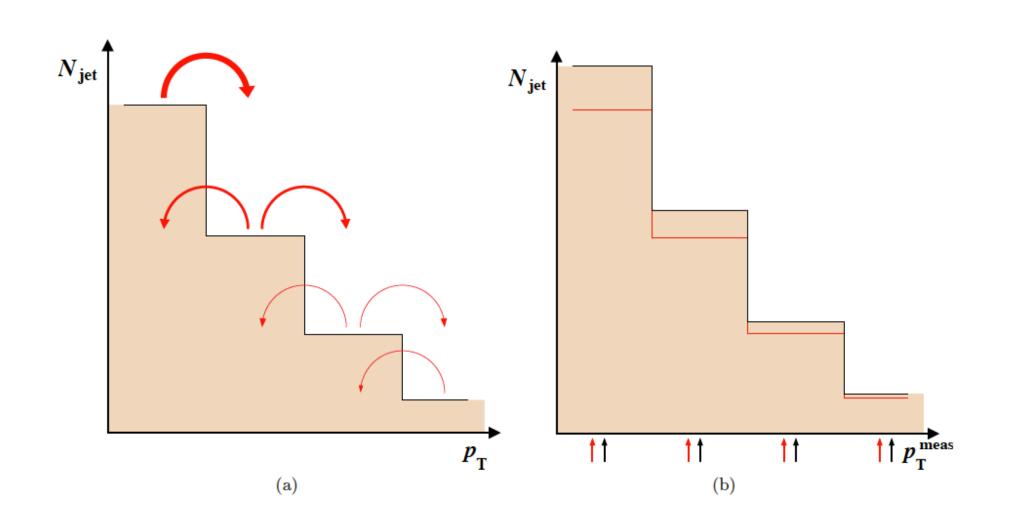
- Ok, looks like we understand our efficiencies and resolutions pretty well in our simulation
- Correct data for backgrounds and efficiency
 Now on to unfolding the resolution effects

Unfolding the resolution effects

$$\begin{split} \mathbf{N_{obs}}(\mathbf{p_T}, \phi, \eta, \mathbf{z}) &= B + \sigma A \epsilon \int \mathcal{L} \mathbf{dt} \\ &= B(p_T, \phi, \eta, z, \mathcal{L}) + \\ &\int \sigma(p_T^0, \phi^0, \eta^0) \times \mathcal{L}(\mathbf{t}, \mathbf{z}) \times A(p_T^0, \phi^0, \eta^0, z^0) \times \\ &\mathbf{R}(\mathbf{p_T^0}, \phi^0, \eta^0, \mathbf{z}^0; \mathbf{p_T}, \phi, \eta, \mathbf{z}, \mathcal{L}) \times \\ &\epsilon(\mathbf{p_T}, \phi, \eta, \mathbf{z}, \mathcal{L}) dp_T^0 d\phi^0 d\eta^0 dt \end{split}$$

- $\sigma(X_{\alpha}...)$ is the true cross section as a function of true variables $X_{a}...$
- $\mathcal{L}(t, z)$ is the luminosity as a function of time.
- $A(X_{\alpha}..)$ is the geometrical Acceptance as a function of true variables.
- $\mathbf{R}(\mathbf{X}_{\alpha}...,\mathcal{L};\mathbf{X}_{\mathbf{a}}...)$ is the Resolution function which smears true $X_{\alpha}...$ to detected $X_{a}...$
- $\epsilon(\mathbf{X_a}..., \mathcal{L})$ is the probability that a particle is actually detected by a physical detector.
- $B(X_a..., \mathcal{L})$ is the background

Resolution Smearing



Unfolding method 1 - matrix

Invert the resolution matrix

$$N_a = B_a + \epsilon_a \mathbf{R}_{a\alpha} A_\alpha \sigma_\alpha \int \mathcal{L}dt$$

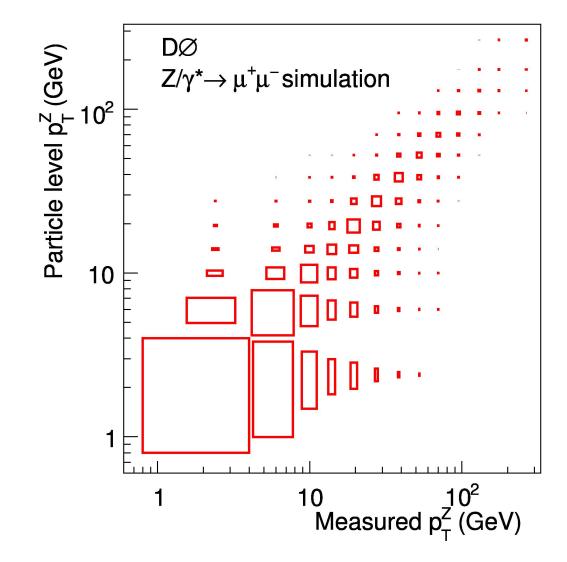
Invert

$$\sigma_{\alpha} = \frac{1}{A_{\alpha}} R^{-1} \alpha a \frac{(N_a - B_a)}{\epsilon_a} \int \mathcal{L} dt$$

Common to regularize the matrix

- Improved stability
- Bias towards smoothing the function

Smearing Matrix $R(X_{\alpha} - X_{\alpha})$



Unfolding method 2 - Ansatz

 \square Convolute a trial unsmeared function with the resolution as a function of (pt, $\eta)$

$$f(p_T, \eta) = N_0 \left(\frac{p_T}{100 \text{ GeV/c}}\right)^{-\alpha} \left(1 - \frac{2p_T \cosh(y_{\min})}{\sqrt{s}}\right)^{\beta} \exp\left(-\gamma p_T\right).$$

- □ Fit this convoluted "smeared" function to your data
- Correct data by the ratio of the unsmeared to smeared "ansatz" function.

$$R(X_{\alpha} \to X_{a}) = \delta_{\alpha a} \frac{\text{Smeared}(X_{a})}{\text{Unsmeared}(X_{\alpha})}$$

This works if a simple functional form can fit your data well and the resolution function is well understood.

Ansatz fit $R(X_{\alpha} \to X_{a}) = \delta_{\alpha a} \frac{\text{Smeared}(X_{a})}{\text{Unsmeared}(X_{\alpha})}$

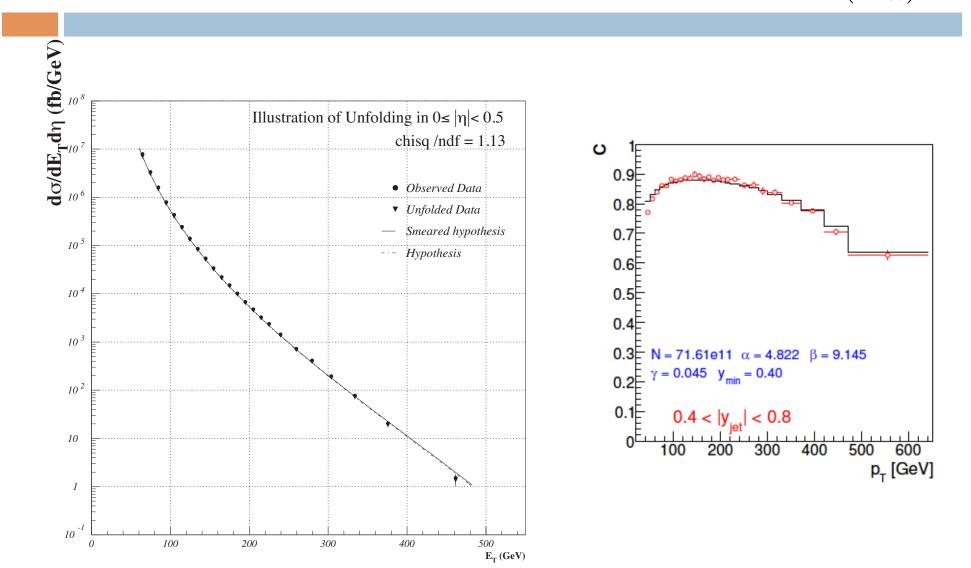
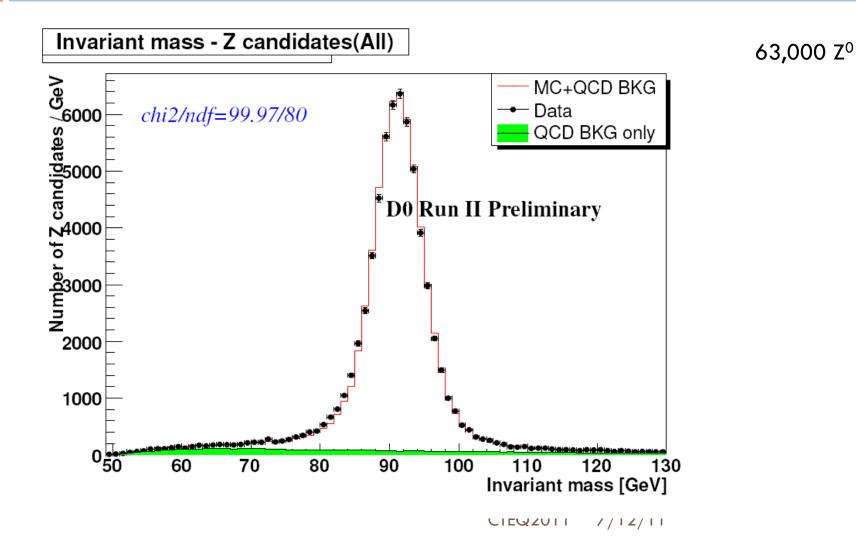


Figure 49: Illustration of the data unfolding procedure in the rapidity region $|\eta| < 0.5$.

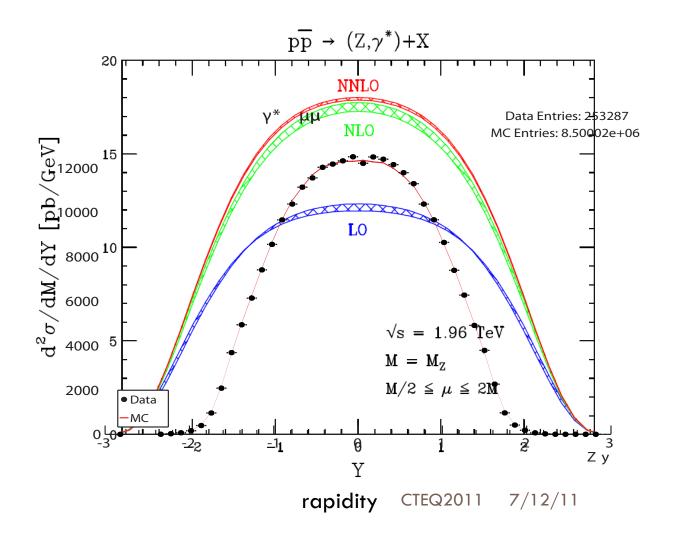
Now for the final step: Acceptance

- Correct from the region we cover to the rest of phase space
- Why do this instead of including in the "efficiency"?
 - This correction is big factor of 3 for Z production $\epsilon \sim 40\%$, A ~ 30%
 - Only depends on generator level quantities
 - Can use the best NNLO simulations

D0 1 fb⁻¹ sample



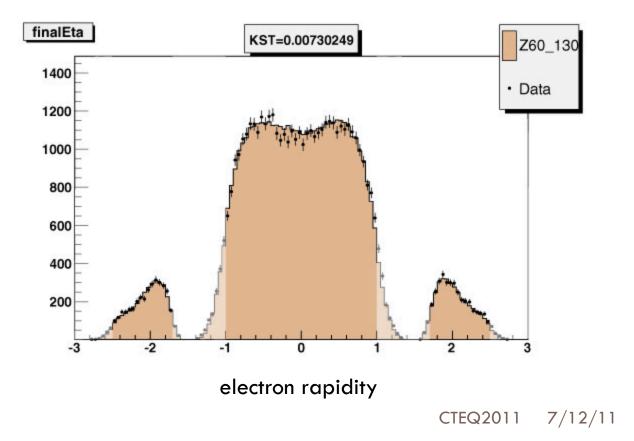
Z Rapidity raw and generated



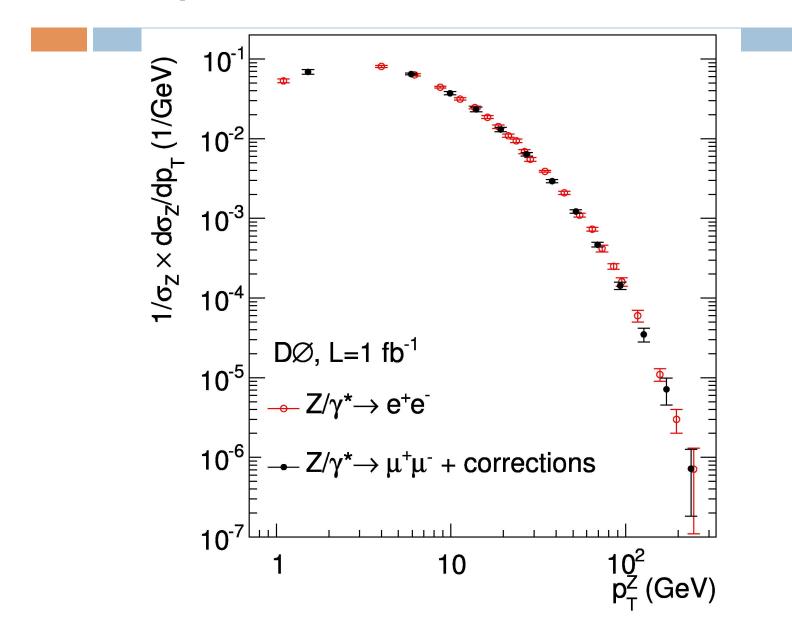
We actually define acceptance as $|\eta_e|<1$ or 1.5 $<|\eta_e|<$ 2.4 and $p_T>$ 25 GeV

53

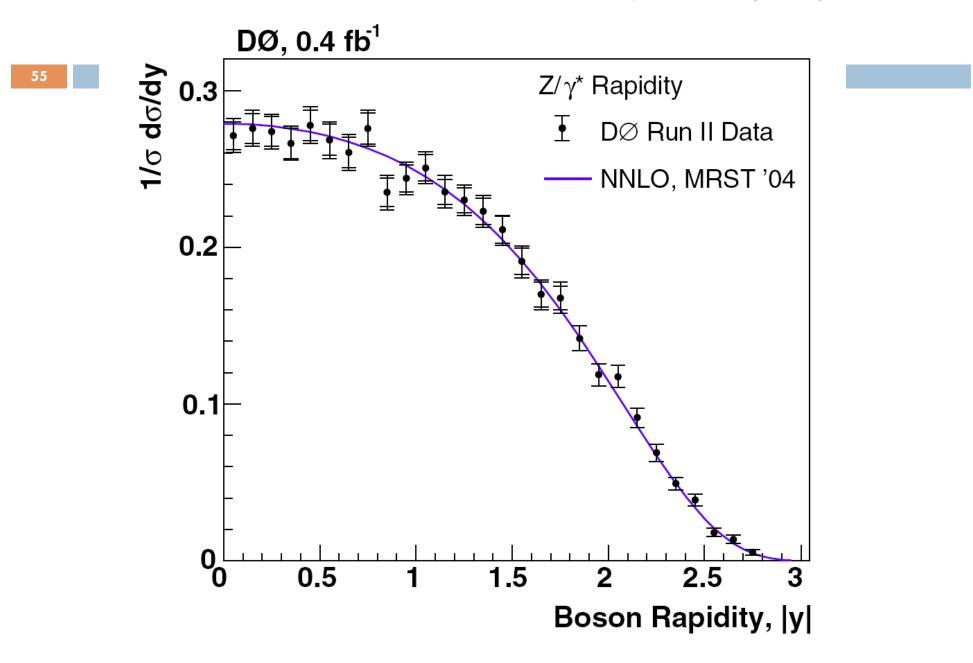
Cut cross section in electron rapidity and pt



Fully corrected normalized distributions



PHYSICAL REVIEW D 76, 012003 (2007)



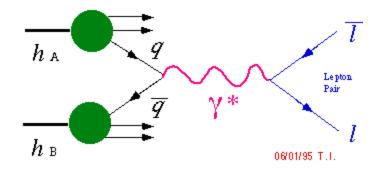


56 $\frac{d\sigma^{h_1h_2}}{dQ^2 dp_T^2 dy \, d\Omega^*} = \sum_{a,b} \int dx_1 \, dx_2 \left(f_a^{h_1}(x_1, \mu_F^2) f_b^{h_2}(x_2, \mu_F^2) \right) \frac{s \, d\hat{\sigma}_{ab}}{dQ^2 \, dt \, du \, d\Omega^*} \left(x_1 P_1, x_2 P_2, \alpha_s(\mu_R^2) \right)$

Parton density functions

Hard cross section

The Drell-Yan Process



Higher order QCD is important

 $gq \rightarrow Zq, qq \rightarrow Zg$ etc.

Cross section*B(Z→ee) LO ~ 180 pb NLO ~ 250 pb NNLO ~ 260 pb CTEQ2011 7/12/11

$Z \rightarrow ee Errors$

57	Statistical	0.62%
	Preselection efficiency	0.85%
	Radiative corrections	<0.5%
	ID eff stat error	0.4%
	Tag-probe bias	0.3%
	Noise corrections	0.22%
	Vertex z	0.6%
	Cut variations	1.5%
	Total systematic error	2.0%
	PDF error on sigma(tot) Luminosity	+1.3% -1.7% 6.1% -1.7%

More fun

- Radiative corrections
- Calibration
- Jet definitions
- B-tagging

□ My list of things to check is around 80 items long